Decentralized collaborative transport of fabrics using micro-UAVs

Ryan Cotsakis, David St-Onge and Giovanni Beltrame

Abstract—Small unmanned aerial vehicles (UAVs) have generally little capacity to carry payloads. Through collaboration, the UAVs can increase their joint payload capacity and carry more significant loads. For maximum flexibility to dynamic and unstructured environments and task demands, we propose a fully decentralized control infrastructure based on a swarm-specific scripting language, Buzz. In this paper, we describe the control infrastructure and use it to compare two algorithms for collaborative transport: field potentials and spring-damper. We test the performance of our approach with a fleet of micro-UAVs, demonstrating the potential of decentralized control for collaborative transport.

I. INTRODUCTION

Transport tasks for industrial, commercial and medical applications are now considering small-scale aerial technologies ready for deployment. New use cases are emerging, but the available payload remains the main limitation for their implementation. Quite a novel use case was presented to our team from a fashion designer: to involve micro-UAVs in the domestic task of getting dressed. This proposal is challenging at many levels, namely: interaction, clothing design, and UAV control. In this paper, we address the first piece of the puzzle: how to use very small and therefore safe devices to transport light pieces of clothing. Indeed, the proximity of the user and the narrow space of a dressing room prevent the use of large devices. However, a decrease in size tends to considerably decrease the payload. We propose to base our solution on the use of multiple micro-UAVs collaborating as a whole for the transport task, as shown in Fig. 1.

The nature of the payload is itself a challenge for aerial transport. A flexible material supported from many points creates a complex dynamic between the carriers [1]. Solving this problem has many other useful applications: supporting a net to catch other UAVs, transporting hoses to drop water over fires, or slings for rescue missions on mountains have comparable dynamics with bigger payload requirements. Nets, fabrics, hoses and slings, can all be approximated with one or many deformable linear objects (DLOs).

Since suspended fabrics are heavily influenced by the motion of air, the forces on the UAVs are largely unpredictable, and can be quite large in magnitude. In this case, the UAVs that share the payload must compensate for any disturbances in real time. This can be achieved by having the UAVs cooperate, and adjust their movement according to the positions and dynamics (acceleration, torque, etc.) of

M. Cotsakis is an undergraduate student from the University of British Colombia, 2329 West Mall, Vancouver, BC. Dr. St-Onge and Prof. Beltrame are with the Department of Computer and Software Engineering, Polytechnique Montréal, 2900 Boul Édouard-Montpetit, Québec CA. e-mail: (name.surname@polymtl.ca).



Fig. 1: Three Crazyflie micro-UAVs collaborating to transport a white fabric without requiring any central controller.

their neighbours. Swarm intelligence has proven that it is possible to construct robust behaviours from local interaction with the robot neighbours only [2]. For payload transport, this approach leads to less processing power required to fully embed the control. The decentralized nature of swarms also removes the weakness of a central control station, an advantage to many critical scenarios [3].

This paper presents adapted Lennard-Jones and Spring-Damper decentralized control algorithms including a DLO model and compares their performance in experiments. It is structured as follows: we discuss inspiring works in Sec. II, then we detail our decentralized architecture in Sec. III and we derive the control algorithms in Sec. IV. Finally, the results of a set of experiments conducted with Crazyflie micro-UAVs are presented and discussed in Sec. V.

II. RELATED WORKS

The transport of DLOs has been recently studied [1]. The feasibility of this task, with three UAVs, was demonstrated in simulations using a particle swarm optimization to get the proper set of PID gains before flight. This work, as with most of the literature of flexible payload transport with UAVs, adopted a centralized controller. Similarly, the impressive body of work of D'Andrea's group at ETH Zurich includes works such three UAVs throwing and catching a ball in a net [4] and the transport of a flexible ring with six UAVs [5]; but they rely on a fully centralized and latency-free state estimator and controller. In realistic field deployments, we can't rely on such an ideal state estimator (GPS noise, lack of resources to deploy beacons, etc.). Furthermore, while centralized approaches have more history in robust control theory, aerial transport benefits from relaxing the need for a reliable communication link with a central ground station.

Distributed or decentralized approaches to the problem are seen less frequently. Distributed collaborative transport was achieved for wheeled robots, pushing a rigid geometric

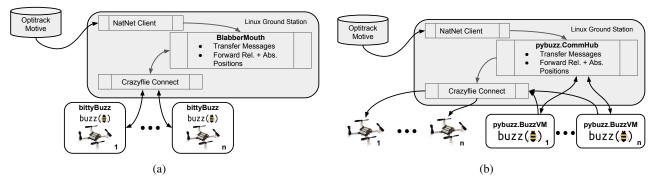


Fig. 2: Buzz deployment architecture for the Crazyflie: (a) Buzz running on-board thanks to its micro-controller ultra-lightweight version: bittyBuzz and the ground station emulating neighbours interactions (relative localisation and communication). (b) Buzz emulated using PyBuzz: similar to (a) but emulating also the onboard Buzz control with multiple instance of pybuzz on the ground station.

object [6], [7] with a combination of potential field forces, similar to the Lennard-Jones algorithm adapted for this work. More recently, two UAVs were designed to cooperate in a fully decentralized configuration for the transport of a rigid body using only inertial measurement and vision [8]. Both are inspiring works, but are still far from the DLO transport targeted by this work.

The same type of micro-UAVs used in our experiments, the Crazyflie, was subject to a detailed swarm controller design [9]. We leverage the stability of their on-board control and the scalability of their configuration to deploy our solution. Nevertheless, collaborative payload transport has never been attempted with such a small and sensitive device. Other solutions for controlling a swarm of Crazyflies have been considered [10], [11], but they are centralized. As with Preiss et al., their scripts are specific to their implementation for this platform, whereas we leveraged a domain-specific programming language for swarms [12] to guarantee portability to other platforms.

III. DECENTRALIZED ARCHITECTURE

The aim of this project is not only to achieve a fleet of Crazyflies collaborating on a transport task, but to study a behaviour that can then be ported to other robotic platforms. The Buzz domain-specific language [12] and its associated virtual machine was successfully deployed on DJI Matrice 100, Intel Aero, and Clearpath Husky to name a few.

A. Buzz

To accelerate the implementation of swarm behaviours, Buzz provides a set of special primitives from which we leverage two essential concepts: a) swarm aggregation and b) neighbour operations. We detail here these concepts to ease the comprehension of the following implementation, but all information can be found in previous publications specific to Buzz [12] and example scripts that are available online¹.

Swarm Aggregation is a primitive which allows for grouping of robots into sub-swarms, through the principle of dynamic labelling [12]. The swarm construct is used to create a group of robots which can be attributed with a specific

1http://the.swarming.buzz

behaviour that differs from the other robots, based either on the task or robot abilities.

Neighbour Operations in Buzz refer to a rich set of functions which can be performed with or on neighbouring robots through situated communication [?]. Neighbours are defined from a network perspective as robots which have a direct communication link with each other. With situated communication, whenever a robot receives a message, the origin position of the message is also known to the receiver.

Finally, any Buzz script is compiled into an optimized, memory-efficient, and platform-agnostic bytecode to be executed on the Buzz Virtual Machine (BVM). To interface the BVM with the robots' actuators and sensors, the integrator needs to write C hooks that are callable from a Buzz script.

B. PyBuzz

Rather than embedding the BVM into the firmware of each robot, we decided to accelerate the development of the control algorithms by creating a centralized emulator for our decentralized configuration. Emulating a distributed software architecture requires a wrapper that can be instantiated to connect to each hardware node. To solve this problem, we created a Python module for wrapping the BVM called PyBuzz, such that in Python, one can construct a BVM as a Python object, and link Python functions as callable functions in Buzz. When interpreting the Buzz object code, the BVM performs calls to these Python functions.

By importing pybuzz in a Python script, one can construct any number of BVM objects. The translation between C and Python was achieved using Cython [13], a tool that is used for writing C extensions for Python.

On a ground station, we have information about the locations of all of the UAVs. These locations are fed into a communication hub, pybuzz.CommHub, as shown in Fig. 2-b. The communication hub emulates the robots' communication range, giving each BVM its current location, and the relative locations of its neighbours. While each BVM continuously steps through the Buzz script, the communication hub automatically manages the transmission of all messages sent between each BVM. The only information sent to the Crazyflie's hardware is its current location (received from

pybuzz.CommHub) and instructions for course adjustments (received from pybuzz.BuzzVM).

In the end, a Buzz script running on the PyBuzz emulator (Fig. 2-b) can work unmodified on a Crazyflie (Fig. 2-a).

IV. DECENTRALIZED CONTROL

The accurate control of a fleet carrying a shared payload is a problem usually approached from a robust low-level control perspective [5]. However, modelling the payload between the robots as a set of local interactions can bring it up to the fleet coordination level, for which decentralized solutions are widely studied in swarm behaviours.

A. Flexible payload model

We take inspiration from the modelling of hoses, approximated for UAVs collaborative payload with discrete linear objects [14]. We also postulate that the transport is achieved only with UAVs at cruise heights, i.e. with the fabric freely hanging. However, our experiments show that this approximation can hold for stable take-off and landing while holding the fabric in certain conditions (see Sec. V).

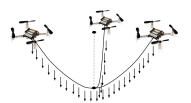


Fig. 3: Modelling fabric transport by connected DLOs. Each of the three Crazyflies carries half of a catenary curve with its vertex at a point directly below the centre of mass of the swarm

The payload can be approximated by cables with uniformly distributed weight creating catenary curves that hang from each robot to a point directly below the centre of mass of the swarm, which we will refer to as the *connecting point*. The vertex of each catenary curve coincides with the connecting point. This model is depicted in Fig. 3. With n robots all at the same height sharing a payload of mass m, the vertical tensile force T_z is approximately

$$T_z = \frac{mg}{n} \tag{1}$$

The horizontal force T_x is a much more involved calculation. We recall the equation of a catenary curve to be

$$z(x) = a \cosh \frac{x}{a} \tag{2}$$

where a is a parameter that describes how stretched the approximating cable is. Remaining under the assumption that the robots have the same altitude, a can be solved for numerically in the transcendental equation

$$L = a \sinh \frac{x_0}{a} \tag{3}$$

where L is the length of the approximating cable from the robot to the connecting point, and x_0 is the horizontal

distance from the robot to the centre of mass of the swarm. If we take the derivative of Eqn. 2, we see that

$$\frac{T_z}{T_x} = z'(x_0) = \sinh\frac{x_0}{a} = \frac{L}{a} \tag{4}$$

If we observe through Eqn. 3 that a is a function of x_0 , we can see that T_x is a function of x_0 and other constants:

$$T_x(x_0) = \frac{mg}{nL}a\tag{5}$$

B. Spring-Damper analogy

An intuitive method of having the UAVs maintain a stable formation is to simulate the forces due to springs and dampers emplaced between each robot. Spring and damper control was shown to work well in a simulated swarm [15]. The repulsive force between two robots as a function of the distance d between them would be

$$F_{ren} = -k(d - l_0) - B\dot{d} \tag{6}$$

where k and l_0 are respectively the spring constant and unstretched length of the spring, and B is the damping constant of the damper.

To tune the parameters k and B, recall the horizontal force on each robot due to the payload from Eqn. 5. Let us define k_p to be the value of T_x' at the equilibrium position of x_0 . Then for small perturbations from the equilibrium position, the horizontal force due to the payload can be approximated by a spring with spring constant k_p connecting the centre of mass of the swarm to the robot. The centre of mass is computed in a decentralized way, by each robot based on the relative position of its neighbours. By relating the payload to a spring force, we can interpret the entire system as a sum of forces from springs and dampers to simplify the analysis.

We can calculate the net force on a particular robot in the x direction, the direction towards the centre of mass of the swarm, for a small change Δx from the equilibrium position. We denote this force as F_x and express it as a sum over the robot's N neighbours:

$$F_x = -\left(k_p + k \sum_{i=1}^{N} \cos^2 \theta_i\right) \Delta x \tag{7}$$

where θ_i is the angle from the centre of mass of the swarm to neighbour i with respect to the robot in question. In the same way we can calculate the force in the y direction resulting from a small perturbation Δy to be

$$F_y = -\left(k\sum_{i=1}^N \sin^2\theta_i\right) \Delta y \tag{8}$$

By substituting equations 7 and 8 with the values

$$k_x = k_p + k \sum_{i=1}^{N} \cos^2 \theta_i, \quad k_y = k \sum_{i=1}^{N} \sin^2 \theta_i$$
 (9)

we are able to simplify the system of springs down to two springs. One parallel to the x-axis with spring constant k_x , and one parallel to the y-axis with spring constant k_y . What remains is to simplify the system of dampers.

Similarly to k_x and k_y in Eqn. 9, we define

$$B_x = B \sum_{i=1}^{N} \cos^2 \theta_i, \quad B_y = B \sum_{i=1}^{N} \sin^2 \theta_i$$
 (10)

and interpret our system separately in dimensions x and y when choosing parameters B and k. In the x direction we have, attached to the Crazyflie's mass, a spring and a damper in parallel with parameters k_x and B_x respectively. In the y direction, we have a spring and damper in parallel with parameters k_y and B_y respectively. We tune B and k such that both of these systems in each dimension are nearly critically damped.

The calculation of \dot{d} in Eqn. 6 is not a straight forward task when d is being sampled discretely in time and space. To mitigate the errors induced by the lack of continuity, we use all n previously measured values of d, and weigh each measurement's importance with a coefficient that's magnitude decays exponentially with time. The simple implementation of this algorithm is to initialize another distance variable \tilde{d} that acts as a representative of all previous values of d.

Given that at time t_n we measure the distance between robots to be d_n , we define

$$\tilde{d}_n = w_n \tilde{d}_{n-1} + (1 - w_n) d_n, \text{ with } \tilde{d}_0 = d_0$$

$$w_n = \exp\left(-\frac{t_n - t_{n-1}}{\tau}\right)$$
(11)

where τ is a time constant that governs how quickly information from previous measurements is lost. We aim to show that we can approximate the velocity as follows:

$$\dot{d} = \frac{d_n - \tilde{d}_n}{\tau} \tag{12}$$

In the case where d is measured at regular time intervals Δt , then $w_n = w = \exp(-\Delta t/\tau)$, $\forall n$ which gives a better approximation:

$$\dot{d} = \frac{(1-w)(d_n - \tilde{d}_n)}{w\Delta t} \tag{13}$$

To show that Eqn. 13 is a reasonable approximation, let us consider the explicit form of \tilde{d}_n :

$$\tilde{d}_n = (1 - w) \sum_{i=0}^{\infty} d_{n-i} w^i$$
 (14)

Since we aim to obtain an expression for d, let us assume that it is relatively constant (at least for a short interval of time). This allows us to make the approximation

$$d_{n-i} \approx d_n - i\Delta t\dot{d} \tag{15}$$

With this approximation, Eqn. 14 can be expressed as

$$\tilde{d}_n = (1 - w) \sum_{i=0}^{\infty} (d_n - i\Delta t \dot{d}) w^i$$
 (16)

which can be evaluated exactly:

$$\tilde{d}_n = d_n - \frac{w\Delta t\dot{d}}{1 - w} \tag{17}$$

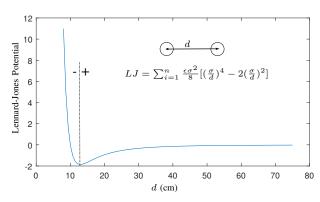


Fig. 4: Lennard-Jones potential adapted for Crazyflies. The 'and '+' domains are respectively the repulsive and attractive parts, for which the pivot point is set with the parameter σ . d is the distance between two robots and ϵ is a parameter acting as a control gain on the potential.

Finally, solving for \dot{d} yields Eqn. 13.

With w being a exponential function of $\frac{\Delta t}{\tau}$, it can easily be shown that as $\frac{\Delta t}{\tau} \to 0$,

$$\frac{1-w}{w\Delta t} = \frac{1}{\tau} \tag{18}$$

which when substituting in Eqn. 13 implies Eqn. 12. One downside to using Eqn. 12 is that we require $\tau\gg\Delta t$, and since \dot{d} is an approximation of the velocity $\frac{\tau}{2}$ seconds in the past, our measurement of \dot{d} is delayed. The lack of w and Δt dependence in Eqn. 12 implies that it can be used over Eqn. 13 whenever d is sampled irregularly, as long as τ is sufficiently large.

C. Bacteria interaction analogy

Among the most popular formalizations of biological swarm behaviours, potential functions are a simple, yet flexible control approach. Artificial potential functions have been used extensively for robot navigation and control [16], [17]. Based on their knowledge of their neighbours' positions, each robot computes a virtual force vector:

$$F = \sum_{i=1}^{N} f(d_i)e^{j\theta_i}$$
(19)

where θ_i and d_i are the direction and the distance to the ith perceived obstacle or robot, and the function $f(d_i)$ is the negative gradient of an artificial potential function. We chose the commonly used Lennard-Jones potential and adapted it for our physical system as shown in Fig. 4.

The two parts of the potential equation represent the attractor and repulsor effect. This potential is driven by two parameters: the target distance between robots σ and the strength ϵ of the potential.

Based on its popularity, we selected Lennard-Jones potential to control the behaviour of our robotic swarm:

$$LJ = \frac{\epsilon \sigma^2}{8} \left[\left(\frac{\sigma}{d} \right)^4 - 2 \left(\frac{\sigma}{d} \right)^2 \right] \tag{20}$$

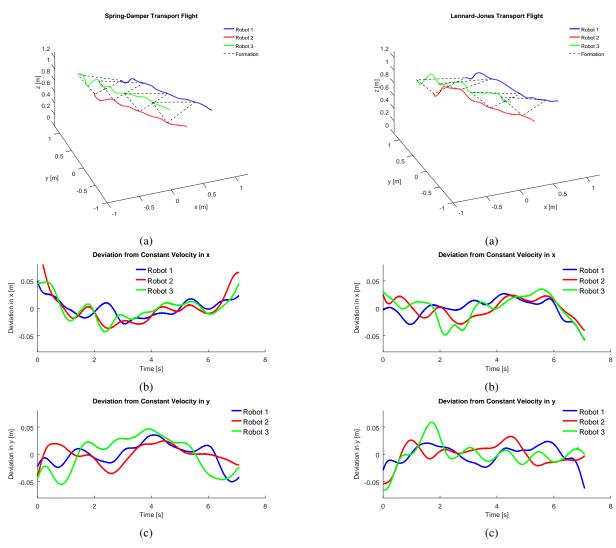


Fig. 5: **Spring-damper** flight test: (a) Overview of flight paths for the three Crazyflies while cooperatively carrying fabric during one of the five flight tests of its kind. The vertices of the dotted triangles indicate where the Crazyflies are at two second intervals. (b) Deviation of each robot's x coordinate from the best approximating constant velocity trajectory for the curves in (a). Each curve is a plot of x vs t with the zeroth and first order terms subtracted, i.e. the average value and average slope of each curve is 0. (c) The same construction as (b), but taking the y vs t graph, rather than the x vs t.

We can compute the resulting force exerted as the negative gradient of Eqn. 20:

$$F_{rep} = \frac{\epsilon \sigma}{2} \left[\left(\frac{\sigma}{d} \right)^5 - \left(\frac{\sigma}{d} \right)^3 \right]$$
 (21)

This representation of the Lennard-Jones repulsive force has a very elegant Taylor expansion about $d = \sigma$, that is

$$F_{rep} = -\epsilon(d - \sigma) + O\left((d - \sigma)^2\right) \tag{22}$$

The first term in Eqn. 22 has the same form as the first term in Eqn. 6. With $\epsilon=k$ and $\sigma=l_0$, the Lennard-Jones

Fig. 6: Lennard-Jones flight test: (a), (b), (c) See Fig. 5 for description of these figures.

potential is, to second order, identical to a spring system at the equilibrium position and can be approximated by a harmonic oscillator. Representing the potential in this form, we are able to compare Lennard-Jones and spring-damper experiments using the parameters derived in Sec. IV-B.

D. Fleet translation

With both control approaches, a goal (target location) is represented as an attractor influencing the whole group. The final displacement vector at each step is then computed from a weighted sum of this attraction force and the forces resulting from the formation algorithm. Under ideal circumstances, the robots move towards their final position at a constant velocity.

V. EXPERIMENTS

To start, we devised an unorthodox method to have a swarm of Crazyflies be seen and distinguished by our motion sensing platform, Optitrack, since the UAVs are too small to allow for many unique configurations of reflective markers. Our solution was to attach a single marker on each Crazyflie, which provides the position of the UAV in 3-D space, but not the orientation. We are able to control the attitude of the Crazyflie with the on-board controller using its Inertial Measurement Unit. To have the robots be uniquely identifiable, we tell the software approximately where each Crazyflie is expected to start, and conduct a grid search before takeoff. For the duration of the flight, we assume continuous motion of the Crazyflie. This way, the software can identify which robot is which, based on the previous frame.

Building on the infrastructure work described in Sec. III, we implemented Buzz scripts for both algorithms. We conducted eleven flight tests, each involving three Crazyflies. The first of these tests demonstrates the flight of the three Crazyflies without any payload, nor any control algorithm facilitating the distances between robots. Five of the flight tests demonstrate the transport of fabric using the spring-damper algorithm, and the last five demonstrate the transport of fabric using the Lennard-Jones algorithm. All other variables were kept constant. In each flight test, the Crazyflies were instructed to take off from slightly elevated platforms to 80 cm and maintain this altitude while travelling horizontally at a constant velocity. For the tests involving the transport of fabric, one of the corners of the fabric was manually tied to each Crazyflie with fishing line before takeoff.

A. Weight limit

From BitCraze specifications, a Crazyflie can lift 42 g and its own weight is 27 g, leaving 15 g for a payload, from which we already use 0.1 g for an Optitrack marker. In order to confirm and test these numbers, we measured the maximum lifting capability of the Crazyflies. Using a digital spring scale with a resolution of 1 g, we took 5 measurements of the maximum payload, and all measurements gave 18 g.

B. Spring-damper control versus Lennard-Jones potential

Post-experiment, we analyzed the robustness and stability of the collaborative transport algorithms, and quantified the cooperation between robots. For each flight test, we took seven seconds of linear translation data, which is depicted in Fig. 5 and Fig. 6. Fig. 5-a is the flight trajectory of one of the collaborative fabric transport tests using a spring-damper algorithm. Fig. 6-a is the flight trajectory of one of the collaborative fabric transport tests using a Lennard-Jones potential. In a decentralized system, the behaviour of one robot can have great impacts on all of its neighbors. In particular, with non-deterministic forces arising from the flexible payload, it is expected that each robot must dynamically accommodate for its neighbours when they undergo instability. Thus, in a statistical analysis we should be able to measure these adjustments.

Parts (b) and (c) of figures 5 and 6 illustrate the instabilities of each robot in a single dimension with respect to time. Without a shared payload the movement is expected to be directed only towards the goal, so these plots show the deviation from that constant velocity trajectory.

It can be observed that there is a similarity between the curves of Fig. 5, and of Fig. 6. Synchronicity of UAV

TABLE I: Comparison of correlation coefficients.

	$\bar{\rho_x}$	$\bar{\rho_y}$	$\frac{\bar{\rho_x} + \bar{\rho_y}}{\bar{\rho_x}}$
No Fabric	0.19	-0.17	0.01
Spring- Damper	0.74	0.44	0.59
	0.66	0.31	0.48
	0.60	-0.02	0.29
	0.31	0.03	0.17
	0.29	-0.06	0.12
Lennard- Jones	0.47	0.67	0.57
	0.43	-0.08	0.17
	0.68	0.20	0.44
	0.55	0.28	0.41
	0.26	-0.06	0.10

movements is ideal in this decentralized system, and it can be quantified by the average of the correlation coefficients of each pair of signals. The correlation between Robot i and Robot j in one dimension is calculated as

$$\rho_{ij} = \frac{\delta_{\mathbf{i}} \cdot \delta_{\mathbf{j}}}{\sqrt{(\delta_{\mathbf{i}} \cdot \delta_{\mathbf{i}})(\delta_{\mathbf{j}} \cdot \delta_{\mathbf{j}})}}$$
(23)

where δ_i is vector of distances of Robot *i* from its constant velocity trajectory at various times. Finally,

$$\bar{\rho} = \frac{1}{3}(\rho_{12} + \rho_{23} + \rho_{13}) \tag{24}$$

 $\bar{\rho}$ was calculated for each of the two dimensions x and y, and for each flight test. The results are presented in Tab. I.

For the Spring-Damper algorithm, the mean μ and the standard deviation σ of the correlation coefficients are:

$$\mu_{SD} = 0.33, \quad \sigma_{SD} = 0.20$$
 (25)

while for the Lennard-Jones algorithm,

$$\mu_{LJ} = 0.34, \quad \sigma_{LJ} = 0.20$$
 (26)

The experiments show that the spring-damper and Lennard-Jones algorithms do in fact allow the robots to maintain safe, stable distances with similar efficacy. The non-linearity of the Lennard-Jones potential induces the same amount of cooperation between robots as the dampers that are present in the spring-damper control algorithm. We successfully tried the same algorithms with four micro-UAVs, as shown in the attached video.

VI. FUTURE WORK

This work is a first step along the way to get a safe fleet of micro-UAVs to dress a human. The decentralized approach was demonstrated to work well for such a scenario. Leveraging PyBuzz, other decentralized behaviours will be tested to succeed in creating autonomous flying dressing aid. In the future, we will try various shapes of clothing, designed specifically to be transported with micro-UAVs. We will also work on enhancing the robustness of the algorithm and on adding infrared proximity detection to add user interaction to the system.

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