

1 **Comparison of wind turbine blade structural models of different levels of complexity**
2 **against experimental data**

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8 As the design process of a wind turbine blade is highly iterative, one needs to do the
9 same calculations several times. During that process, the kind of structural model
10 use must be chosen carefully trying to obtain a good compromise between precision
11 and model setup and computation time. This paper compares four blades structural
12 models of different levels of complexity. These models are compared to each other
13 and with experimental results with respect to their abilities to analyze blade cross-
14 sectional properties, natural frequencies, deflection, strains, buckling strength and
15 composite strength. This comparison shows that even if the 3D shell finite element
16 model is the more precise and is the only one that can manage the regions of the blade
17 where the cross-sectional shape changes quickly, strength of material based models
18 give accurate results. Even the simpler model, based on blade shape simplification,
19 gives conservative and accurate results at a very low computational cost.

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20 I. INTRODUCTION

21 Like many other components, the structural design process of a wind turbine blade is an
22 iterative process. From one iteration to another, the structural designer has to provide blade
23 cross-sectional properties (for instance, \overline{EA} , \overline{EI} , mass per unit length) to the aeroelastic
24 analysis model and to perform the blade structural dimensioning and validation using the
25 loads outputted by the aeroelastic model. Either of those can come first depending on the
26 information available and the hypothesis used.

27 As a few iterations can be needed to reach the final design, the tools and models used
28 may not be the same at each iteration. For example, the models needed for final iterations
29 are not necessarily well suited for preliminary stages of the design process where simple and
30 easy-to-solve models are generally preferred.

31 Two types of structural models are mainly used: beam models and 3D shell finite element
32 models. Beam models are based on the strength of material theory. They are the type of
33 model used in most of the aeroelastic codes, and they are well suited and often used for the
34 blade cross-sectional property evaluation. However, they can also yield valuable information
35 for the blade structural validation. Their capabilities for evaluating stresses and strains,
36 deflection, buckling, and eigenfrequencies are not as good as those of a 3D finite element
37 model (especially in the areas of rapid cross section change like maximum chord to root
38 transition region), but they are sufficient for preliminary dimensioning and validation and
39 they are much faster to set up and get results than 3D shell finite element models. It is also
40 interesting to note that the most recent aeroelastic codes tend to include beam models that
41 are able to manage all the material and geometric coupling between the different deformation
42 modes of general composite beams. Therefore, complex cross-sectional analysis tools are
43 needed in order to get all the required cross-sectional properties.

44 3D shell finite element models are based on surface elements with nodes having transla-
45 tional and rotational degrees of freedom. Commercial software offers shell finite elements
46 with through the thickness layered material definition that are well suited for composite
47 materials. This type of model is often used for the structural dimensioning and validation
48 of the blade. However, with proper post-processing, a 3D shell finite element model can also
49 be used to get the blade cross-sectional properties¹.

50 The blade designer then has to choose the right structural analysis model to use in a

51 given situation. That being said, the purpose of this article is to compare wind turbine
52 blade structural models of different levels of complexity with respect to their capability of
53 blade design and analysis and conclude on their use at the different stages of the blade design
54 process. These models are compared with each other and with experimental data.

55 The blade used as reference and the test setup are first presented. Then, four different
56 blade structural models, from the simplest to the most complex, are described. Finally, the
57 comparison is made about blade cross-sectional properties, natural frequencies, deflection,
58 strains, buckling strength and composite strength.

59 II. WESNET BLADE AND TEST METHODOLOGY

60 The blades used for the experimental part of this article are those of a wind turbine
61 designed at École de technologie supérieure and manufactured in the course of a project
62 funded by the Natural Sciences and Engineering Research Council of Canada's Wind En-
63 ergy Strategic Network (NSERC/WESNet). This 8.08 m rotor diameter turbine reaches its
64 nominal power of 10 kW at a wind speed of 9.5 m/s. It uses the direct-drive technology
65 and is controlled using variable speed and active pitch. The rotor auto-aligns with the wind
66 direction due to a free yaw and downwind rotor configuration.

67 The 3.76 m long blades are made of epoxy-glass fiber composite and foam core. A
68 schematic representation of the blade cross section is presented in Figure 1. The upper
69 and lower surfaces of the airfoil and the shear web are bonded together. Both aerodynamic
70 shells are thicker in the maximum thickness region of the airfoil (between 15 % and 45 %
71 of chord length) to form the spar caps that support most of the blade loads. The blade
72 external geometry is presented in Table I and the blade composite layup, in Table II. The
73 longitudinal positions of the first column of Table II are the beginning of the ply drops. One
74 ply is dropped each 13 mm (0.5 in). More details on blade structural design can be found
75 in an authors' previous work².

76 For the experimental validation of the tools developed in this article, different versions of
77 bending tests were performed. These experiments were done on a steel frame as illustrated
78 in Figure 2. Two blades were instrumented with strain gauges on their exterior surface.
79 On the blade surface, reflective targets were also stuck for the blade to be scanned with
80 a 3D scanner. An EXAscan from Creaform was used. The comparison of deflected and

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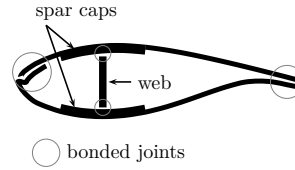


FIG. 1. Blade internal structure

TABLE I. Blade aerodynamic shape. z is the distance from the blade root, c is the airfoil chord length and θ_T is the airfoil twist angle

| z [m] | θ_T [deg.] | c [m] | Airfoil |
|------------|----------------------|------------|--------------|
| 0.000 | 15.3 | 0.200 | Circle |
| 0.160 | 15.3 | 0.200 | Circle |
| 0.560 | 15.3 | 0.334 | DU-97-W-300 |
| 0.760 | 14.5 | 0.318 | transition |
| 0.960 | 13.7 | 0.303 | transition |
| 1.160 | 12.9 | 0.287 | transition |
| 1.360 | 12.1 | 0.271 | DU-91-W2-250 |
| 1.560 | 11.2 | 0.256 | transition |
| 1.760 | 10.4 | 0.240 | transition |
| 1.960 | 9.6 | 0.224 | transition |
| 2.160 | 8.8 | 0.208 | transition |
| 2.360 | 8.0 | 0.193 | transition |
| 2.560 | 7.2 | 0.177 | DU-96-W-180 |
| 2.810 | 6.2 | 0.157 | transition |
| 3.060 | 5.2 | 0.138 | transition |
| 3.260 | 4.3 | 0.122 | transition |
| 3.460 | 3.2 | 0.106 | transition |
| 3.660 | 1.3 | 0.090 | transition |
| 3.760 | 0.0 | 0.083 | DU-96-W-180 |

TABLE II. Blade layup

| Blade region, distance from blade root | Laminate |
|---|--|
| Blade root circular region | |
| 0 mm–210 mm | [GC/CSM/(+45/-45) ₃ /0 ₆ /(+45/-45) ₃ /S/ (+45/-45) ₃ /0 ₈ /(+45/-45) ₃] |
| Spar cap, from 15 % to 45 % of chord length | |
| 210 mm–1074 mm | [GC/CSM/(+45/-45) ₂ /0 ₁₁ /(+45/-45) ₂] |
| 1074 mm–2624 mm | [GC/CSM/(+45/-45) ₂ /0 ₁₀ /(+45/-45) ₂] |
| 2624 mm–3173 mm | [GC/CSM/(+45/-45) ₂ /0 ₆ /(+45/-45) ₂] |
| 3173 mm–3776 mm | [GC/CSM/(+45/-45) ₂ /0 ₂ /(+45/-45) ₂] |
| Aerodynamic shells, outside spar cap | |
| 210 mm–960 mm | [GC/CSM/(+45/-45) ₂ /0 ₁₁ /(+45/-45) ₂] |
| 960 mm–3776 mm | [GC/CSM/(+45/-45) ₂ /0 ₁ /(+45/-45) ₂] |
| Shear web | |
| 210 mm–3776 mm | [(+45/-45) ₃ /C/(+45/-45) ₃] |
| GC: gelcoat, thickness: 0.51 mm. | |
| CSM: chopped strand mat, glass fiber with vinylester resin, thickness: 0.65 mm. | |
| S: steel studs or 0° unidirectional glass-epoxy filler. | |
| C: foam core, thickness : 19.05 mm. | |
| Thickness of +45° and –45° glass-epoxy layers: 0.23 mm. | |
| Thickness of 0° glass-epoxy layers: 0.50 mm. | |

81 undeflected scanned shapes allows the computation of the blade deflection. Note that for
82 all flapwise bending tests, the upper surface of the airfoils is oriented towards the floor and
83 the lower surface oriented upward.

84 The first test performed was a modal analysis of the blades. For that purpose, the blade
85 root was fixed on the test support. The blade tip was manually deflected and then, quickly
86 released. The signal of a strain gauge was recorded during the free vibration phase and a
87 Fourier transform of this signal was used to get the eigenfrequencies of the blade.

88 A second test was performed in order to recreated the critical load case of the blade.



FIG. 2. Design load test setup

89 This was done by loading the blade at two blade stations ($z = 2.360$ m and $z = 3.500$) as
90 shown in Figure 2. The loads were applied with manual winches and were transferred to the
91 blade by aluminum saddles. Loads cells were also installed on the winch cables to monitor
92 the applied loads.

93 A third test was intended to measure the blade cross-sectional stiffness properties. For
94 that purpose, the blade was loaded with a mass applied near its tip. Flapwise, edgewise and
95 torsional load cases were realized.

96 Finally, destructive tests were done. These tests were performed in two steps. In the first
97 step, the blade, fixed at the root, was also simply supported in its central part ($z = 2.360$ m,
98 the saddle closest to the root in Figure 2 was supported by a column) and loaded near its
99 tip ($z = 3.500$ m, saddle closest to the tip in Figure 2) to create a failure in its outer part. In
100 the second step of the destructive test, the blade was loaded at $z = 2.360$ m (saddle closest
101 to the root in Figure 2) to generate a failure in its inner part.

102 III. WIND TURBINE BLADE STRUCTURAL MODELS

103 The four different models for blade structural analysis are now presented. From the
104 simplest to the more complex, we have: (1) simple model, (2) classical strength of materials
105 model, (3) cross-sectional finite element model and (4) 3D shell finite element model. For
106 each model, the methods for the computation of eigenfrequencies, deflection, stresses, strains,
107 buckling, strength and cross-sectional properties are presented, where applicable. The first

108 two models are built based on literature. The third model was developed by the authors³.
109 The fourth model uses commercial finite element software.

110 Before the description of these four models, a section is dedicated to a review of some
111 generalities about beam cross-sectional properties.

112 A. Generalities about beam cross-sectional properties

113 For the purpose of this work, the coordinate system attached to the blade is defined as
114 follows. The z -axis is the blade longitudinal axis. When the blade pitch angle is 0, the
115 x -axis is in the rotor plane and point towards the direction of blade rotation, and the y -axis
116 is normal to the rotor plane and is pointing downwind. This coordinate system is fixed to
117 the blade and defines the flapwise and edgewise direction that correspond to the out-of-plane
118 and in-plane directions, respectively, only when the blade pitch angle is 0.

119 That being said, we can define the vector of beam internal loads as

$$\mathbf{V} = [V_x \ V_y \ N \ M_x \ M_y \ M_t]^T \quad (1)$$

120 where V_x and V_y are the shear forces, N is the axial force, M_x and M_y are the bending
121 moments and M_t is the torsional moment. Defining the beam reference axis displacements
122 as χ_x , χ_y and χ_z and the beam reference axis rotations as φ_x , φ_y and φ_z , we can express
123 the beam generalized strain vector as

$$\boldsymbol{\kappa} = [\gamma_{zx}^0 \ \gamma_{yz}^0 \ \epsilon_z^0 \ \kappa_x \ \kappa_y \ \kappa_z]^T \quad (2)$$

124 where

$$\begin{aligned} \gamma_{zx}^0 &= \frac{\partial \chi_x}{\partial z} - \varphi_y \\ \gamma_{yz}^0 &= \frac{\partial \chi_y}{\partial z} + \varphi_x \\ \epsilon_z^0 &= \frac{\partial \chi_z}{\partial z} \\ \kappa_x &= \frac{\partial \varphi_x}{\partial z} \\ \kappa_y &= \frac{\partial \varphi_y}{\partial z} \\ \kappa_z &= \frac{\partial \varphi_z}{\partial z} \end{aligned} \quad (3)$$

125 This describes the behavior of a Timoshenko beam.

126 The relation between beam internal loads \mathbf{V} and beam generalized strains $\boldsymbol{\kappa}$, for a given
127 cross section is

$$\mathbf{V} = \mathbf{K}_s \boldsymbol{\kappa} \quad (4)$$

128 or

$$\boldsymbol{\kappa} = \mathbf{F}_s \mathbf{V} \quad (5)$$

129 where \mathbf{K}_s is the cross-sectional stiffness matrix, which is symmetric and $\mathbf{F}_s = \mathbf{K}_s^{-1}$ is the
130 cross-sectional compliance matrix.

131 For an Euler beam, $\gamma_{zx}^0 = \gamma_{yz}^0 = 0$ in Eq. 3. The cross-sectional stiffness matrix then
132 reduces to a 4×4 matrix, the terms associated with transverse shear due to shear forces
133 being eliminated. The compliance matrix of an Euler beam model is equal to the compliance
134 matrix of a Timoshenko beam model from which the first two rows and columns are removed.

135 In the particular case where the origin of the cross-sectional coordinate system is coin-
136 cident with the elastic and shear centers, and the x and y -axes are the principal axes of
137 bending, Eq. 5 reduces to

$$\begin{bmatrix} V_x \\ V_y \\ N \\ M_x \\ M_y \\ M_t \end{bmatrix} = \begin{bmatrix} k_x \overline{GA}_x & 0 & 0 & 0 & 0 & 0 \\ 0 & k_y \overline{GA}_y & 0 & 0 & 0 & 0 \\ 0 & 0 & \overline{EA} & 0 & 0 & 0 \\ 0 & 0 & 0 & \overline{EI}_x & 0 & 0 \\ 0 & 0 & 0 & 0 & \overline{EI}_y & 0 \\ 0 & 0 & 0 & 0 & 0 & \overline{GJ} \end{bmatrix} \begin{bmatrix} \gamma_{zx}^0 \\ \gamma_{yz}^0 \\ \epsilon_z^0 \\ \kappa_x \\ \kappa_y \\ \kappa_z \end{bmatrix} \quad (6)$$

138 where

$$\begin{aligned} \overline{GA}_x &= \int G_{zx} \, dA \\ \overline{GA}_y &= \int G_{yz} \, dA \\ \overline{EA} &= \int E_z \, dA \\ \overline{EI}_x &= \int E_z y^2 \, dA \\ \overline{EI}_y &= \int E_z x^2 \, dA \end{aligned} \quad (7)$$

139 k_x and k_y are the correction factors for transverse shear⁴ and \overline{GJ} is obtained from the
140 solution of the torsion problem. Note that for an axisymmetric cross section:

$$\overline{GJ} = \int (x^2 G_{yz} + y^2 G_{zx}) \, dA \quad (8)$$

141 G_{zx} and G_{yz} are the material's shear moduli and E_x , E_y and E_z are the material's elastic
142 moduli. The integrations are performed over the beam cross section A .

143 We can show that the relationship between the beam generalized strains expressed in two
144 parallel coordinate systems (Figure 3) is

$$\boldsymbol{\kappa} = \mathbf{T}_s \boldsymbol{\kappa}' \quad (9)$$

145 where

$$\mathbf{T}_s = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & b \\ 0 & 1 & 0 & 0 & 0 & -a \\ 0 & 0 & 1 & -b & a & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix} \quad (10)$$

146 The relationship between two sets of internal loads expressed in these coordinate systems
147 is

$$\mathbf{V} = \mathbf{T}_s^{-T} \mathbf{V}' \quad (11)$$

148 where $\mathbf{T}_s^{-T} = (\mathbf{T}_s^{-1})^T = (\mathbf{T}_s^T)^{-1}$. The compliance matrix in the prime coordinate system is
149 thus

$$\mathbf{F}'_s = \mathbf{T}_s^{-1} \mathbf{F}_s \mathbf{T}_s^{-T} \quad (12)$$

150 so that

$$\boldsymbol{\kappa}' = \mathbf{F}'_s \mathbf{V}' \quad (13)$$

151 The cross-sectional stiffness matrix in the prime coordinate system is accordingly

$$\mathbf{K}'_s = \mathbf{T}_s^T \mathbf{K}_s \mathbf{T}_s \quad (14)$$

152 so that

$$\mathbf{V}' = \mathbf{K}'_s \boldsymbol{\kappa}' \quad (15)$$

153 Two sets of generalized strains and internal forces, after rotation of the coordinate system
154 (see Fig. 3), are related a follows:

$$\boldsymbol{\kappa} = \mathbf{T}_\theta \boldsymbol{\kappa}'' \quad (16)$$

155 and

$$\mathbf{V} = \mathbf{T}_\theta \mathbf{V}'' \quad (17)$$

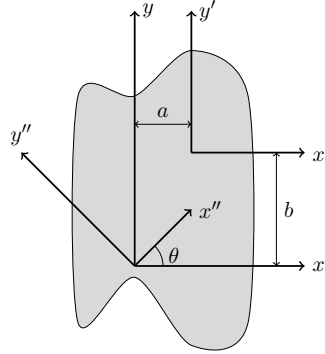


FIG. 3. Coordinate systems used to transform the cross-sectional stiffness and compliance matrices

156 where

$$\mathbf{T}_\theta = \begin{bmatrix} c & -s & 0 & 0 & 0 & 0 \\ s & c & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & c & -s & 0 \\ 0 & 0 & 0 & s & c & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix} \quad (18)$$

157 and where $s = \sin \theta$ and $c = \cos \theta$.

158 The cross-sectional compliance matrix in the double prime coordinate system is then:

$$\mathbf{F}_s'' = \mathbf{T}_\theta^{-1} \mathbf{F}_s \mathbf{T}_\theta \quad (19)$$

159 so that

$$\boldsymbol{\kappa}'' = \mathbf{F}_s'' \mathbf{V}'' \quad (20)$$

160 and the cross-sectional stiffness matrix in the double prime coordinate system is

$$\mathbf{K}_s'' = \mathbf{T}_\theta^{-1} \mathbf{K}_s \mathbf{T}_\theta \quad (21)$$

161 so that

$$\mathbf{V}'' = \mathbf{K}_s'' \boldsymbol{\kappa}'' \quad (22)$$

162 The shear center of the cross section is defined as the point where an applied shear force
163 does not cause torsion. So in order for a and b to be the coordinates of this point, the

164 shear-twist coupling terms of the compliance matrix (F'_{s16} and F'_{s26} of Eq. 12) must be null.
165 The shear center coordinates are then:

$$\begin{aligned} x_c &= -\frac{F_{s26}}{F_{s66}} \\ y_c &= \frac{F_{s16}}{F_{s66}} \end{aligned} \quad (23)$$

166 The elastic center of the cross section is defined as the point where an applied axial
167 force does not cause bending. In order for a and b to be the coordinates of this point, the
168 extension-bending coupling terms (F'_{s34} and F'_{s35} of Eq. 12) of the cross-sectional compliance
169 matrix must be zero. The elastic center coordinates are then:

$$\begin{aligned} x_e &= \frac{F_{s44}F_{s35} - F_{s45}F_{s34}}{F_{s44}F_{s55} - F_{s45}^2} \\ y_e &= \frac{F_{s45}F_{s35} - F_{s55}F_{s34}}{F_{s44}F_{s55} - F_{s45}^2} \end{aligned} \quad (24)$$

170 When the origin of the coordinate system is also the elastic center, this coordinate system
171 can be rotated by an angle $\theta = \theta_1$ to get the principal axes of bending which are characterized
172 by the absence of coupling between both directions, i.e., $F''_{s45} = 0$ in Eq. 19. The orientation
173 of the principal axes of bending is then:

$$\theta_1 = \frac{1}{2} \arctan \left(\frac{-2F_{s45}}{F_{s55} - F_{s44}} \right) \quad (25)$$

174 B. Model 1: Simple model

175 Several simple models for blade preliminary analysis have been proposed⁵⁻⁷ (for instance).
176 The one proposed here is based on Hansen's book⁸, where the cross section is represented as
177 shown in Figure 4. The only parts modelled are the spar caps idealized as two rectangular
178 strips. c is the chord length and t is the airfoil thickness.

179 The cross-sectional inertia (about an axis parallel to the chord line and passing through
180 the mid distance between both spar caps) is

$$I = \frac{bt^3}{12} - \frac{b(t-2h)^3}{12} \quad (26)$$

181 and the maximum strain (and absolute value of minimum strain) is

$$\epsilon = \frac{Mt}{2EI} \quad (27)$$

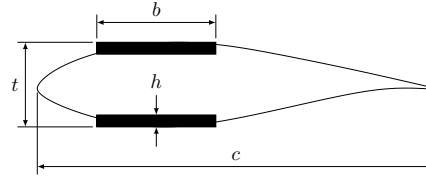


FIG. 4. Model 1: Simplified blade model

182 where M is the bending moment. The term "local flapwise" is used to describe this bending
 183 direction. As a simple and conservative assumption, the resultant of the in-plane and out-
 184 of-plane bending moments is supposed to be applied about an axis parallel to the chord line.
 185 E is the spar cap elastic modulus.

186 Using Eq. 26 and 27, the required spar cap thickness as a function of airfoil shape, material
 187 properties and bending moment is

$$h = \frac{t}{2} - \frac{1}{2} \sqrt[3]{t^3 - \frac{6Mt}{Eb}} \quad (28)$$

188 where ϵ is the failure strain.

189 This model can then be used for the dimensioning of the blade spar caps with Eq. 28.
 190 Once the spar cap thickness is known, at several cross sections along the blade length, a
 191 blade mass can be estimated using typical relations between spar cap mass and whole blade
 192 mass. The model can also be used for the stress and strain analysis with equations 27 and
 193 26.

194 As just demonstrated, this model is well suited for preliminary structural dimensioning
 195 of wind turbine blades based on strength. However, it is not usable to evaluate natural
 196 frequencies, deflection and buckling. Also, the only cross-sectional stiffness property it
 197 computes is the local flapwise bending stiffness.

198 C. Model 2: Classical strength of materials model

199 This model is based on classical strength of materials theory. Figure 5 shows the different
 200 coordinate systems used for this analysis. Taking into account the thin-walled nature of the
 201 wind turbine blade, the integral over the area A of the cross section is computed as a line

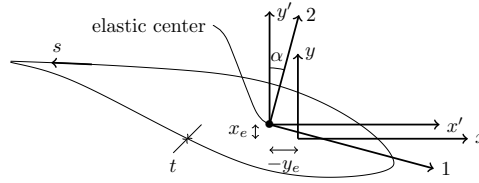


FIG. 5. Model 2: Coordinate systems for the classical strength of materials model

202 integral over the walls:

$$\int_A f \, dA = \int_s f t \, ds \quad (29)$$

203 where t is the wall thickness, s is the variable along the walls and f is the quantity to
204 integrate.

205 This integral is then evaluated numerically by discretizing the contour in several segments
206 and summing over these segments:

$$\int_s f t \, ds \approx \sum_i f_i t_i \Delta s_i \quad (30)$$

207 Before performing these calculations, the contour points are translated towards the inter-
208 rior of the cross section by a distance of half of the wall thickness. The coordinates of the
209 contour point then represent the mid thickness surface of the walls.

210 For the analysis of the extension and bending behaviour of the blade, a first set of cross-
211 sectional properties relative to the reference coordinate system (xy) can be computed as⁸:

$$\begin{aligned} \overline{EA} &= \int_A E \, dA \\ \overline{ES}_x &= \int_A y E \, dA \\ \overline{ES}_y &= \int_A x E \, dA \\ \overline{EI}_x &= \int_A y^2 E \, dA \\ \overline{EI}_y &= \int_A x^2 E \, dA \\ \overline{EI}_{xy} &= \int_A xy E \, dA \end{aligned} \quad (31)$$

212 where E is the material's elastic modulus in the z -direction. For laminates, the effective

213 elastic modulus is used. From these properties, the coordinates of the elastic center are:

$$\begin{aligned} x_e &= \frac{\overline{ES}_y}{EA} \\ y_e &= \frac{\overline{ES}_x}{EA} \end{aligned} \quad (32)$$

214 Knowing the location of the elastic center, the bending stiffness relative to this point are:

$$\begin{aligned} \overline{EI}_{x'} &= \int_A (y')^2 E dA = \overline{EI}_x - y_e^2 \overline{EA} \\ \overline{EI}_{y'} &= \int_A (x')^2 E dA = \overline{EI}_y - x_e^2 \overline{EA} \\ \overline{EI}_{x'y'} &= \int_A x' y' E dA = \overline{EI}_{xy} - x_e y_e \overline{EA} \end{aligned} \quad (33)$$

215 Finally, the angle between the x' -axis and the closest principal axes of bending is:

$$\alpha = \frac{1}{2} \arctan \left(\frac{2\overline{EI}_{x'y'}}{\overline{EI}_{x'} - \overline{EI}_{y'}} \right) \quad (34)$$

216 and the bending stiffnesses around these axes (referred as local axes) are

$$\begin{aligned} \overline{EI}_1 &= \frac{\overline{EI}_x + \overline{EI}_y}{2} + \frac{\overline{EI}_x - \overline{EI}_y}{2} \cos 2\alpha + \overline{EI}_{xy} \sin 2\alpha \\ \overline{EI}_2 &= \frac{\overline{EI}_x + \overline{EI}_y}{2} - \frac{\overline{EI}_x - \overline{EI}_y}{2} \cos 2\alpha - \overline{EI}_{xy} \sin 2\alpha \end{aligned} \quad (35)$$

217 \overline{EI}_1 is the local flapwise bending stiffness and \overline{EI}_2 is the local edgewise bending stiffness.

218 Generally, for a wind turbine blade, $\overline{EI}_2 > \overline{EI}_1$.

219 For a given cross section, if the local flapwise bending moment M_1 , the local edgewise
220 bending moment M_2 and the axial force N are known, it is possible to compute the axial
221 strain of a point using:

$$\epsilon_z = \frac{M_1 y_p}{\overline{EI}_1} - \frac{M_2 x_p}{\overline{EI}_2} + \frac{N}{EA} \quad (36)$$

222 where x_p and y_p are the coordinates along the 1- and 2-axes respectively and M_1 and M_2
223 are computed as:

$$\begin{aligned} M_1 &= M_x \cos \alpha - M_y \sin \alpha \\ M_2 &= M_x \sin \alpha + M_y \cos \alpha \end{aligned} \quad (37)$$

224 For the analysis of torsion, from strength of materials textbooks^{9,10}, the unit torsion angle
225 of one of the cells of a multicell thin-walled beam is:

$$\kappa_z = \frac{1}{2A_i} \left(q_i \oint_i \frac{ds}{Gt} - q_{i-1} \int_{i-1,i} \frac{ds}{Gt} - q_{i+1} \int_{i,i+1} \frac{ds}{Gt} \right) \quad (38)$$

226 t and G are respectively the wall thickness and shear modulus. For laminates, the effective
227 shear modulus is used. A_i is the surface enclosed by the i th cell and q_i , q_{i-1} and q_{i+1} are
228 respectively the shear flow in the walls of the i th cell, the shear flow in the web left to the
229 i th cell and the shear flow in the web right to the i th cell (see Figure 6). Integrals over (i) ,
230 $(i-1, i)$ and $(i, i+1)$ are respectively integral over the i th cell, the web to the left to the
231 i th cell and the web to the right to the i th cell. Using the following notation,

$$\begin{aligned}\delta_i &= \oint_i \frac{ds}{Gt} \\ \delta_{i,j} &= \int_{i,j} \frac{ds}{Gt}\end{aligned}\quad (39)$$

232 Eq. 38 can be written as

$$2A_i\kappa_z = q_i\delta_i - q_{i-1}\delta_{i-1,i} - q_{i+1}\delta_{i,i+1}\quad (40)$$

233 Knowing that each cell must have the same unit torsion angle κ_z , the previous equation is
234 repeated for each of the n cells of the beam. This results in a system of n equations to
235 compute the shear flow q_i in each cell:

$$2\kappa_z\mathbf{A} = \boldsymbol{\delta}\mathbf{q}\quad (41)$$

236 where

$$\begin{aligned}\mathbf{A} &= [A_1 \ A_2 \ A_3 \ \dots \ A_n]^T \\ \mathbf{q} &= [q_1 \ q_2 \ q_3 \ \dots \ q_n]^T \\ \boldsymbol{\delta} &= \begin{bmatrix} \delta_1 & -\delta_{1,2} & 0 & 0 & \dots & 0 & 0 \\ -\delta_{1,2} & \delta_2 & -\delta_{2,3} & 0 & \dots & 0 & 0 \\ 0 & -\delta_{2,3} & \delta_3 & -\delta_{3,4} & \dots & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -\delta_{n-1,n} & \delta_n \end{bmatrix}\end{aligned}\quad (42)$$

237 From equilibrium,

$$M_t = 2\mathbf{A}^T\mathbf{q}\quad (43)$$

238 where M_t is the torsion moment. Introducing the torsional stiffness \overline{GJ} :

$$M_t = \overline{GJ}\kappa_z\quad (44)$$

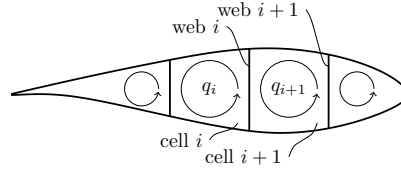


FIG. 6. Model 2: Torsion analysis of a multicell thin-walled beam

239 From Eq. 41, 43 and 44, the torsional rigidity can be computed as:

$$\overline{GJ} = 4\mathbf{A}^T \boldsymbol{\delta}^{-1} \mathbf{A} \quad (45)$$

240 The deflection analysis procedure is as follows⁸. Assuming that the applied forces (p_x ,
241 p_y and p_z) and moments (m_x , m_y and m_z) per unit length are known at n different points
242 from the blade root to the blade tip and assuming a linear variation between these points,
243 the internal load distribution at each of these points (V_x , V_y , N , M_x , M_y and M_t) can be
244 computed numerically.

245 For each point, the beam curvature can be computed by first transferring the bending
246 moments in the principal axis of bending and then, by evaluating the curvatures around
247 these axes. Finally, these curvatures are converted into the xy -coordinate system.

248 The axial and torsional beam generalized deformations can be evaluated with:

$$\begin{aligned} \epsilon_x^{0,i} &= \frac{N^i}{EA^i} \\ \kappa_z^i &= \frac{M_t^i}{GJ^i} \end{aligned} \quad (46)$$

249 Assuming a linear variation of these beam deformations between loading points, the beam
250 rotations and deflections can be computed numerically by solving Eq. 3 assuming that
251 $\gamma_{zx}^0 = \gamma_{yz}^0 = 0$. By doing so, the effect of transverse shear is neglected, following the Euler-
252 Bernoulli hypothesis.

253 For the modal analysis, the blade mass distribution is required. Following the same
254 procedure as for the calculation of the cross-sectional stiffness properties, the mass per unit
255 length m' can be computed at some location along the blade length as:

$$m' = \int_A \rho \, dA \quad (47)$$

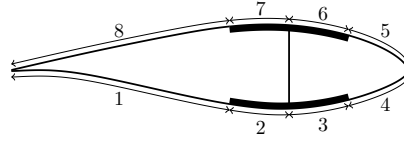


FIG. 7. Model 2: Example of panels for buckling analysis

256 where ρ is the material density. The modal analysis procedure used here⁸ allows the com-
257 putation of the first flapwise and edgewise bending modes, the most important for blade
258 design.

259 For the buckling analysis, at a given cross section, the blade exterior surface is separated
260 in different panels. A separation between two adjacent panels occurs each time (1) a web is
261 connected to the blade surface and (2) there is a change in the laminate. For instance, in
262 Figure 7, the cross section is separated in 8 panels. For each panel, the width b is computed
263 as the sum of the length of the elements forming the panel and the critical compressive force
264 per unit length is calculated using the conservative infinite length (in the blade longitudinal
265 direction) buckling solution for flat panel simply supported on all sides¹¹ :

$$N_{cr} = \frac{2\pi^2}{b^2} \left(\sqrt{D_{11}D_{22} + D_{12} + 2D_{66}} \right) \quad (48)$$

266 where D_{11} , D_{22} , D_{12} and D_{66} are the panel bending stiffnesses obtained from the classical
267 lamination theory¹¹⁻¹⁴.

268 For each panel, the mean (along its width) compressive force per unit length N_z is also
269 computed. For each element, N_z is computed as:

$$N_z = \sigma_z t = E_z \epsilon_z t \quad (49)$$

270 where t is the panel thickness and E_z is its modulus of elasticity (effective modulus for
271 laminates).

272 A buckling failure index can then be computed as:

$$F_{\text{buckling}} = \frac{N_z}{N_{cr}} \quad (50)$$

273 D. Model 3: Cross-sectional finite element model

274 The third model for blade structural analysis is based on a finite element discretization
275 of the cross section. This framework, sometimes called Nonhomogeneous Anisotropic Beam

FIG. 8. Model 3: Example of blade cross-sectional mesh for analysis

276 Section Analysis (NABSA) was proposed by Giavotto et al.¹⁵ and used by Blasques et
277 al.¹⁶⁻¹⁹. It uses triangular and quadrilateral elements to discretize the beam cross section.
278 Capitalizing on the fact that wind turbine blades are thin-walled structures, the authors have
279 added line finite elements to this framework, allowing modelling thin-walled cross sections³.
280 Figure 8 shows an example of blade cross-sectional mesh for this method.

281 This method allows the computation of a 6×6 cross-sectional stiffness matrix at different
282 points along the blade length. The post-processing of the results also allows the computation
283 of the stresses and strains in the cross sections.

284 When it is needed to model the whole blade behavior (for deflection and modal analysis),
285 the blade is discretized using 3-node Timoshenko beam finite elements.

286 E. Model 4: 3D shell finite element model

287 The use of 3D finite element models for structural analysis of wind turbine blades is
288 common. Most of the time, due to the thin-walled topology of these structures, shell fi-
289 nite elements are employed. As the exterior shape of the blade is the reference shape for
290 aerodynamics, the finite element mesh is frequently located on this surface and the element
291 thickness is built towards the blade interior.

292 The particularity of this model over the three previous is that it is able to manage the
293 effects of the variation of the cross-sectional shape of the beam along its length. However,
294 it takes more work to set up the model.

295 The use of this kind of model is well documented for the computation of natural frequen-
296 cies, deflection, stresses, strains and buckling²⁰⁻²⁴ (for instance). However, a difficulty arises
297 for computing the blade cross-sectional properties needed for the aeroelastic analysis and it
298 is to this task that the remaining of this section is dedicated.

299 To compute the cross-sectional properties from a 3D finite element model, we need to
300 know, at different locations along blade length, the internal loads and the beam generalized
301 deformations. Knowing the internal loads is the easiest task, but extracting the generalized
302 deformation is not straightforward. Usually, to do so, the blade reference axis displacements

303 and rotations are computed from the nodal displacements and rotations of the nodes forming
304 the cross section. The longitudinal distribution of these reference axis displacements is then
305 derived using Eq. 3 to compute the generalized deformations.

306 The method proposed here is highly inspired by Malcolm et al.¹. The particularity of the
307 proposed method is that the computation of cross-sectional displacements and rotations are
308 performed using a formulation similar to an interpolation element like NASTRAN's RBE3
309 element.

310 The procedure to compute the distribution of cross-sectional properties along the blade
311 length (illustrated in Fig. 9) is as follows:

- 312 1. The blade is subjected to a set of linearly independent loads cases. N load cases are
313 needed for a $N \times N$ cross-sectional stiffness matrix.
- 314 2. n stations where the cross-sectional properties are to be computed are determined.
- 315 3. For each of these stations, the displacements (translation and rotation) of the reference
316 point is computed using the kinematic relation linking the reference node to the con-
317 nected nodes of a RBE3 element. The reference node of the element is the point that
318 is located at the intersection of the beam's reference axis and cross-sectional plane.
319 The connected nodes are all the nodes within a given distance on both sides of the
320 cross-sectional plane. This distance should be as small as possible (but large enough
321 to get nodes to cover the entire cross section) to get a good representation of the
322 transverse shear strains.
- 323 4. Once the displacements of the beam's reference axis are known at each stations along
324 its length, Eq. 3 is used to compute the generalized strains at each station. The
325 derivatives are computed using second order numerical derivations on 3 points.
- 326 5. For each station, the cross-sectional stiffness matrix is computed by solving the fol-

FIG. 9. Model 4: Procedure for the cross-sectional properties evaluation from a 3D finite element model from a set of unit loads at blade tip

327 following system:

$$\begin{bmatrix} \begin{bmatrix} V_x \\ V_y \\ N \\ M_x \\ M_y \\ M_z \end{bmatrix}_1 & \dots & \begin{bmatrix} V_x \\ V_y \\ N \\ M_x \\ M_y \\ M_z \end{bmatrix}_6 \end{bmatrix} = \mathbf{K}_s \begin{bmatrix} \begin{bmatrix} \gamma_{zx}^0 \\ \gamma_{zy}^0 \\ \epsilon_z^0 \\ \kappa_x \\ \kappa_y \\ \kappa_z \end{bmatrix}_1 & \dots & \begin{bmatrix} \gamma_{zx}^0 \\ \gamma_{zy}^0 \\ \epsilon_z^0 \\ \kappa_x \\ \kappa_y \\ \kappa_z \end{bmatrix}_6 \end{bmatrix} \quad (51)$$

328 where the index of a vector designates the load case to which it is associated.

329 F. Application to the experimental cross-sectional analysis

330 The method used to compute the blade cross-sectional properties of model 4 can also be
 331 used for the evaluation of the cross-sectional properties of a blade during a bending test.
 332 As described in Section II, this test has been performed on the WESNet blade. Instead
 333 of using the displacements of connected nodes to compute the cross-sectional generalized
 334 deformations, the displacements of the 3D scanner's reflective target were used.

335 Only 3 modes of deformation were considered in this analysis: flapwise bending, edgewise
 336 bending and torsion. Then, the following system has to be solved to get a 3×3 stiffness
 337 matrix:

$$\begin{bmatrix} \begin{bmatrix} M_x \\ M_y \\ M_z \end{bmatrix}_1 & \dots & \begin{bmatrix} M_x \\ M_y \\ M_z \end{bmatrix}_3 \end{bmatrix} = \mathbf{K}_s \begin{bmatrix} \begin{bmatrix} \kappa_x \\ \kappa_y \\ \kappa_z \end{bmatrix}_1 & \dots & \begin{bmatrix} \kappa_x \\ \kappa_y \\ \kappa_z \end{bmatrix}_3 \end{bmatrix} \quad (52)$$

338 The 3 load cases used was a flapwise force near the blade tip, an edgewise force near the
 339 blade tip and an excentric force close to the blade tip (coupled torsion and flapwise bending).

340 G. Summary of model capabilities

341 To conclude this section, Table III presents a summary of the capabilities of the 4 struc-
 342 tural models described above.

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TABLE III. Summary of model capabilities

| Analysis types | Model 1: beam model; cross section idealized as two rectangular strips | Model 2: beam model; strength of ma- terials theory for thin-walled composite beams | Model 3: beam model; thin-walled cross-sectional finite elements | Model 4: shell finite ele- ment model of the whole blade |
|-------------------------------|---|---|--|---|
| Cross-sectional properties | Local flapwise bending | Extension, two- plane bending, and their cou- plings; torsion | Full 6×6 cross-sectional stiffness matrix | Full 6×6 cross-sectional stiffness matrix |
| Natural frequencies | No | Edgewise and flapwise modes for Euler beam | Eigenvalue so- lution for Tim- oshenko beam finite elements | Eigenvalue so- lution for 3D shell finite ele- ments |
| Deflection | No | Edgewise and flapwise deflec- tion for Euler beam | Beam axis displacements and rotations for Timoshenko beam finite elements | Full field dis- placements and rotations for 3D shell finite ele- ments |
| Strains | Longitudinal strain | Longitudinal strain due to extension and bending, shear strain due to torsion | Full 3D strains | Full 3D strains |
| Buckling | No | Linear buckling of long plates | Not imple- mented | Linear buckling |
| Strength | Based on longi- tudinal strain | Based on longi- tudinal strain | Based on fiber direction stress | Based on fiber direction stress |

343 IV. COMPARISON OF DIFFERENT WIND TURBINE BLADE MODELS 344 AGAINST EXPERIMENTAL RESULTS

345 The four different structural models presented in the previous section are now compared
346 with each other and with the WESNet blade experimental results. First, the characteristics
347 of each model as well as the resulting cross-sectional properties are presented. Then, the
348 results with respect to modal analysis, deflection, strains, buckling and composite strength
349 are discussed.

350 A. Models description

351 The simple model (model 1) of the WESNet blade is built by evaluating the local flapwise
352 moment of inertia I (using Eq. 26) and local flapwise bending stiffness EI . These charac-
353 teristics are given at each blade longitudinal station of Table I as well as at other points of
354 interest like ply drops.

355 The classical strength of materials model (model 2) is used to compute the WESNet
356 blade cross-sectional properties at the same longitudinal stations as the model 1. Each of
357 these longitudinal station is discretized using 100 (for circular sections) to 210 (for airfoil
358 sections) segments.

359 The WESNet blade is also modelled using the cross-sectional finite element method
360 (model 3). Each section is discretized using 100 to 117 quadratic elements (depending
361 on whether the shear web is present or not). The aerodynamic surface elements use the
362 offset node option, i.e., the nodes are on the exterior surface of the blade and the element
363 thickness is built towards the blade interior. The shear web elements use the conventional
364 mid thickness surface definition. The blade cross-sectional properties are computed at the
365 same longitudinal stations as for models 1 and 2.

366 The 3D shell finite element model (model 4) of the blade is built using the Altair Hyper-
367 Works suite. Hypermesh, OptiStruct and Hyperview are used respectively as pre-processor,
368 solver and post-processor. The OptiStruct solver uses the same input format as NASTRAN.
369 The model uses 4-node linear shell elements. As seen in Figure 10, the blade is discretized
370 with 38 elements along its chords length and the element size reduces towards the blade
371 tip to keep their aspect as square as possible. The model uses a total of 46 763 nodes and

FIG. 10. Mesh of the WESNet 3D shell finite element model

372 47 272 elements (all but 3 are quadrilateral). Laminates are defined using the PCOMPP
 373 method of OptiStruct. As for the cross-sectional finite element model, nodes are on the
 374 blade's exterior surface and the element thickness is built towards the blade interior. The
 375 shear web elements use the conventional mid thickness surface definition. With this 3D
 376 shell finite element model, the methodology presented earlier is used to compute the cross-
 377 sectional properties of the WESNet blade. 75 equally spaced computation points are used
 378 along the blade length.

379 Using the methodology presented for model 4, the experimental WESNet blade cross-
 380 sectional properties are also computed. As the test jig does not allow the application of
 381 axial load cases and as the precision of the measurements do not allow the computation of
 382 transverse shear deformation, a 3×3 cross-sectional stiffness matrix is obtained (torsion and
 383 two axes of bending). The cross-sectional characterization test was performed on one blade
 384 only.

385 B. Cross-sectional properties

386 The cross-sectional stiffness properties computed from the different models of the WES-
 387 Net blade are now compared. Due to the blade configuration, for which there are no signif-
 388 icant couplings at the laminate level, the blade cross-sectional stiffness matrix should take
 389 the following form:

$$\mathbf{K}_s = \begin{bmatrix} K_{s11} & K_{s12} & 0 & 0 & 0 & K_{s16} \\ K_{s12} & K_{s22} & 0 & 0 & 0 & K_{s26} \\ 0 & 0 & K_{s33} & K_{s34} & K_{s35} & 0 \\ 0 & 0 & K_{s34} & K_{s44} & K_{s45} & 0 \\ 0 & 0 & K_{s35} & K_{s45} & K_{s55} & 0 \\ K_{s16} & K_{s26} & 0 & 0 & 0 & K_{s66} \end{bmatrix} \quad (53)$$

390 There is coupling between extension and bending because the reference axes are not neces-
 391 sarily centered at the elastic center (K_{s34} and K_{s35} terms) nor aligned with the principal axes
 392 of bending (K_{s45} term). There is also coupling between both transverse shear deformation
 393 (K_{s12} term). Finally, there is a coupling between the transverse shears and the torsion de-

394 formations because the reference axes are not necessarily centered at the shear center (K_{s16}
395 and K_{s26} terms).

396 Models 3 and 4 are able to evaluate all the terms of the 6×6 cross-sectional stiffness
397 matrix. Model 2 only evaluates those associated with the extension, bending and torsion
398 deformations (K_{s33} , K_{s44} , K_{s55} , K_{s34} , K_{s35} , K_{s45} and K_{s66}). Model 1 only evaluate the local
399 flapwise bending stiffness (K_{s44} , when the x are is the flapwise principal axis of bending).
400 Finally, the experimentations allow computing terms associated with the bending and torsion
401 behaviour only (K_{s44} , K_{s55} , K_{s45} and K_{s66}).

402 The terms of the cross-section stiffness matrix that should be null are effectively null
403 when computed by model 3. Model 4 returns values that are not null but are small when
404 compared to the other terms.

405 The detailed results for each of the non-null terms are then presented. Figure 11a, b and c
406 present the results obtained for the terms associated with the transverse shear deformation.
407 We can see that both models show the same trends but a significant error is observable.
408 This is caused by the imprecision of model 4 to evaluate transverse shear properties. This
409 imprecision is due to the fact that transverse shear deformation is small in this blade and
410 hard to capture with the used method so a small error results in a larger relative difference.

411 As shown in Figure 11d, the axial stiffness from models 2, 3 and 4 are similar in the
412 outboard region of the blade. In the inboard region, model 4 shows the same trends but
413 with important differences. This is due to the fact that this model (3D shell finite element)
414 is able to take into account the effects of the rapidly changing cross section shape in this
415 part of the blade. On their side, models 2 and 3 suppose a constant cross-sectional shape.
416 The same conclusion can be made when looking at the bending stiffness of Figure 11e and
417 f.

418 As the tests performed do not include axial loads on the blade, the extension-bending
419 couplings cannot be evaluated. So, the results of models 2, 3 and 4 have to be transferred to
420 the elastic center to be compared with the experimental results. This is shown in Figure 11g
421 and h where we can see a difference between the models of up to 30 % for the flapwise
422 bending stiffness (K_{s44}) and of up to 20 % for the edgewise bending stiffness (K_{s55}). Note
423 that the experimental results were not available for approximately the first 1 m closest to
424 the blade root.

425 Starting from the cross-sectional stiffness matrices computed at the elastic center, it is

426 possible to compute the orientation of the principal axes of bending. The results of the
427 angle θ_1 between the x -axis and the flapwise principal axis of bending (or local flapwise
428 axis) are shown in Figure 11*i*. All models return values close to each other. Transferring the
429 K_{s44} term to this axis allows to compare with the local flapwise bending stiffness obtained
430 from model 1. As shown in Figure 11*j*, model 1 return values that are close to those of the
431 other models in the outboard region of the blade. In fact, model 1 computes values that
432 are approximately 2 to 3 % smaller than those of the other models. A higher difference is
433 observed from blade root to $z = 1$ m. This is because, in that region, the aerodynamic shells
434 outside the spar caps are as thick as the spar cap, contributing to the blade stiffness, which
435 is not taken into account in model 1. However, we can say that model 1 gives conservative
436 results that are close to the values obtained from the other models.

437 The results of torsional stiffness are presented in Figure 11*k*, where they are transferred
438 to the shear center for models 3 and 4 in order to be able to compare with model 2 and
439 the experimental results. The results from models 2, 3 and 4 show similar trends. Around
440 $z = 0.5$ m, Model 4 differs from the two others. This is due to the fact that this model takes
441 into account the changing cross-sectional shape. When zooming in the outboard section
442 of the blade (see Figure 11*l*, we can observe differences between the different models of
443 up to 50 %. At the opposite of what was observed for bending, model 4 underestimates
444 the torsional stiffness. This can be explained by the difficulty of obtaining good results for
445 torsion from a shell finite element model using offset shells as reported in the literature²⁵⁻²⁷.

446 As shown in Figures 11*l* and 11*m*, the shear center position as computed by models 3 and
447 4 are quite different. However, these differences are still small relative to the cross-sectional
448 dimensions. It illustrates the difficulties associated with the computation of the transverse
449 shear properties using the 3D shell finite element model.

450 Results for the K_{s16} , K_{s26} , K_{s34} , K_{s35} and K_{s45} terms are not presented here as they are
451 used to compute the elastic and shear center as well as the orientation of the principal axes
452 of bending.

453 Comparing the blade cross-sectional properties from the different numerical models
454 against the experimental data leads to the following observations. Model 1 gives a conser-
455 vative but fair estimation of the local flapwise bending properties at low calculating cost.
456 Model 2 gives very good results for extension, bending and torsion. Model 3 seems to be the
457 most reliable according to the validation performed on it³. Model 4 gives very good results

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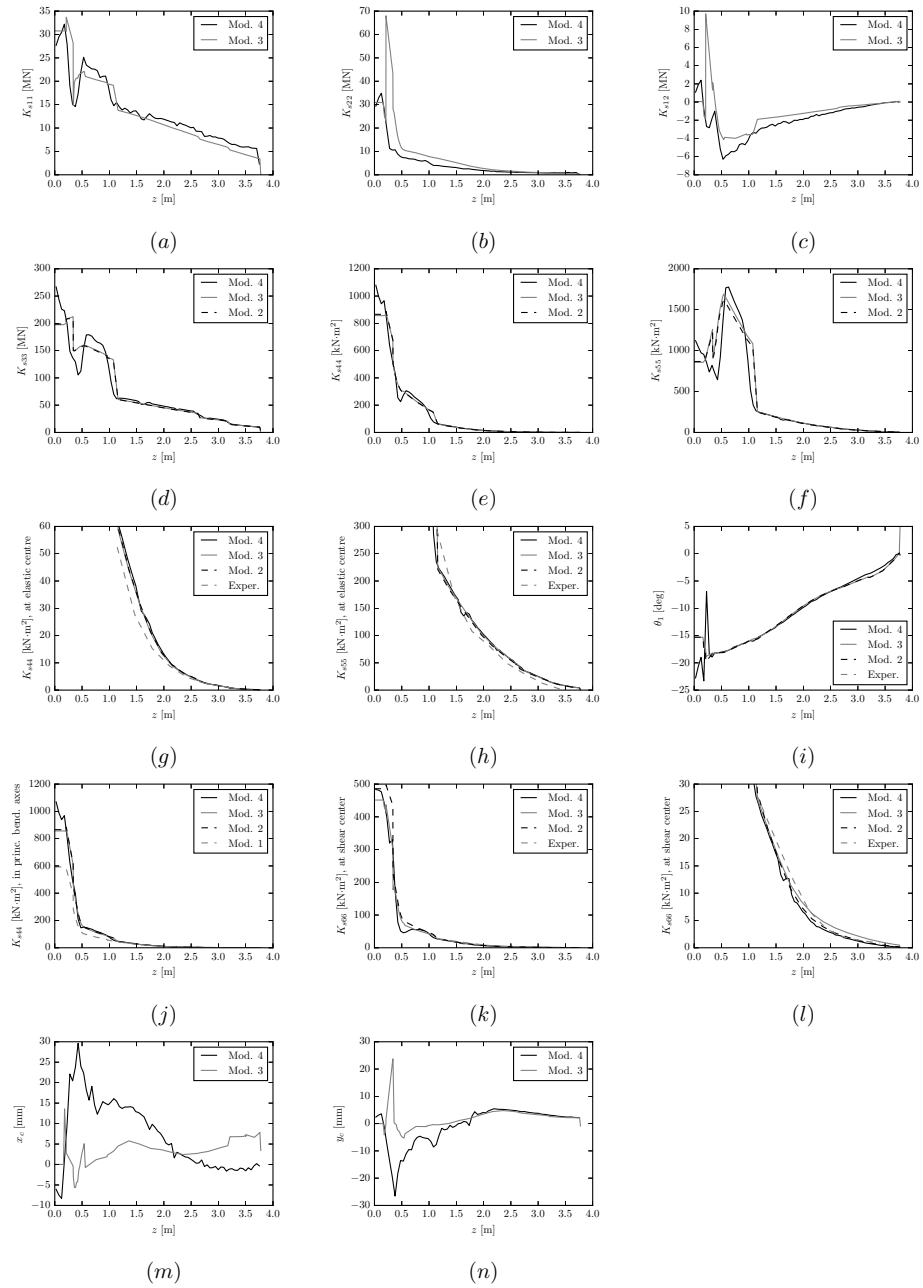


FIG. 11. Cross-sectional stiffness terms and associated properties. Experimental results come from blade 1. (a) K_{s11} . (b) K_{s22} . (c) K_{s12} . (d) K_{s33} . (e) K_{s44} . (f) K_{s55} . (g) K_{s44} at elastic centre. (h) K_{s55} at elastic center. (i) θ_1 . (j) K_{s44} in principal bending axes. (k) K_{s66} at shear center. (l) K_{s66} at shear center (zoom). (m) x_c . (n) y_c

458 for axial and bending behaviours. For a wind turbine blade with no bend-twist coupling,
459 models 2, 3 and 4 are correct. Model 1 is usable for preliminary analysis based on flapwise
460 bending behaviour.

461 C. Natural frequencies

462 Blade natural frequencies computed using models 2, 3, and 4 are now compared with the
463 experimental data. The experimental method consists in deflecting the blade tip and then,
464 releasing it suddenly. The Fourier transform of the time-domain signal of one of the strain
465 gauges is computed to obtain the natural frequency. Two natural frequencies can clearly be
466 identified at 6.4 Hz and 15.0 Hz. Both blades tested give the same results.

467 Table IV shows the comparison of model data with experimental data. In each model,
468 the materials' density was adjusted to get a blade mass of 21.6 kg, equivalent to the real
469 blade. This mass does not include the 6.40 kg of steel parts at blade root. Each of these
470 models predict a similar center of gravity location, but these values are approximately 3 %
471 lower than the one of the real blade. For the first blade natural frequency, all three models
472 are within 3.5 % difference relative to the experimental value. This difference increases up
473 to 14.6 % for the second natural frequency. The differences between models 2, 3 and 4 are
474 within a 2 to 3 % range and they are satisfactorily predicting the experimental data.

475 As we can see in Figure 12 for mode 3, each model predicts mode shapes that are similar
476 to each other. Mode shapes computed by model 4 are also similar.

477 D. Deflection

478 Looking now at the blade deflections during the design load test (defined in Section II),
479 we can see in Figure 13 that all models predict well the flapwise deflection. At tip, all values
480 are within a 5 % interval. The experimental data were obtained from tests on two different
481 blades (results for both two blades are shown in Figure 13). For the edgewise deflection, we
482 can observe some scatter in the experimental data, which is normal due to the low deflection
483 values and to the precision of the method used to compute them (see Section II).

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TABLE IV. Results of the blade modal analysis. m is the blade mass, Z_{cg} is the distance from its root to its center of gravity along its length and f_i is the i th blade natural frequency. δ_j is the relative difference (in %) between model j and experimental data. Experimental results are the same for both blades.

| | | Experimental | Model 2 | Model 3 | Model 4 | δ_2 | δ_3 | δ_4 |
|----------|------|--------------|---------|---------|---------|------------|------------|------------|
| m | [kg] | 21.60 | 21.60 | 21.60 | 21.60 | 0.0 | 0.00 | 0.00 |
| Z_{cg} | [m] | 1.180 | 1.143 | 1.143 | 1.140 | -3.1 | -3.1 | -3.4 |
| f_1 | [Hz] | 6.400 | 6.616 | 6.474 | 6.559 | 3.4 | 1.2 | 2.5 |
| f_2 | [Hz] | 15.00 | 17.19 | 17.05 | 16.72 | 14.6 | 13.7 | 11.5 |
| f_3 | [Hz] | - | 20.86 | 20.32 | 20.21 | - | - | - |
| f_4 | [Hz] | - | 42.69 | 40.42 | 40.55 | - | - | - |
| f_5 | [Hz] | - | 58.32 | 57.08 | 56.18 | - | - | - |

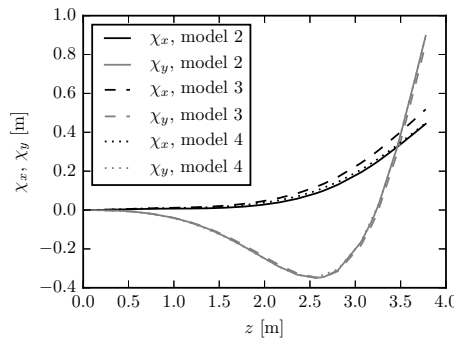


FIG. 12. Comparison of the 3rd blade mode shape. χ_x and χ_y are respectively the edgewise and flapwise deflections

484 **E. Strains**

485 All four models allow computing the blade strains. Figure 14 shows the maximum and
 486 minimum longitudinal strains computed over the blade length by these four models for the
 487 design load (defined in Section II). In addition to these data, the experimental results from
 488 the two blades tested are shown. The strain gauges were placed on the upper and lower

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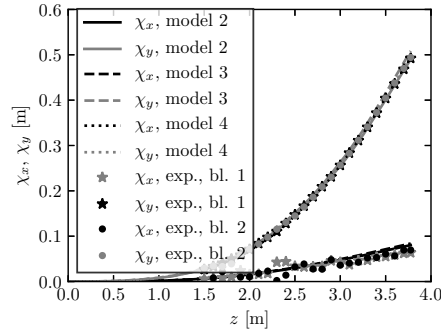


FIG. 13. Comparison of experimental and models blade deflection under the design load. χ_x and χ_y are respectively the edgewise and flapwise deflections

blade surface at 30 % of the chord length (which correspond to the chordwise blade reference
axis location). As seen in Figure 14, models 2, 3 and 4 predict the strains obtained from
the experiments relatively well (within a 10 % range). As expected and desirable, model 1
gives a conservative evaluation of the strain levels by overestimating them.

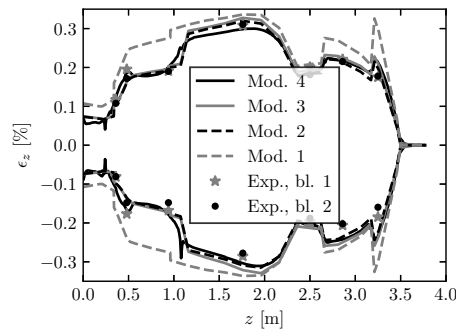


FIG. 14. Comparison of experimental and models blade strains under the design load. Top curves are lower surface data and lower curves are upper surface data

493 F. Buckling

494 During the experimentations, buckling has been observed in both steps of the destructive
495 tests (performed on both blades tested).

496 During the first step (where the blade was simply supported at $z = 2360$ mm and loaded
497 at $z = 3500$ mm), buckling was observed on the first blade tested at location $z = 2800$ mm,
498 on the panel near the trailing edge on the airfoil upper surface. From the 3D scanner data,
499 we know that buckling occurs between 2000 and 2250 N of load applied on the blade. When
500 looking at the signal of the strain gauge closest to the buckling location, we can see a change
501 in the slope at a load of 2160 N.

502 For the second blade tested, no buckling was observed with the 3D scanner before the
503 blade failure under a load of 2000 N. But when looking at the signal of the strain gauge
504 closest to the buckling location (for blade 1), a change of slope is observed at a load of around
505 1800 N. This value corresponds to the intersection point between a line passing through the
506 initial slope and a line passing through the final slope.

507 For the second step of the destructive test (where the force was applied at $z = 2360$ mm),
508 the 3D scanner indicates that buckling occurs between 4000 N and 6000 N applied on the
509 blade, and this is the case for both blades. Figure 15 shows the buckling of blade 1 as
510 recorded by the 3D scanner (blade 2 is similar). As we can see, the center of the wave that
511 has the maximum amplitude is located at $z = 1200$ mm, where a strain gauge was installed.
512 Again, buckling occurs on the panel near the trailing edge on the airfoil upper surface of
513 the blade. When looking at the signal of this strain gauge for both blades, we can see a
514 change in the slope (intersection point between lines passing through the initial and the
515 second linear parts of the curve) at a load of 4600 N for blade 1 and at a load of 5500 N for
516 blade 2.

FIG. 15. Buckling at $z = 1200$ mm on blade 1 during the second step of the destructive test
recorded by the 3D scanner

517 As summarized in Table V, for the first step of the destructive tests, buckling occurs at a
518 load level between 1800 and 2160 N at a radial location $z = 2800$ mm. For the second step
519 of the destructive tests, buckling begins at a load between 4600 and 5500 N at 1200 mm

TABLE V. Comparison of buckling results

| Model | Step 1 | | Step 2 | |
|--------------|--------------------|-------------------------|--------------------|-------------------------|
| | Buckling force [N] | Buckling loc., z [mm] | Buckling force [N] | Buckling loc., z [mm] |
| Experimental | 1800–2160 | 2800 | 4600–5500 | 1200 |
| Model 2 | 2326 | 2662 | 3702 | 1160 |
| Model 4 | 3066 | 2695 | 6059 | 1135 |

520 from the blade root.

521 The shell finite element model of the blade (model 4) allows computing buckling loads.
 522 Figure 16 shows the first buckling mode for both steps of the destructive test. The buckling
 523 loads are 3066 N and 6059 N for the first and second steps of the destructive test respectively.
 524 These results are also presented in Table V.

(a)

(b)

FIG. 16. Buckling results from model 4. (a) step 1, (b) step 2

525 The only other model able to compute buckling loads is the classical strength of materials
 526 model (model 2). For the first step of the destructive test, this model predicts a buckling
 527 load of 2326 N at the section located at 2662 mm from the blade root. For the second step,
 528 buckling occurs at 1160 mm from the root at a load level of 3702 N. These results are also
 529 summarized in Table V.

530 When comparing the buckling results, we can first see that the buckling locations are
 531 relatively well predicted by both models. All results are within ranges of 105 mm for step 1
 532 and 65 mm for step 2. If we compare the buckling loads, we can see that model 2 predict
 533 lower loads than model 4. However, the experimental data for step 1 show that buckling
 534 occurs at a load that is lower than the one predicted by both models. For the second step,
 535 the experimental results range between the results of models 2 and 4.

536 On one hand, the buckling hypothesis for model 2 is conservative. As the blade panels are
 537 curved, the flat plate assumption leads to lower buckling loads than if a solution for curved

538 panel was used²⁸. However, on the other hand, the assumption on the panel boundary con-
539 ditions can be non-conservative in some case. Here, simply supported boundary conditions
540 are used on all sides. But with the one-shear-web configuration, the trailing edge panel is
541 supported only by the spar cap (a thicker laminate) and not by a shear web. This leads
542 to a boundary condition that lies somewhere between a free edge and a simply supported
543 edge, which results in less support than the model used to get the buckling load. It probably
544 explains the fact that model 2 gives a non-conservative buckling load for step 1. For step
545 2, the buckling occurs at a location where the rigidity difference between the trailing edge
546 panel and the spar cap is higher, resulting in boundary conditions that are closer to the
547 simply supported assumption and explaining the conservative result obtained from model 2.

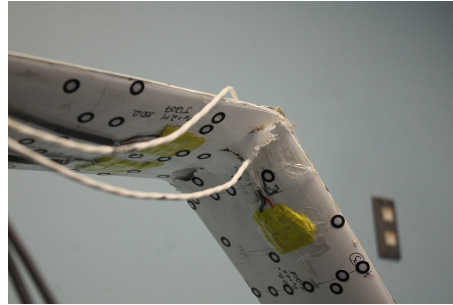
548 When looking at the results of model 4, the ratio of shell finite element buckling loads to
549 experimental minimum buckling loads is 1.70 for step 1 and 1.32 for step 2. The differences
550 between the experimental data and the results from the shell finite element model can be
551 explained partly by the fact that, as seen in the previous results (frequencies, deflections,
552 strains), the models overestimate the blade stiffness. The buckling results go in the same
553 way. Another important aspect explaining the results of model 4 is that, as reported in
554 the literature^{20,23}, finite element linear buckling analyses are non-conservative. For instance,
555 Bergreen et al.²⁰ report a non-linear buckling load as low as 65 % of the linear buckling
556 load depending on the induced imperfections. According to that, the buckling load values
557 obtained from model 4 seems reasonable. We can also note that these buckling behaviors
558 occur in non-structural areas and do not lead to blade failure. Finally, it is worth noting
559 that the non-conservative aspect of the linear finite element buckling results is formalized
560 by some standards^{29,30} where a partial safety factor of 1.25 is specified when using this type
561 of buckling analysis.

562 G. Composite Strength

563 The last object of comparison between the different blade structural models is about the
564 blade strength. For the first step of the destructive test, both blades failed in similar ways.
565 A compressive failure occurs in the spar cap on the upper side. For blade 1, the failure occurs
566 for a load around 2650 N (continuous recording of the load cell data was not available so an
567 estimate of the failure load is given) at 2760 mm from the root. Blade 2 fails at a load level

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(a)



(b)

FIG. 17. Failure of blade 1 in the first (a) and second (b) step of the destructive test

568 of 2110 N and the failure is located at 2690 mm from the root. These values are summarized
569 in Table VII presented at the end of this section. Figure 17a shows images of the failure of
570 blade 1.

571 During the second step of the test, blade 1 fails at a load level of 10 528 N and the failure
572 is located at 1210 mm from the blade root in the spar cap of the upper side (see Figure 17b).

573 The second blade fails in a different way. The failure process starts by a crack appearing
574 on the leading edge around 450 mm from the root at a load level of 8450 N. After reaching a
575 maximum load of 9040 N, the trailing edge suddenly opens at 700 mm from the blade root.
576 At this moment, the load slightly decreases. After a small increase in the applied load, the
577 blade fails at 500 mm from the root. This results in a failure of the shear web and of both
578 upper and lower skins near the trailing edge.

579 As the failure process of the second blade during step 2 of the destructive test is hard to

580 analyze with the numerical model used in this article, only the results of the first blade are
581 used for the comparison of this section. They are reported in Table VII.

582 For the evaluation of the blade strength with models 1 and 2, the strength of the different
583 laminates is needed. This is obtained by using the classical lamination theory with the
584 following procedure as proposed in composite textbooks^{13,14} :

- 585 1. The first ply failure stress is computed using the Tsai-Wu criterion.
- 586 2. For that ply, the maximum stress failure criterion is used to get the failure mode (longi-
587 tudinal tension, longitudinal compression, transverse tension, transverse compression
588 or shear).
- 589 3. If the failure mode is in the transverse direction or in shear, the stiffness dominated
590 by the matrix properties (E_2 , G_{12} and ν_{12}) are set to 0. If the failure mode is in the
591 longitudinal direction, the stiffness dominated by the fiber properties (E_1 , and ν_{12}) are
592 set to 0.
- 593 4. This process is repeated until the maximum load is reached.
- 594 5. The failure strain is computed as the failure stress divided by the initial longitudinal
595 modulus.
- 596 6. The failure analysis is applied to the 0° and $\pm 45^\circ$ plies only.

597 Table VI summarizes the tension and compression longitudinal failure strains for all
598 laminates along the blade length. All compressive strains are lower than tensile strength.

599 When using the compressive failure strains within structural model 1, for the first step of
600 the destructive test, the blade failure is predicted at $z = 3211$ mm under a force of 1288 N.
601 A second possible failure point is located at $z = 2662$ mm and arise when the force reaches a
602 value of 1705 N. This second failure point is interesting because it is located near the failure
603 location observed during the tests. For the second step of the destructive test, the failure is
604 predicted at $z = 1360$ mm under a force of 9661 N. These results are reported in Table VII.

605 The results obtained from structural model 2 are similar to those of model 1. For model
606 2, due to the asymmetry of the blade cross section, the extremum cross-sectional strains are
607 not the same in tension and in compression. The predicted failures are in compression due
608 to the fact that the failure strains are smaller in compression than in tension. As presented

TABLE VI. Failure strains of the laminates along the blade length. e_x^T and e_x^C are respectively the tensile and compressive longitudinal failure strains

| longitudinal position | e_x^T [%] | e_x^C [%] |
|---------------------------|-------------|-------------|
| $z < 337$ mm | 1.32 | 0.89 |
| 337 mm $< z < 960$ mm | 1.61 | 1.02 |
| 960 mm $< z < 2624$ mm | 1.59 | 1.01 |
| 2624 mm $< z < 3173$ mm | 1.45 | 0.94 |
| 3173 mm $< z$ | 1.07 | 0.84 |

609 in Table VII, the predicted failures of structural model 2 are at the same locations as those
 610 of model 1, but for higher load values. So, as for the bending stiffness evaluation, model 1
 611 gives conservative estimation of the failure load.

612 For models 3 and 4, the method used for the evaluation of blade failure load is based on
 613 a method described by Barbero¹². In order to get the failure loads, all the materials' stiff-
 614 nesses were given a value close to zero, except the longitudinal properties of the glass/epoxy
 615 unidirectional plies, the transverse shear of the glass/epoxy unidirectional plies and the core
 616 material properties. By doing so, the blade behaves as if transverse failures have occurred
 617 so that all the load is carried by the fibers. To avoid numerical problems associated with
 618 zero stiffness deformation modes, the transverse shear properties of the glass/epoxy unidi-
 619 rectional plies and the core properties were also kept unchanged. All the material strengths
 620 were set to really high values excepted the tensile and compressive longitudinal strength of
 621 the glass/epoxy unidirectional plies to force the solver to compute failure indices associated
 622 with these failure modes. The Tsai-Wu failure index F_{TW} is used. This index is the inverse
 623 of the safety factor. Note that, in this case, Tsai-Wu and maximum stress failure criterion
 624 give the same results. This gives a conservative estimation of the last ply failure strength
 625 by using a linear model.

626 When performing this analysis with structural model 3, for the first step of the test, a
 627 failure is predicted at $z = 3211$ mm under a force of 1452 N. Another possible failure point
 628 is located at $z = 2662$ mm and the failure occurs at a 1942 N load level. For the second step
 629 of the destructive test, a failure is predicted under a force of 11 270 N at 1360 mm from the
 630 blade root. For both steps of the test, the failure occurs in compression on the upper side's

631 spar cap. As an example, Figure 18 shows the distribution of the composite failure index in
632 the cross section located at $z = 1360$ mm for the second step of the destructive test under
633 a unit load. The failure results from model 3 are presented in Table VII.

FIG. 18. Distribution of the composite failure index from model 3 (maximum value among each layer) in the cross section located at $z = 1360$ mm for the second step of the destructive test under a unit load. The inverse of the maximum failure index gives the failure load: $1/8.87 \times 10^{-5} = 11\,270$

634 The failure results from model 4 are summarized in Figure 19, showing the composite
635 failure index distribution for unit loads. For both steps, a compressive failure is predicted
636 on the upper side of the blade. Two points are identified for both steps, they correspond to
637 the two highest values of the failure index. The failure results of model 4 are summarized
638 in Table VII

(a)

(b)

FIG. 19. Failure indices from model 4 (maximum value among each layer) under unit loads. The failure load can be computed as the inverse of the failure index. (a) Step 1. (b) Step 2

639 When comparing the results of Table VII, we can see that the failure locations are rela-
640 tively well predicted by all models. Sometimes, the first failure predicted is not exactly at
641 the experimental failure point, but a second point of high failure index is located close to
642 the experimental failure point. The predicted failure loads are conservative or close to the
643 observed values. As expected, the results from model 1 are the most conservative, but give
644 very good insight on the failure behaviour despite the model simplicity. As also expected,
645 model 4 is the most precise.

646 V. CONCLUSION

647 In conclusion, the simple model (model 1), despite its simplicity, gives fair results for the
648 local flapwise bending stiffness, the strain distribution and the blade failure. In addition to

TABLE VII. Comparison of blade strength results

| Model | Step 1 | | Step 2 | |
|--------------|------------------|------------------------|------------------|------------------------|
| | Failure load [N] | Failure loc., z [mm] | Failure load [N] | Failure loc., z [mm] |
| Experimental | 2110–2650 | 2690–2760 | 10 530 | 1210 |
| Model 1 | 1288, 1705 | 3211, 2662 | 9661 | 1361 |
| Model 2 | 1723, 1925 | 3211, 2662 | 11 010 | 1360 |
| Model 3 | 1452, 1942 | 3211, 2662 | 11 270 | 1360 |
| Model 4 | 1704, 2249 | 3215, 2687 | 10 850, 13 080 | 1135, 1273 |

649 that, it returns conservative results. However, this model does not allow the evaluation of
 650 the blade's natural frequencies, deflection (a conservative evaluation of the deflection of the
 651 untwisted blade could be performed, but was not done here) and buckling. This model is
 652 well suited for preliminary design based on blade strength.

653 The classical strength of materials model (model 2) gives good results for cross-sectional
 654 properties, natural frequencies, deflection, strain and composite failure. Buckling has to be
 655 handled with care as non-conservative results are obtained despite the conservative hypoth-
 656 esis of the model. Also, this model is limited to blades using orthotropic laminates (i.e., no
 657 material couplings at the laminate level). Care must also be paid to the region where the
 658 shape of the cross section is varying quickly. The stress, strain and failure index are not
 659 accurate in these regions.

660 The cross-sectional finite element model (model 3) gives the more reliable results for the
 661 cross-sectional properties, especially for properties associated with transverse shear. This
 662 model is well suited for beams that use material coupling (bend-twist coupling for wind
 663 turbine blades for instance). It gives good results for natural frequencies, deflections, strains
 664 and composite strength, but as the classical strength of materials model, it suffers from a
 665 lack of precision in the region where the shape of the cross section is varying quickly. No
 666 buckling analysis was implemented within model 3, but it could be possible to implement
 667 something similar to model 2, with similar results.

668 These latter two models are well suited for more detailed design and for providing the
 669 blade cross-sectional properties needed for the aeroelastic analysis without having to build
 670 a 3D shell finite element model, which is much more time consuming.

671 The 3D shell finite element model (model 4) is the more precise and is the only one
672 that can manage the regions of the blade where the cross-sectional shape changes quickly.
673 However, it is less reliable than model 3 for the computation of cross-sectional properties
674 associated with transverse shear. The results of the buckling analysis also have to be used
675 with caution since they are not conservative. The 3D finite element model is the ideal model
676 for the final validation of the blade.

677 When looking at the time to build and run these different models, we can say that models
678 1, 2 and 3 has similar build times. Both of these models require a tabular description of the
679 blade. Assuming that the blade aerodynamic shape is known (in a tabular format), that
680 the airfoil contour point files are available and that the material properties and layups are
681 defined, these models can be built in minutes. As model 1 require only the informations
682 about spar caps, it is even shorter than that. Model 4 can be built as quickly as model
683 2 and 3 if a tool generating finite element model from a tabular description of the blade
684 is available. However, for the final stage of the design process, when a geometric model
685 needs to be done using a CAD software and then, meshed, the built time increases to hours.
686 In regards to the simulation time, models 1, 2, 3 and 4 take respectively 0.1 s, 4.5 s, 28
687 min and 3 min 25 s to solve two load cases on the same computer. These times include
688 buckling analysis for models 2 and 4. It is important to note that the considerable solving
689 time for model 3 is due to the fact it uses a fine mesh and that this in-house code uses an
690 interpreted language (Python) and no code optimization has been done yet. Its conversion
691 into a compiled language would reduce the simulation time.

692 To conclude this paper, the wind turbine blade design process presented in Figure 20 is
693 proposed. The inner circle represents the very first stages of the design process where model
694 1 can be used to validate the feasibility of an aerodynamic design and get an idea of the mass
695 distribution. A set of loads can be obtained on a standstill blade under the extreme wind
696 model without information about the blade mass or stiffness. Once the aerodynamic design
697 seems feasible, the process can go to the second circle. If no blade stiffness information
698 is available, an aeroelastic model with rigid blades can be used to get the loads. Model 2
699 or 3 can be used to get a preliminary structural design and to compute the blade stiffness
700 properties. The process can then enter the outer circle, where model 4 is used for the blade
701 dimensioning and validation and where the aeroelastic model uses flexible blades.

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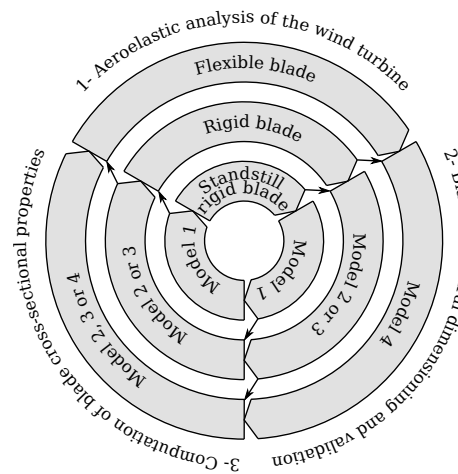


FIG. 20. Proposed wind turbine blade design process: The inner circle represents the very first stages of the design process that will evolve towards the outer circles as the blade design refines

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705 **AUTHOR DECLARATIONS**

706 **Declaration of conflicts of interest**

707 The authors declared no potential conflicts of interest with respect to the research, au-
708 thorship, and/or publication of this article.

709 **REFERENCES**

- 710 ¹David J. Malcolm and Daniel L. Laird. Extraction of equivalent beam properties from
711 blade models. *Wind Energy*, 10:135–157, 2007.
- 712 ²Louis-Charles Forcier, Jon Sumner, Tommy Gagnon, Jean-François Charron, and Simon
713 Joncas. Structural design and validation of a 10 kW wind turbine blade. In *19th Inter-*

This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.

PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0087613

Accepted to J. Renew. Sustain. Energy 10.1063/5.0087613

- 714 *national Conference on Composite Materials*, Montréal (Québec, Canada), 28 July – 2
715 August, 2013.
- 716 ³Louis-Charles Forcier. *Development and validation of tools for the structural design of wind*
717 *turbine blades*. PhD thesis, École de technologie supérieure, Montréal (Canada), 2020.
- 718 ⁴G. R. Cowper. The shear coefficient in Timoshenko's beam theory. *Journal of Applied*
719 *Mechanics*, June:335–340, 1966.
- 720 ⁵Philippe Giguère and Michael S. Selig. Blade geometry optimization for the design of wind
721 turbine rotor. In *2000 ASME Wind Energy Symposium*, Reno (Nevada, United States),
722 10–13 January, 2000.
- 723 ⁶B. Hillmer, T. Borstelmann, P.A. Schaffarczyk, and L. Dannenberg. Aerodynamic and
724 structural design of multiMW wind turbine blades beyond 5 MW. In *The Science of*
725 *Making Torque from Wind*, Lyngby (Denmark), 28–31 August, 28–31 août 2007.
- 726 ⁷Lars P. Mikkelsen. A simplified model predicting the weight of the load carrying beam in a
727 wind turbine blade. In *37th Risø International Symposium on Materials Science*, Roskilde
728 (Denmark), 5–8 September, 2016.
- 729 ⁸Martin O. L. Hansen. *Aerodynamics of Wind Turbines*. Earthscan, London (United
730 Kingdom), 2nd edition, 2008.
- 731 ⁹Robert M. Rivello. *Theory and Analysis of Flight Structures*. McGraw-Hill, New York
732 (New York, United States), 1969.
- 733 ¹⁰André Bazergui, Thang Bui-Quoc, André Biron, Georges McIntyre, and Charles Laberge.
734 *Résistance des matériaux*. Presses internationales Polytechnique, Montréal (Québec,
735 Canada), 3rd edition, 2002.
- 736 ¹¹László P. Kollár and George S. Springer. *Mechanics of Composite Structures*. Cambridge
737 University Press, Cambridge (United Kingdom), 2003.
- 738 ¹²E. J. Barbero. *Introduction to Composite Materials Design*. Taylor & Francis, New York
739 (New York, United States), 1999.
- 740 ¹³Jean-Marie Berthelot. *Matériaux composites: Comportement mécanique et analyse des*
741 *structures*. Lavoisier, Paris (France), 5th edition, 2012.
- 742 ¹⁴Ronald F. Gibson. *Principles of composite material mechanics*. CRC Press, Boca Raton
743 (Florida, United States), 2nd edition, 2007.
- 744 ¹⁵V. Giavotto, M. Borri, P. Mantegazza, and G. Ghiringhelli. Anisotropic beam theory and
745 applications. *Computers & Structures*, 16(1–4):403–413, 1983.

This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.

PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0087613

Accepted to J. Renew. Sustain. Energy 10.1063/5.0087613

- 746 ¹⁶José Pedro Blasques and Mathias Stolpe. Multi-material topology optimization of lami-
747 nated composite beam cross sections. *Composite Structures*, 94(11):3278–3289, 2012.
- 748 ¹⁷José Pedro Blasques. User's Manual for BECAS : A cross section analysis tool for
749 anisotropic and inhomogeneous beam sections of arbitrary geometry. Technical Report
750 Risø-R-1785(EN), Risø DTU – National Laboratory for Sustainable Energy, Roskilde (Den-
751 mark), 2012.
- 752 ¹⁸José Pedro Blasques. *Optimal Design of Laminated Composite Beams*. PhD thesis, Tech-
753 nical University of Denmark, Lyngby (Denmark), 2011.
- 754 ¹⁹J. P. Blasques, R. D. Bitsche, V. Fedorov, and B. S. Lazarov. Accuracy of an efficient
755 framework for structural analysis of wind turbine blades. *Wind Energy*, 19:1603–1621,
756 2016.
- 757 ²⁰C. Berggreen, K. Branner, J. F. Jensen, and J. P. Schultz. Application and analysis
758 of sandwich elements in the primary structure of large wind turbine blades. *Journal of*
759 *Sandwich Structures and Materials*, 9:525–551, 2007.
- 760 ²¹Find Mølholt Jensen. *Ultimate Strength of a Large Wind Turbine Blade*. PhD thesis,
761 Risø National Laboratory for Sustainable Energy and Technical University of Denmark,
762 Roskilde and Lyngby (Denmark), 2008.
- 763 ²²Patricio Lillo. Static and fatigue analysis of wind turbine blades subject to cold weather
764 conditions using finite element analysis, 2011. Master's thesis, University of Victoria,
765 Victoria (British Columbia, Canada).
- 766 ²³J. A. Paquette and P. S. Veers. Increased strength in wind turbine blades through innova-
767 tive structural design. In *AWEA Windpower 2008*, Houston (Texas, United States), 1–4
768 June, 2008.
- 769 ²⁴M. Tarfaoui, H. Khadimallah, A. Imad, and J. Y. Pradillon. Design and finite element
770 modal analysis of 48m composite wind turbine blade. *Applied Mechanics and Materials*,
771 146:170–184, 2012.
- 772 ²⁵V. Fedorov and C. Berggreen. Bend-twist coupling potential of wind turbine blades. In
773 *The Science of Making Torque from Wind 2014*, Lyngby (Denmark), 17–20 June, 2014.
- 774 ²⁶Kim Branner, Peter Berring, Christian Berggreen, and Henrik W. Knudsen. Torsional
775 performance of wind turbine blades – part II: Numerical validation. In *16th International*
776 *Conference on Composite Materials*, Kyoto (Japan), 8–13 July, 2007.

This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.

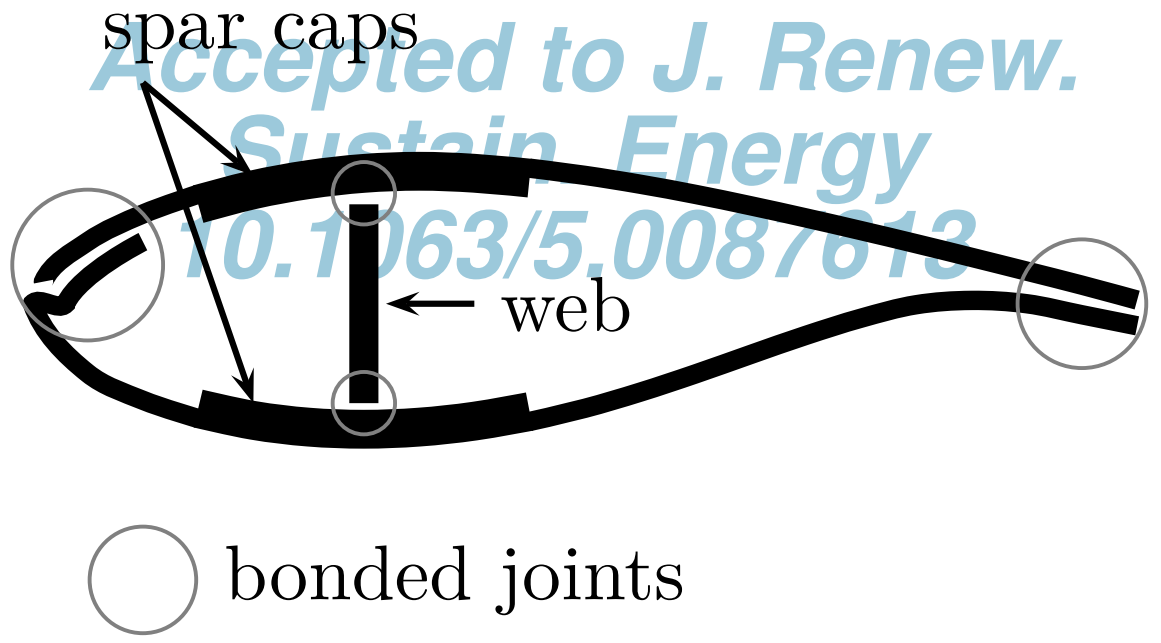
PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0087613

Accepted to *J. Renew. Sustain. Energy* 10.1063/5.0087613

- 777 ²⁷Daniel L. Laird, Felicia C. Montoya, and David J. Malcolm. Finite element modeling of
778 wind turbine blades. In *2005 ASME Wind Energy Symposium / 43rd AIAA Aerospace*
779 *Sciences Meeting and Exhibit*, Reno (Nevada, United States), 10–13 January, 2005.
- 780 ²⁸C. Lindenburg and G. D. de Winkel. Buckling load prediction tools for rotor blades. mode
781 description of tools for buckling of thin-walled beams. Technical Report ECN-C--05-103,
782 Energy Research Center of The Netherlands, Petten (The Netherlands), 2005.
- 783 ²⁹Germanischer Lloyd. Guideline for the certification of wind turbines. Guideline Edition
784 2010, Germanischer Lloyd, Hamburg (Germany), 2010.
- 785 ³⁰DNV GL. Rotor blades for wind turbines. Standard DNVGL-ST-0376 Edition December
786 2015, DNV GL, Høvik (Norway), 2015.

This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.

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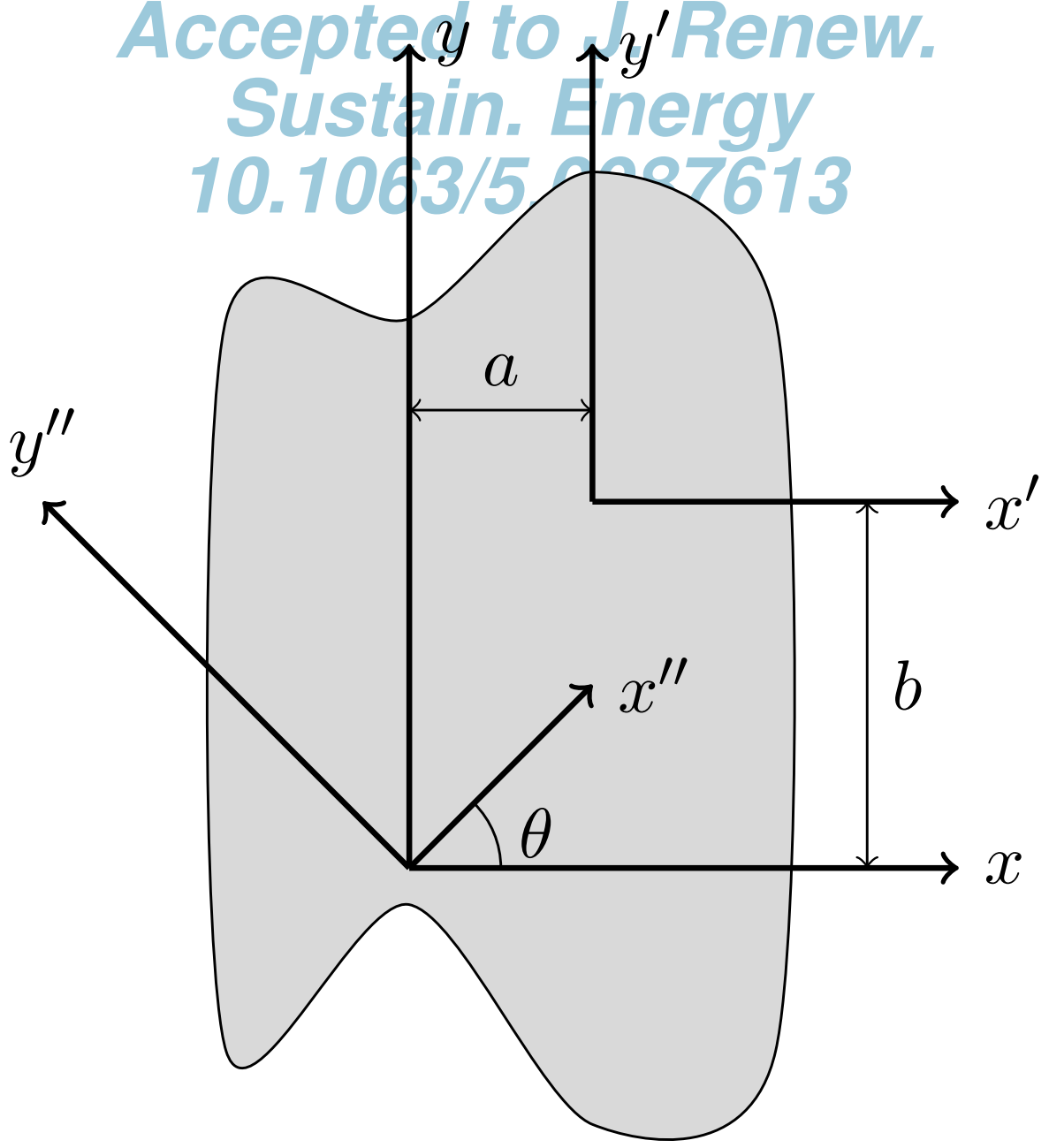
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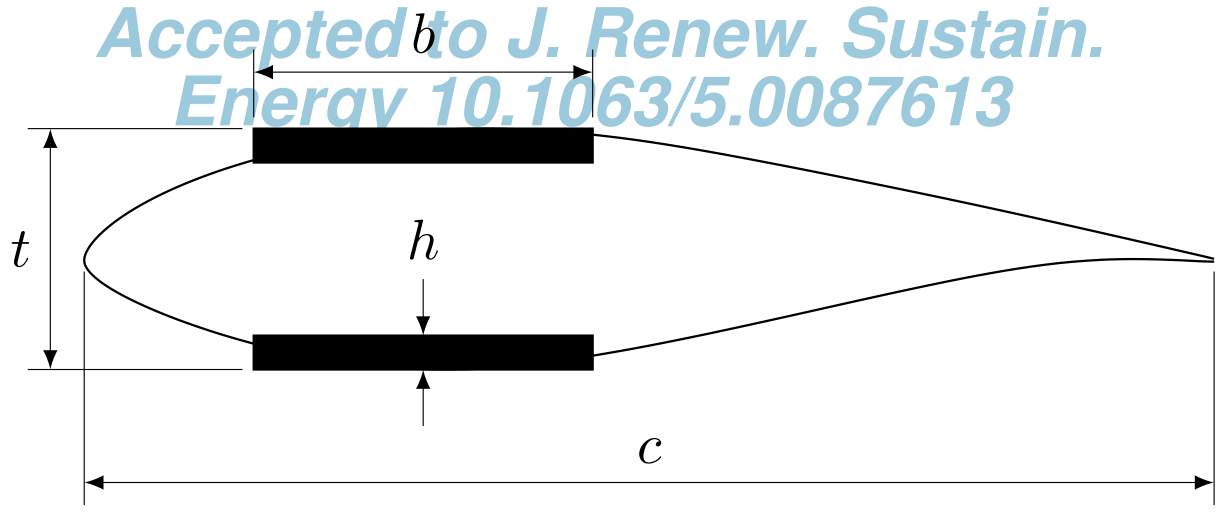
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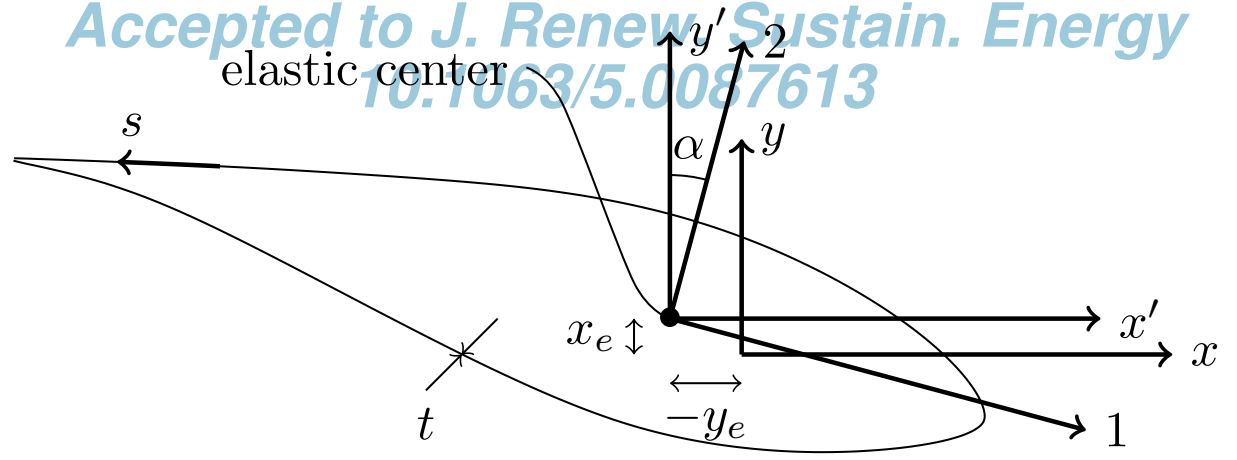
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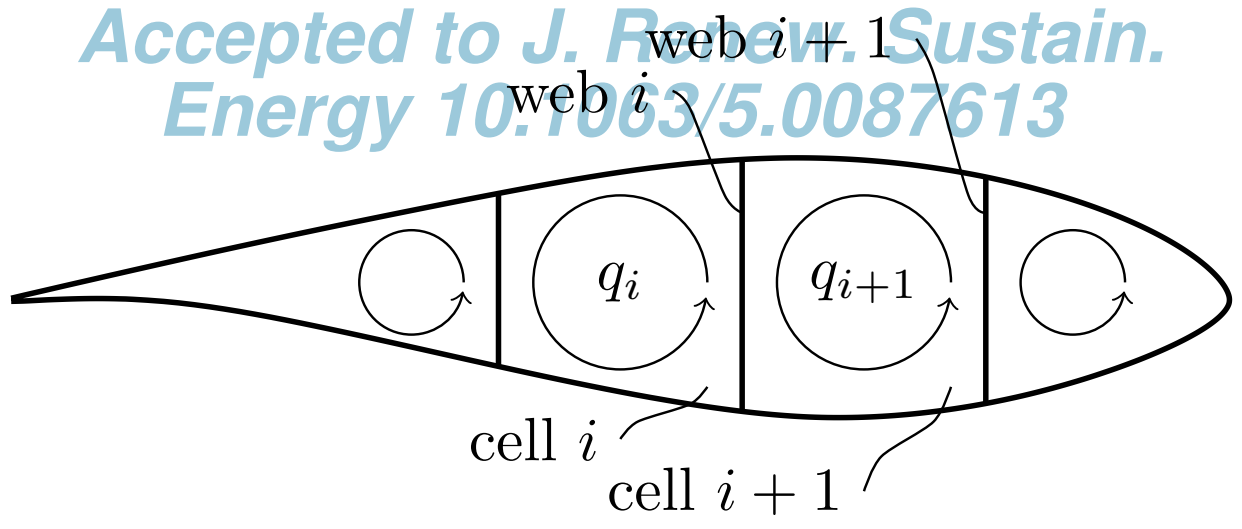
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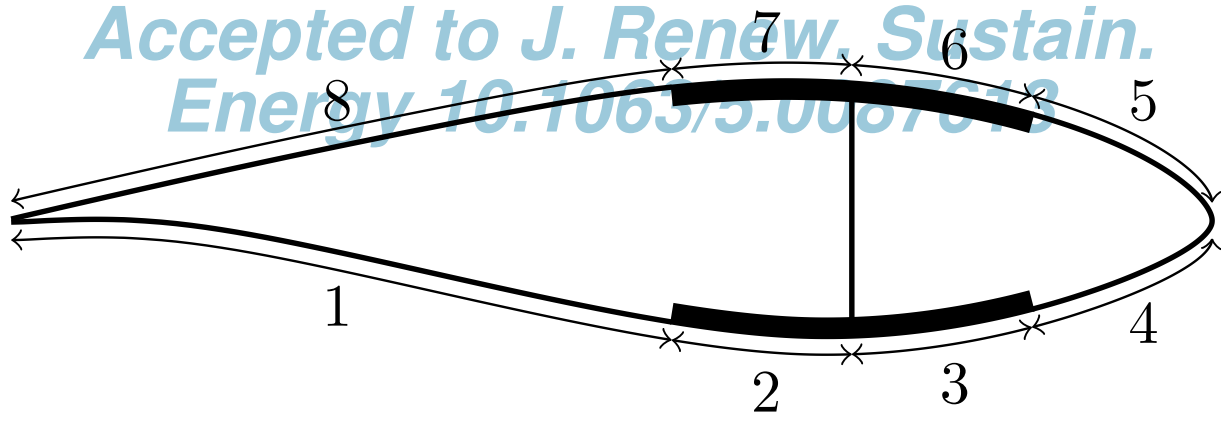
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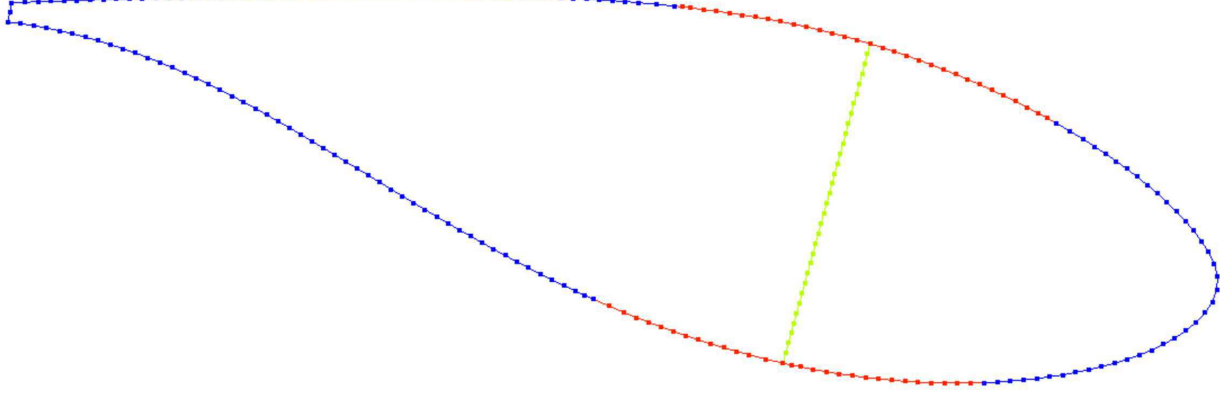
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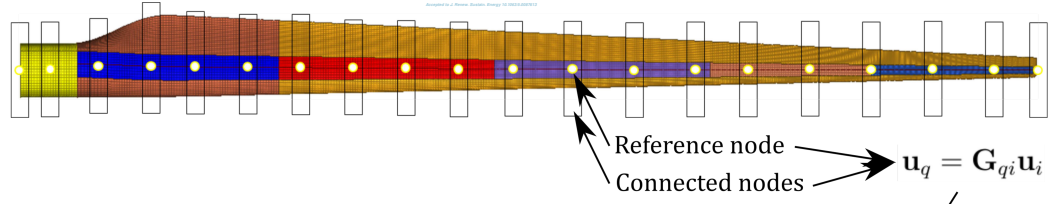
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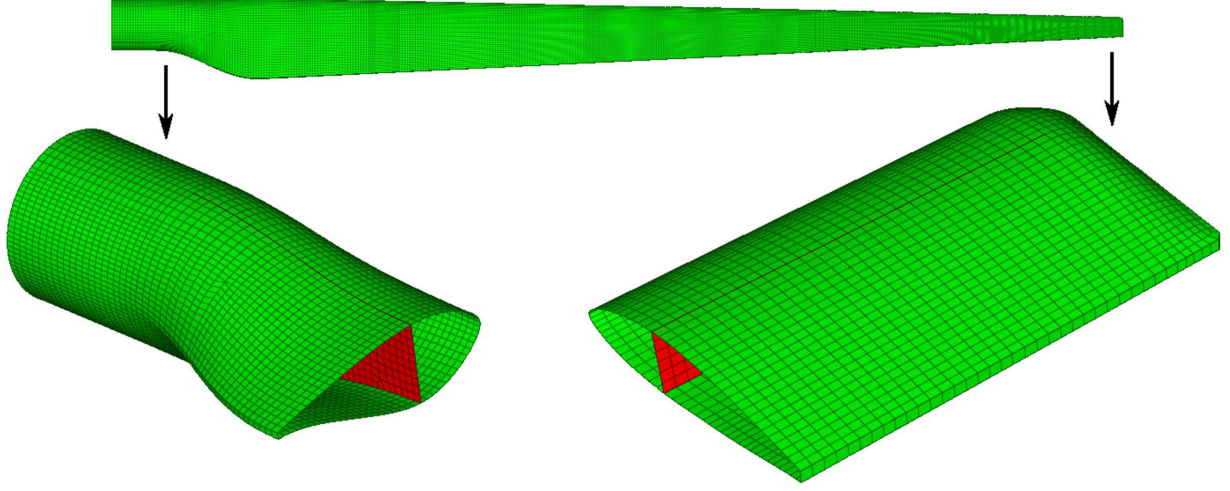
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$$\mathbf{K}_s = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & z - L & 0 & 1 & 0 & 0 \\ L - z & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \gamma_{zx}^0 \\ \gamma_{zy}^0 \\ \epsilon_z^0 \\ \kappa_x \\ \kappa_y \\ \kappa_z \end{bmatrix}_1 \dots \begin{bmatrix} \gamma_{zx}^0 \\ \gamma_{zy}^0 \\ \epsilon_z^0 \\ \kappa_x \\ \kappa_y \\ \kappa_z \end{bmatrix}_6^{-1} \leftarrow \begin{matrix} \gamma_{zx}^0 = \frac{\partial u_{0x}}{\partial z} - \theta_y \\ \gamma_{yz}^0 = \frac{\partial u_{0y}}{\partial z} + \theta_x \\ \epsilon_z^0 = \frac{\partial u_{0z}}{\partial z} \\ \kappa_x = \frac{\partial \theta_x}{\partial z} \\ \kappa_y = \frac{\partial \theta_y}{\partial z} \\ \kappa_z = \frac{\partial \theta_z}{\partial z} \end{matrix} \text{ Numerical derivatives}$$



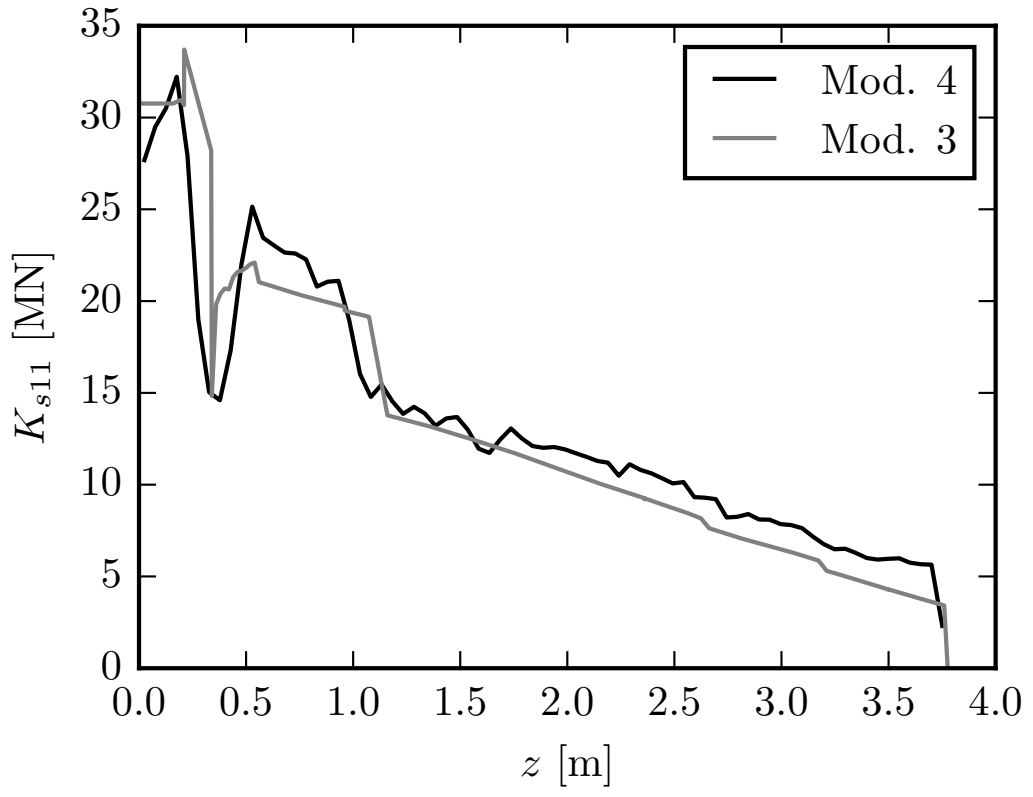
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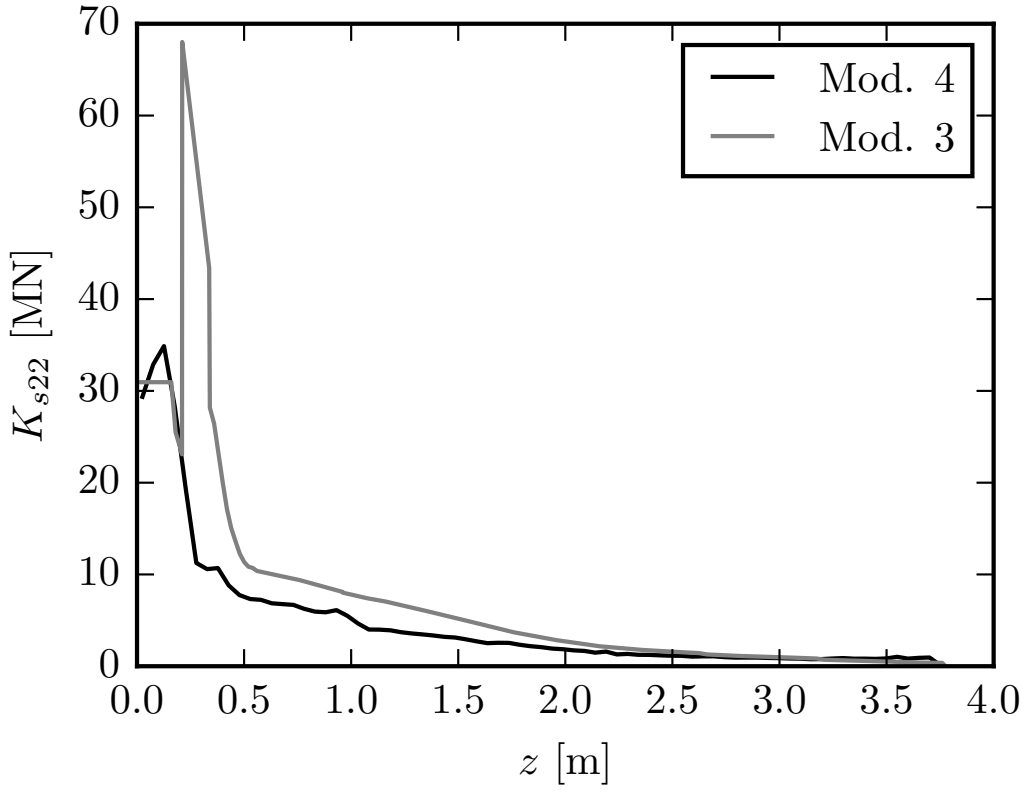
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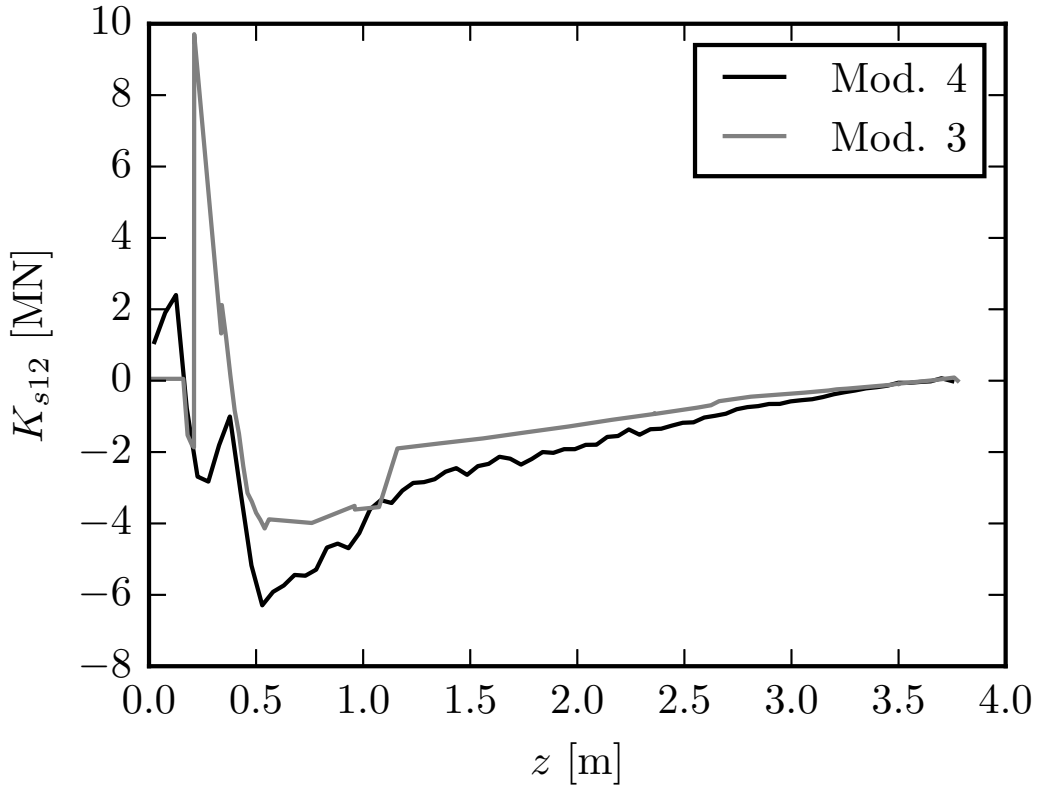
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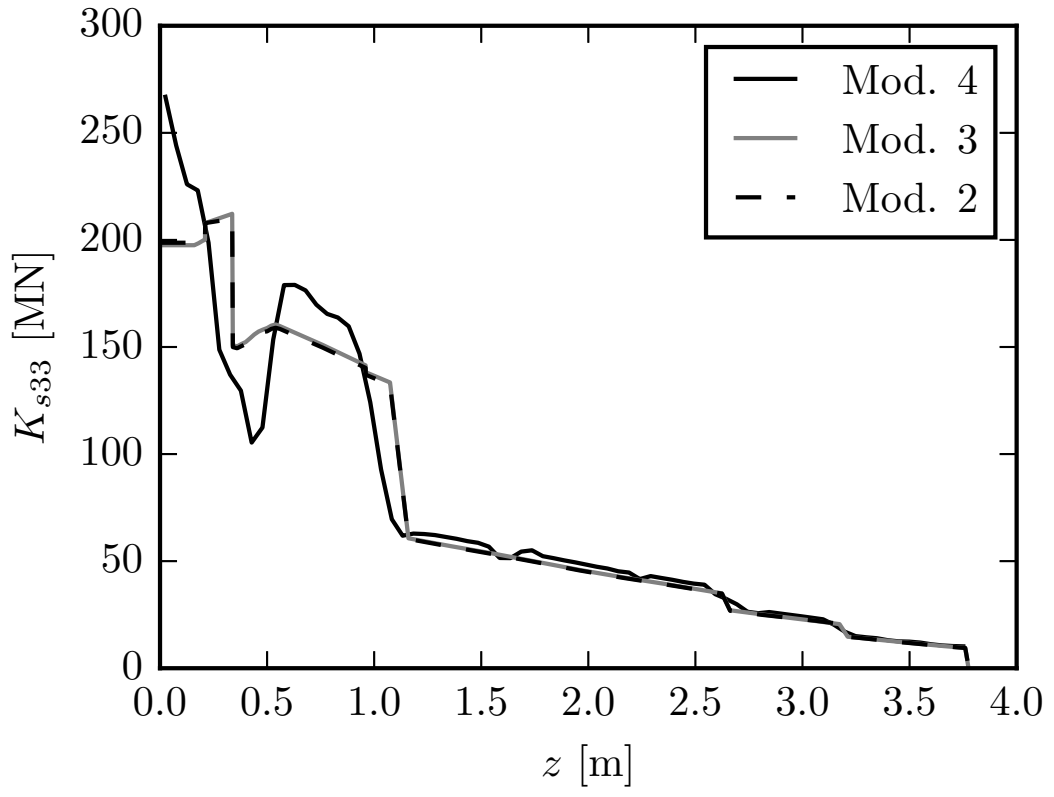
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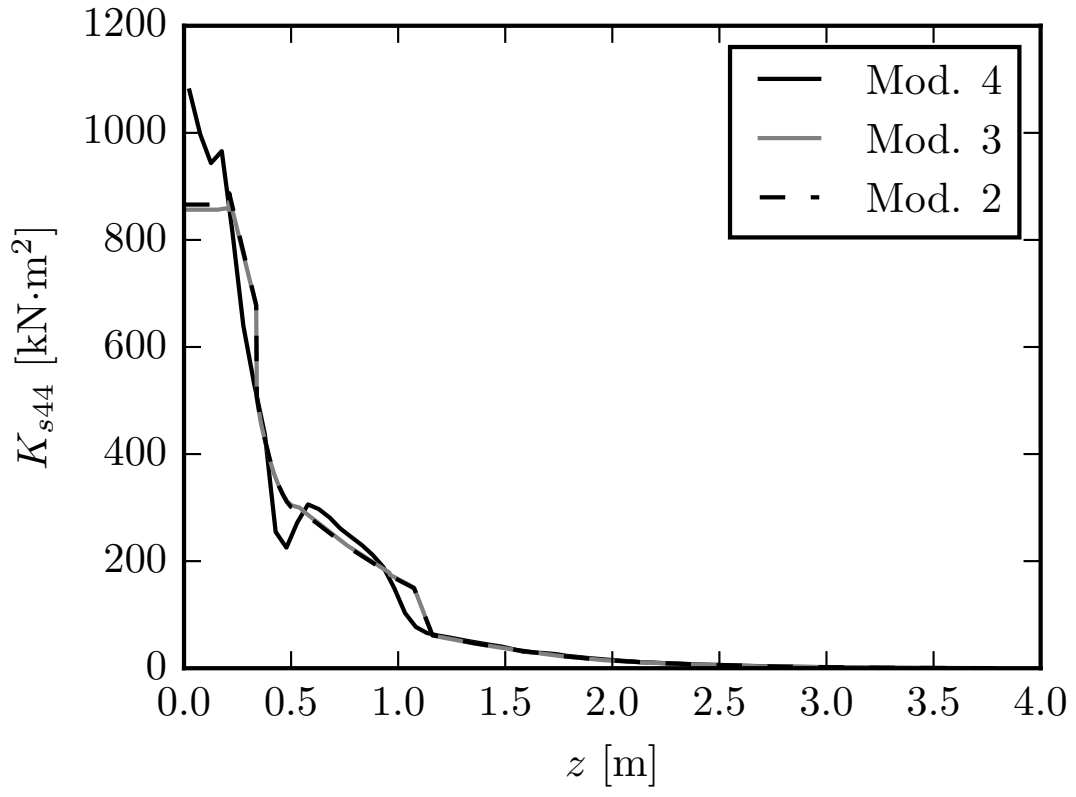
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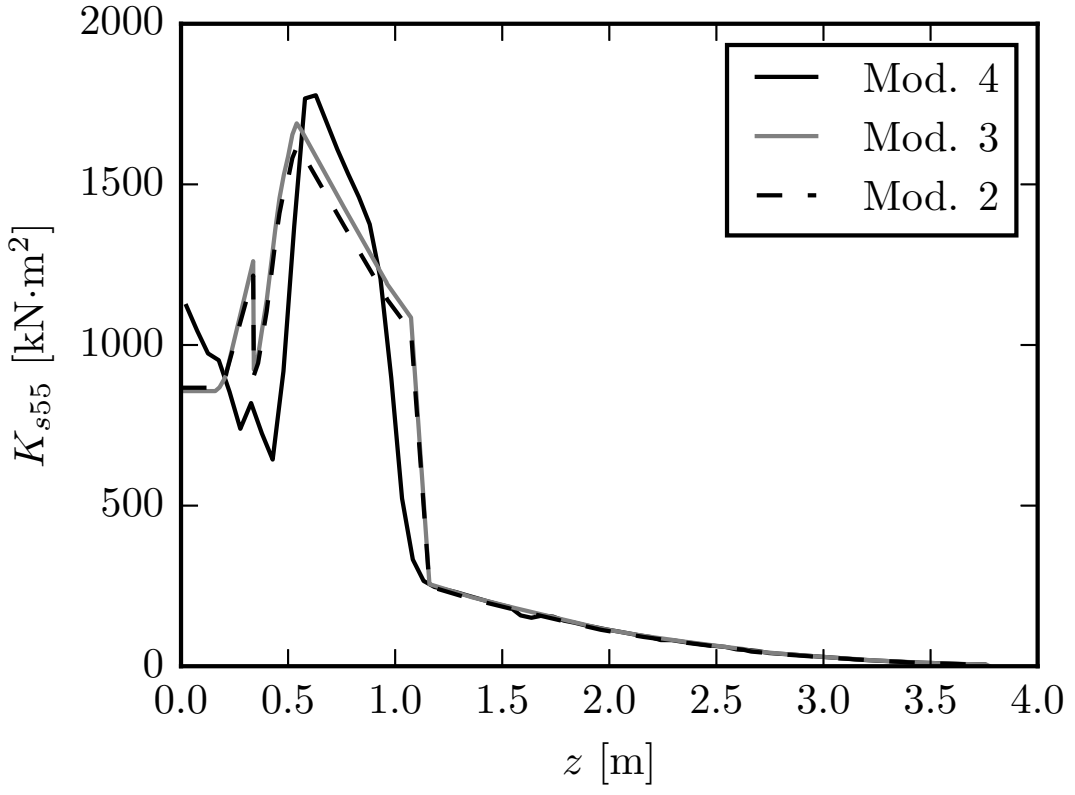
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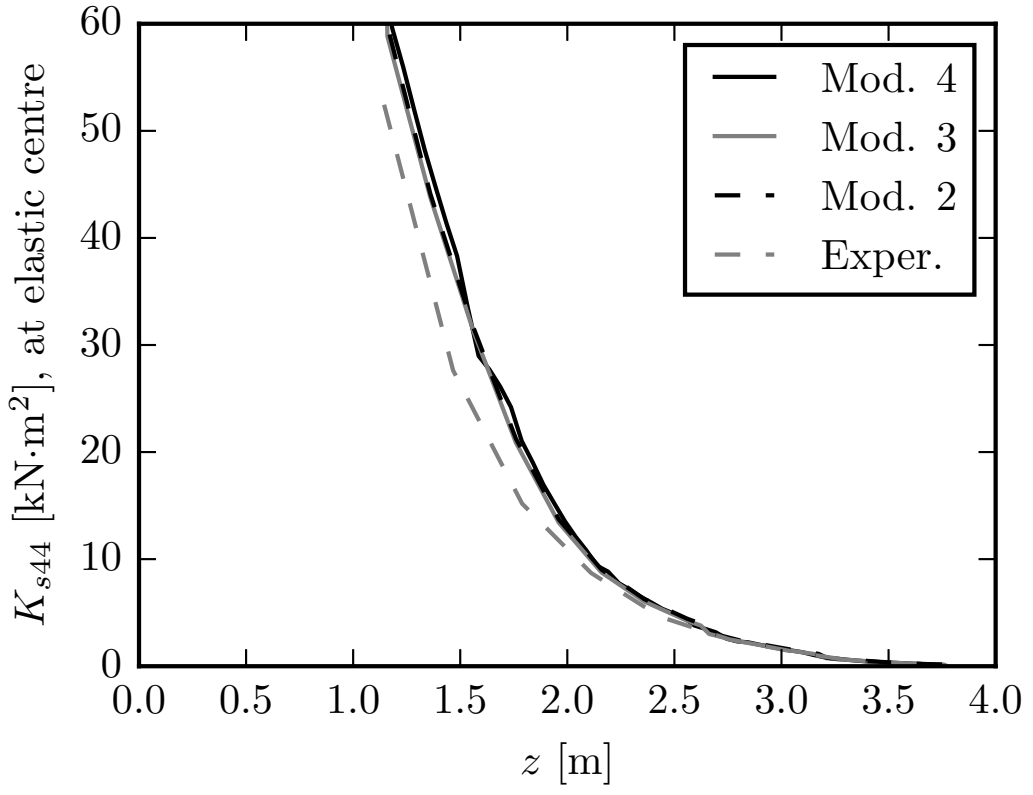
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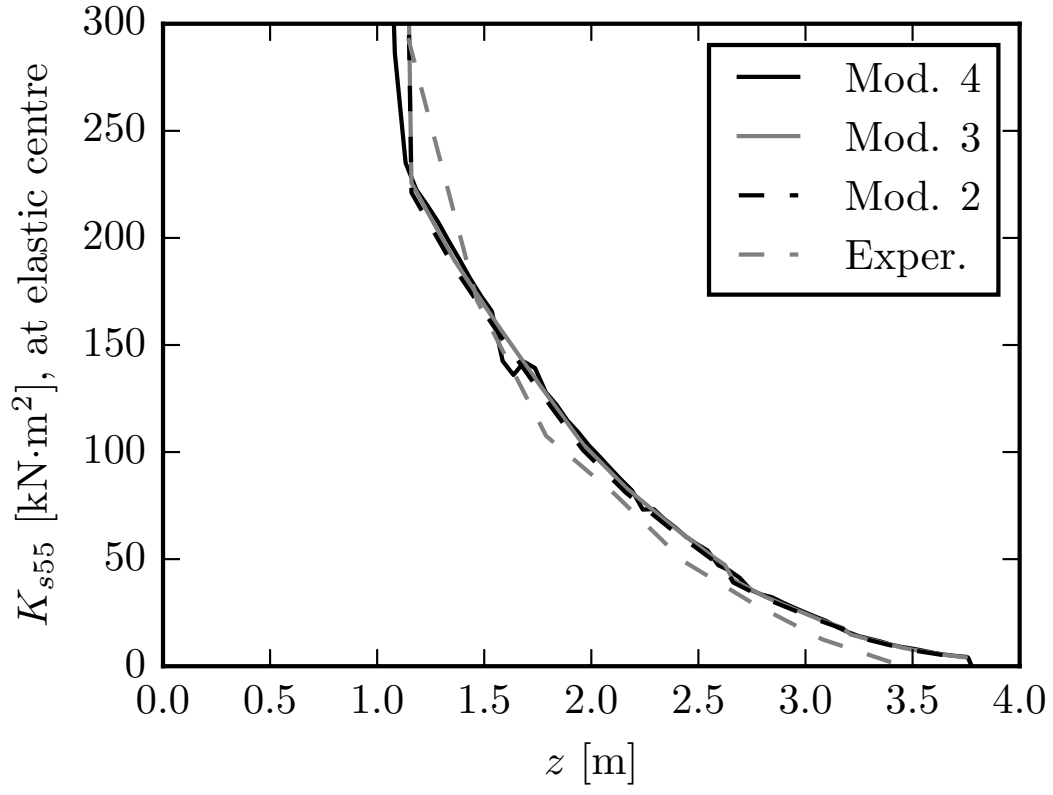
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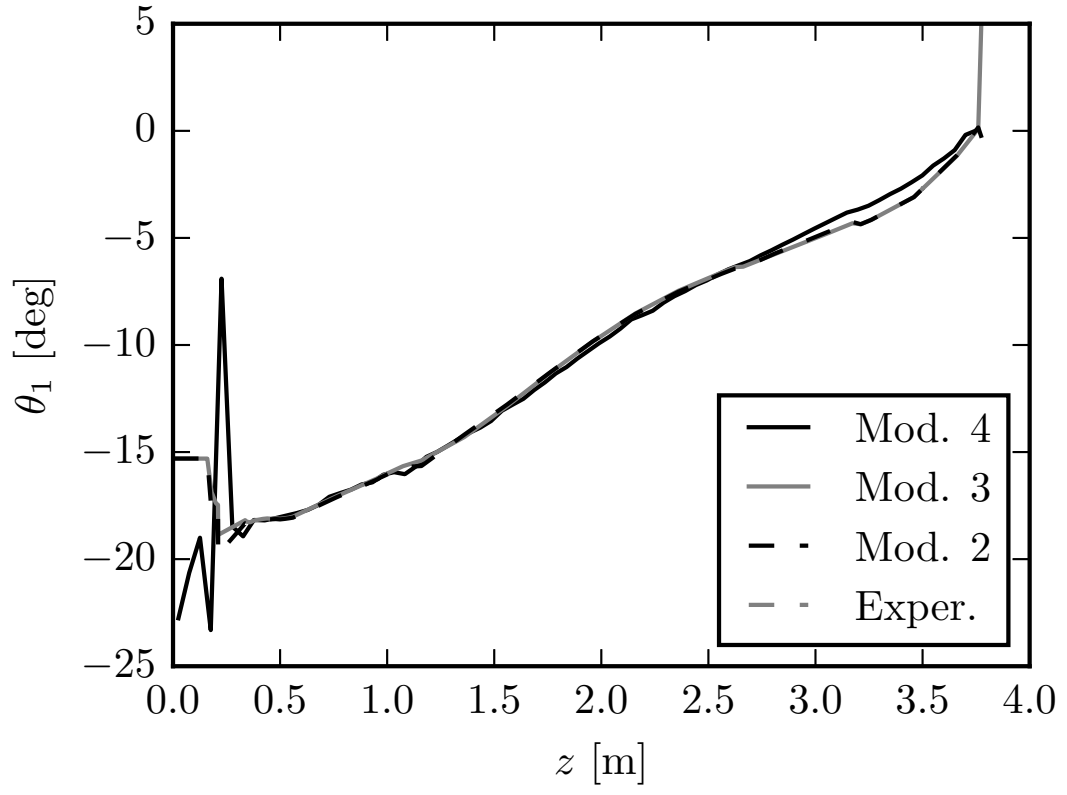
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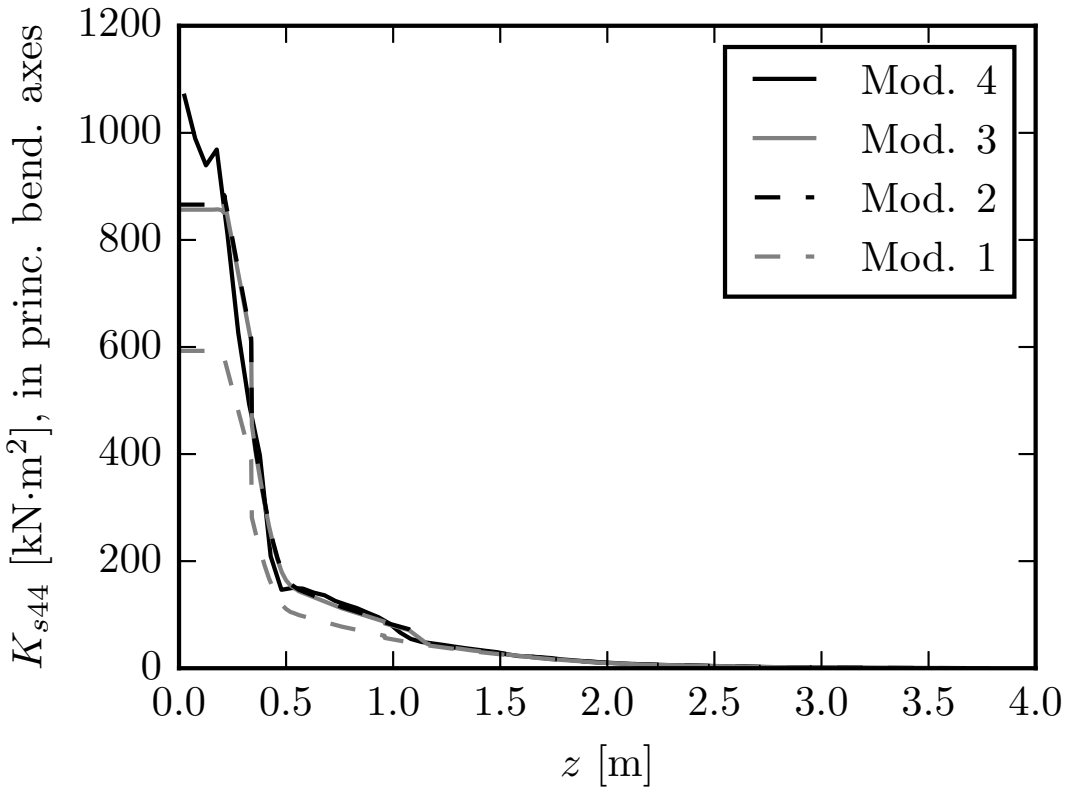
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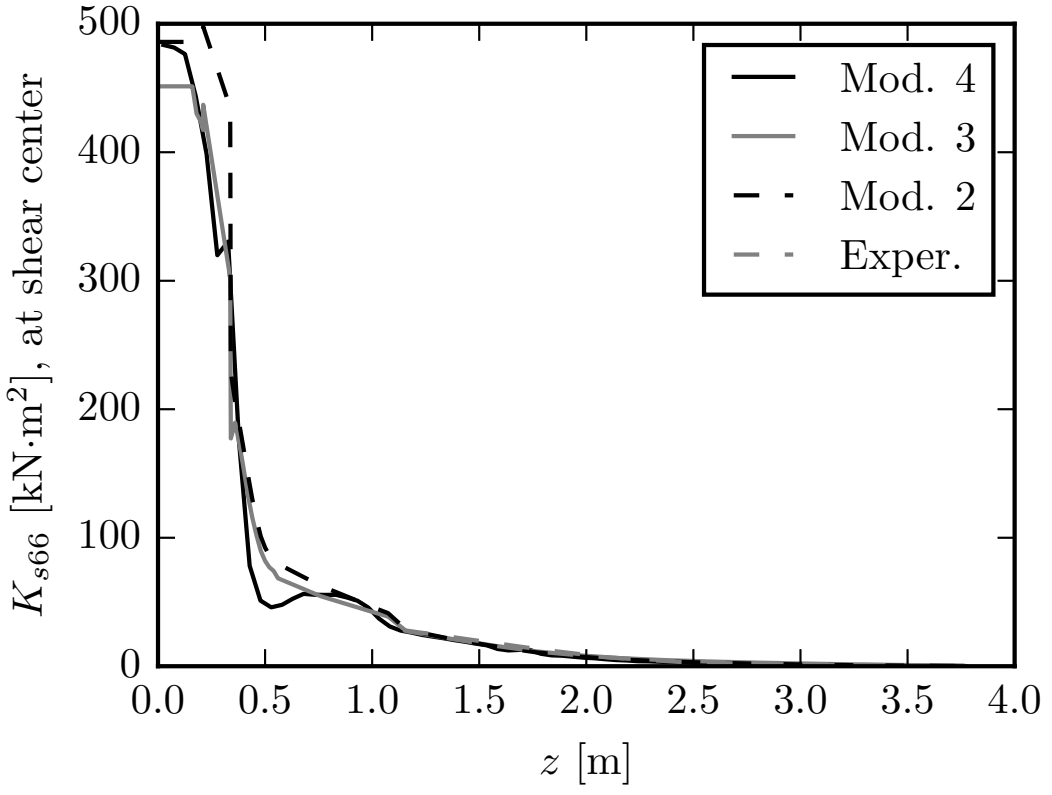
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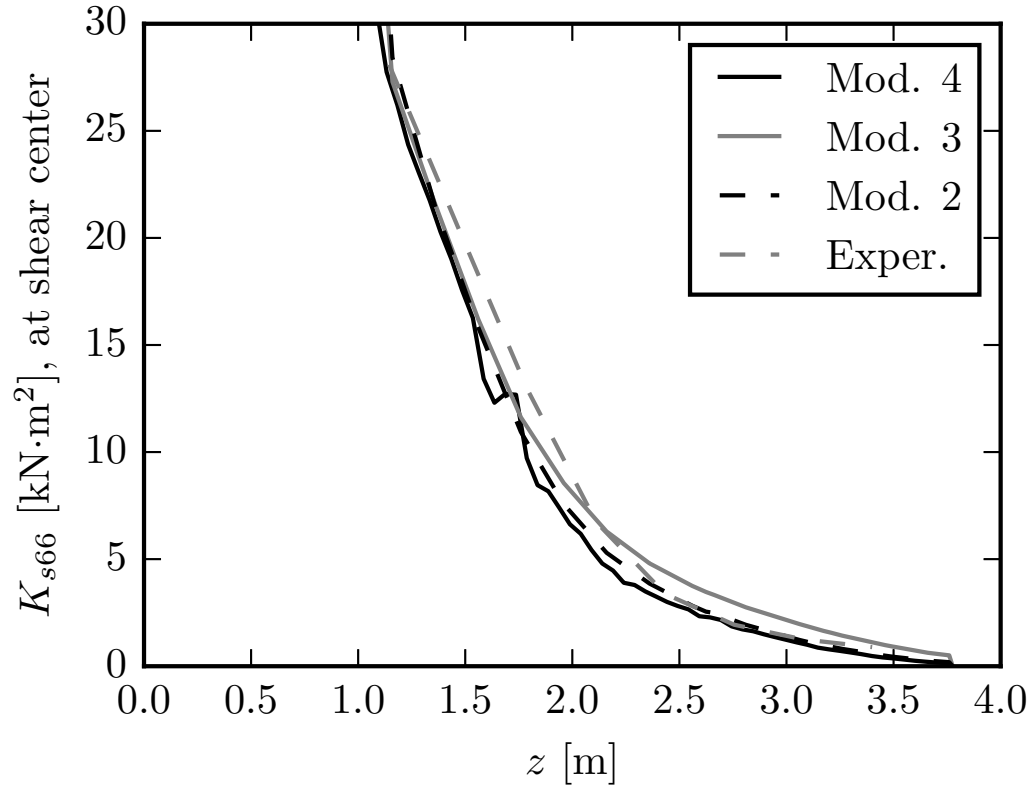
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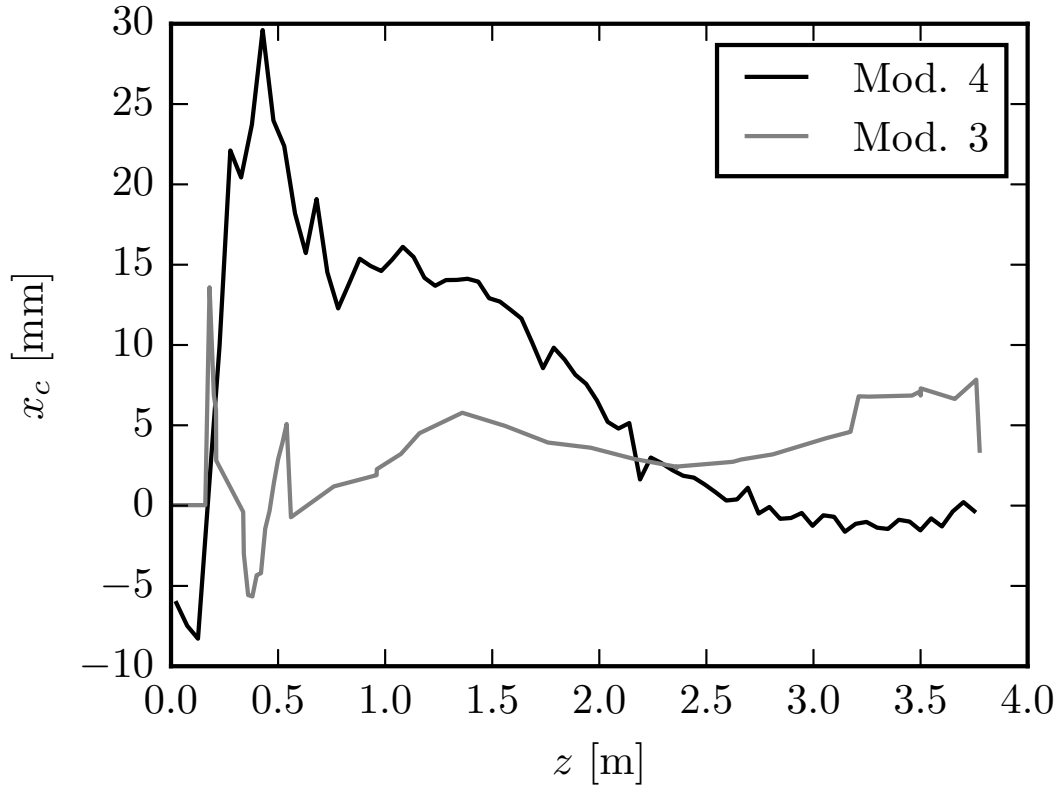
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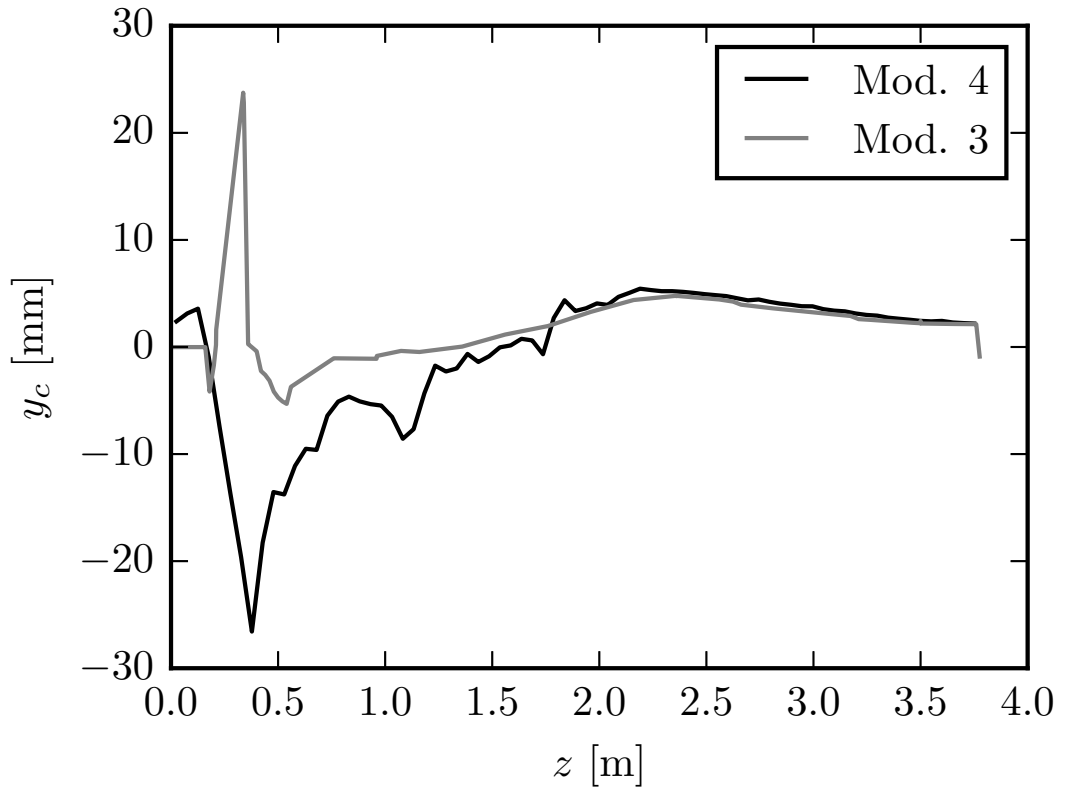
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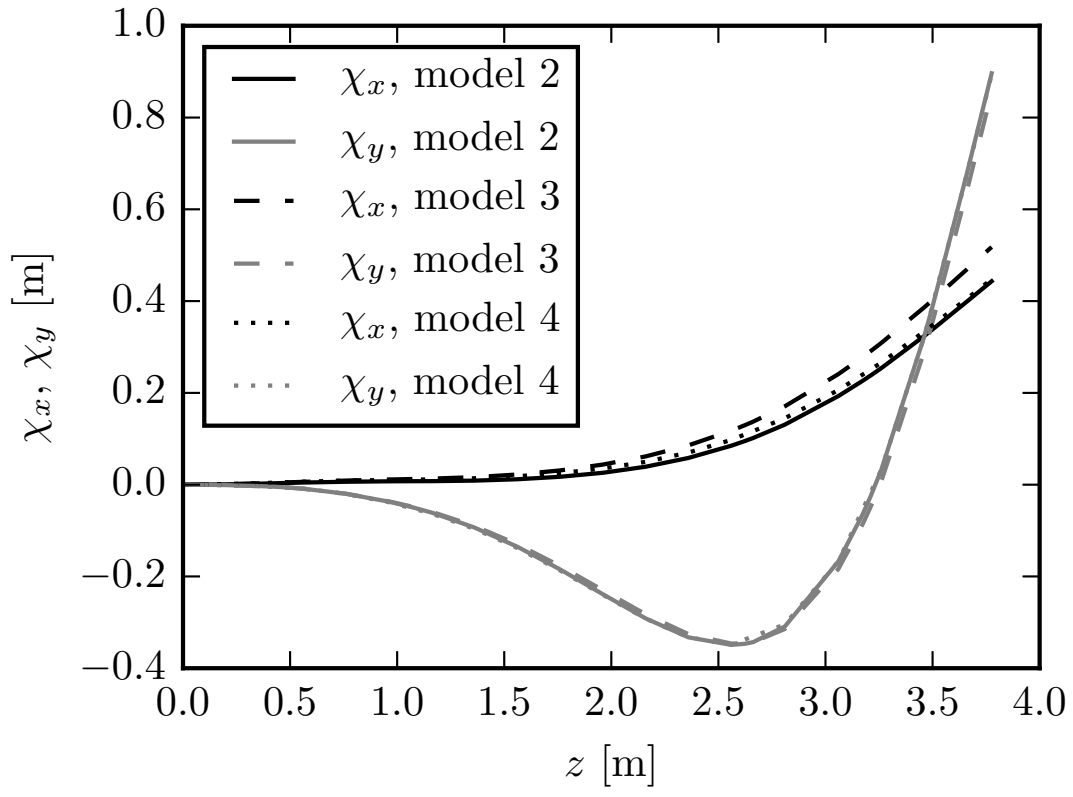
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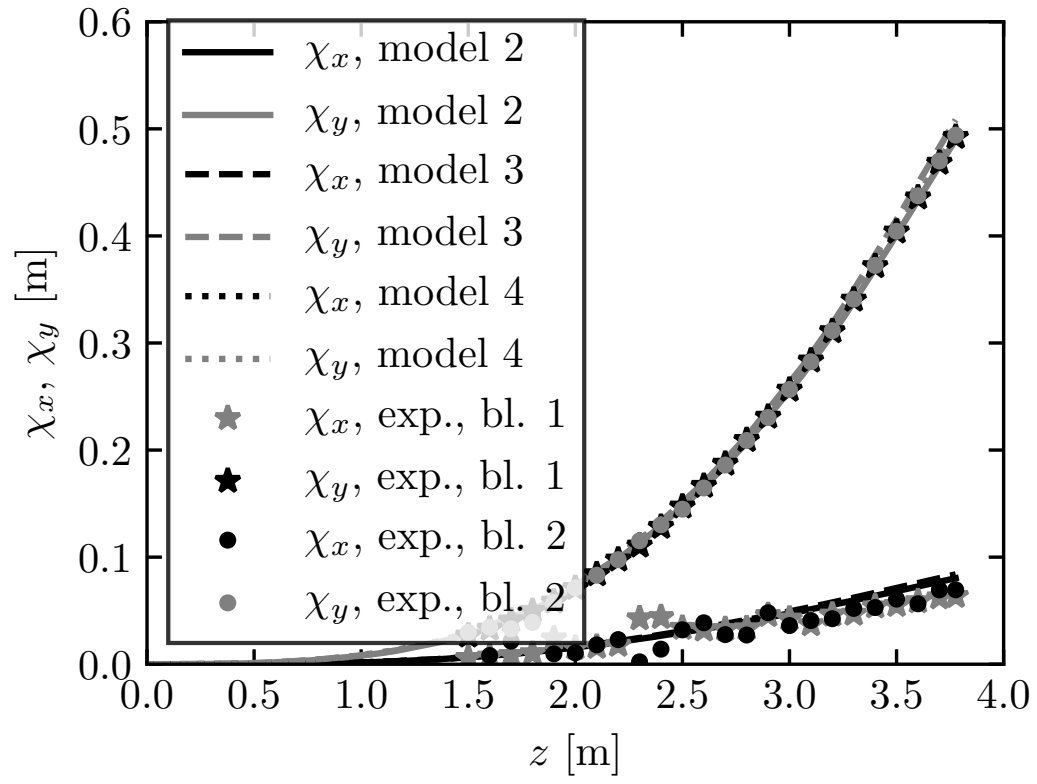
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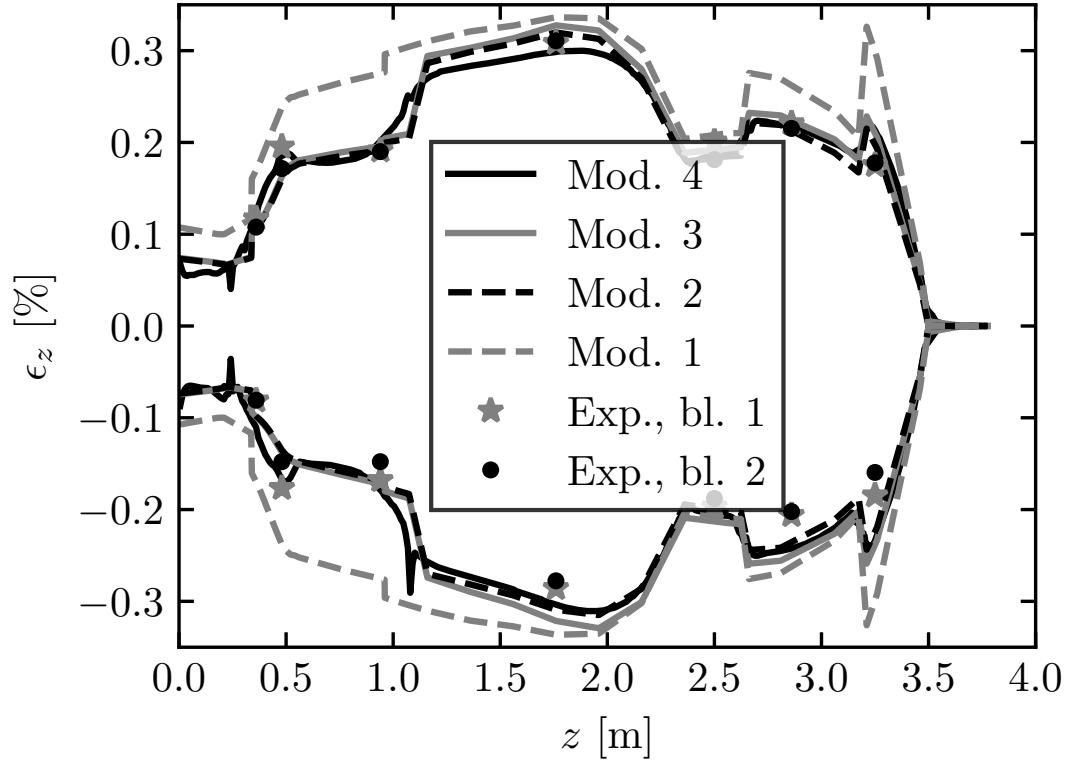
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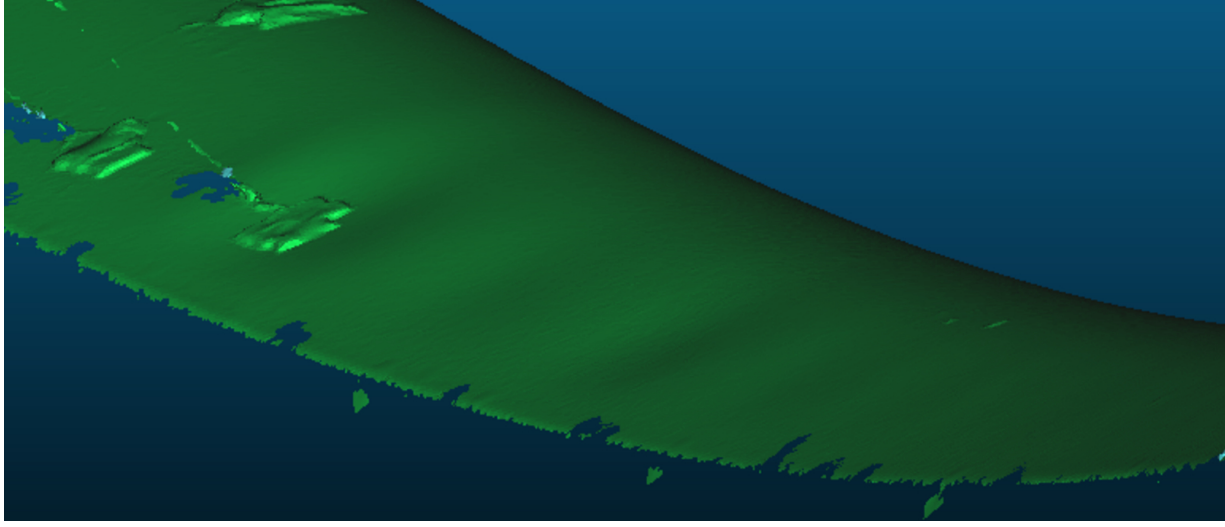
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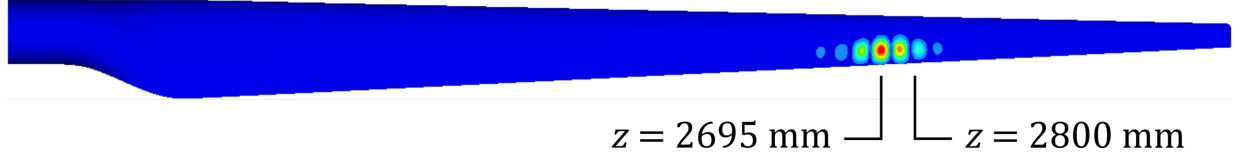
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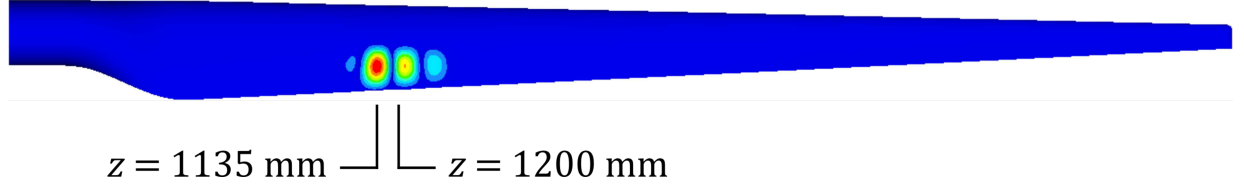
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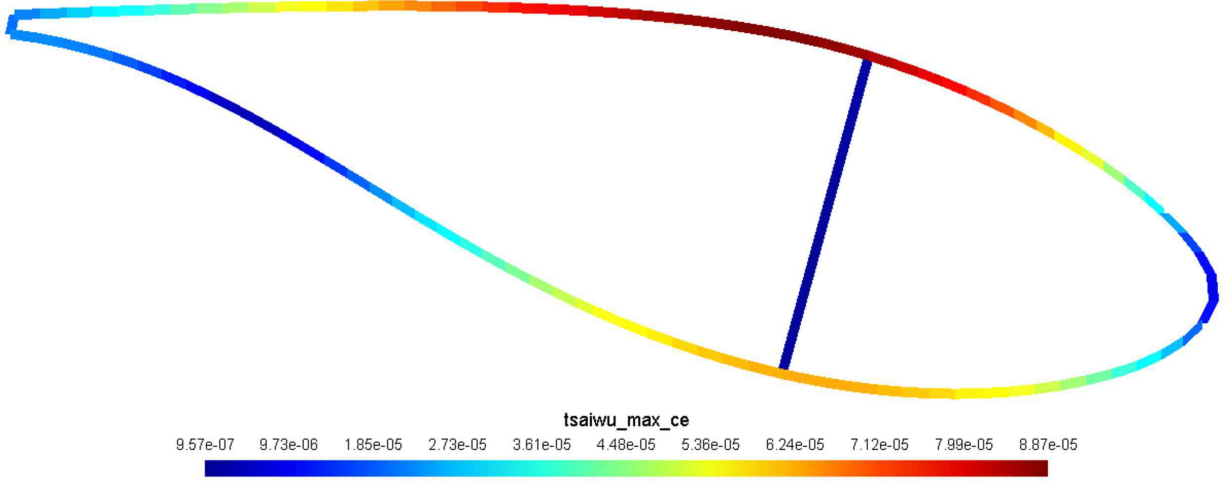
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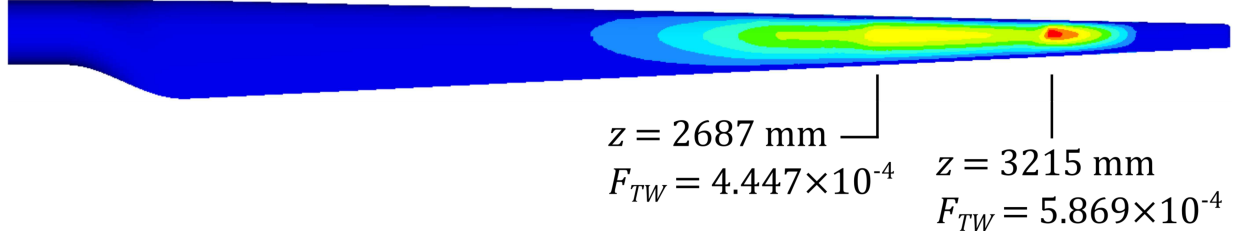
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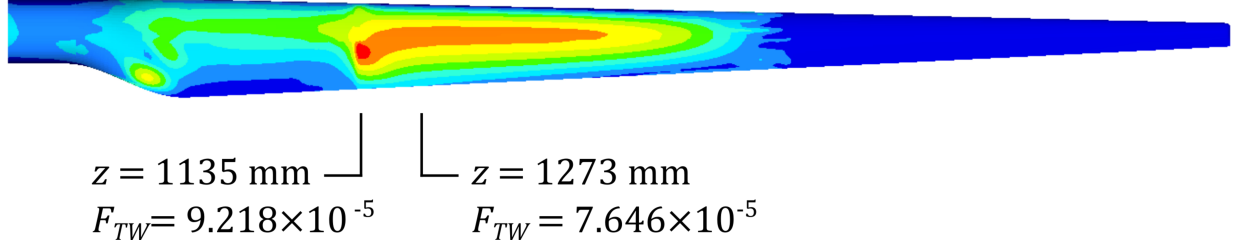
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