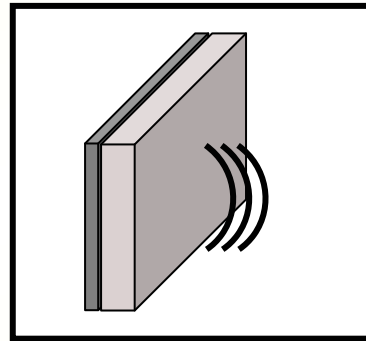


158th Meeting of the Acoustical Society of America
San Antonio, Texas

Acoustic radiation of a vibrating wall covered by a porous layer

Transfer impedance concept and effect of compression



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LAUM



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Université du Maine, Le Mans, France*

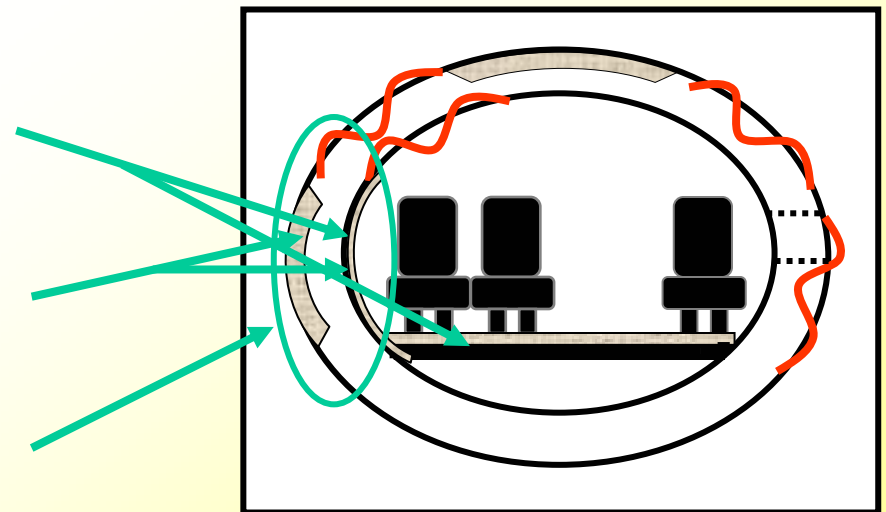
29 october 2009

Context

- Porous materials used in industrial applications
(*automotive, aeronautics,...*)

➔ for noise reduction

- ✓ sound absorption
(*trim panel, floor,...*)
- ✓ vibration damping
(*fuselage*)
- ✓ sound insulation
(*fuselage*)



Introduction

Part I

Part II

Part III

Conclusion

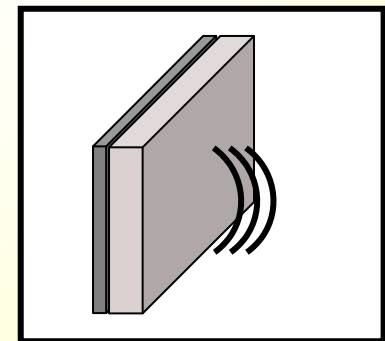
Context

➔ Porous material attached to a vibrating structure

Influence of a porous layer on the acoustic radiation of a plate ?

Method

- Analytical model using transfert impedance concept
- Experimental validation



➔ **Porous layer impedance applied to a moving wall:
Application to the radiation of a covered piston,**
Doutres, Dauchez, Genevaux, J. Acoust. Soc. Am. **121**(1), 2007

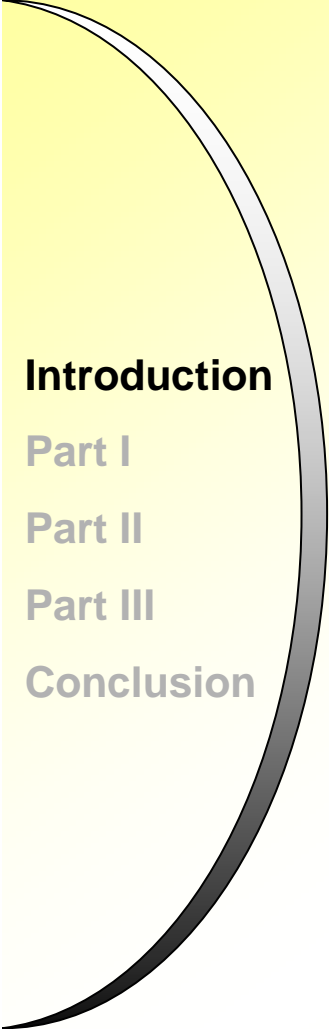
Introduction

Part I

Part II

Part III

Conclusion



Introduction

Part I

Part II

Part III

Conclusion

Introduction

1. Transfert impedance concept

2. Acoustic radiation efficiency

2.1. Infinite plate

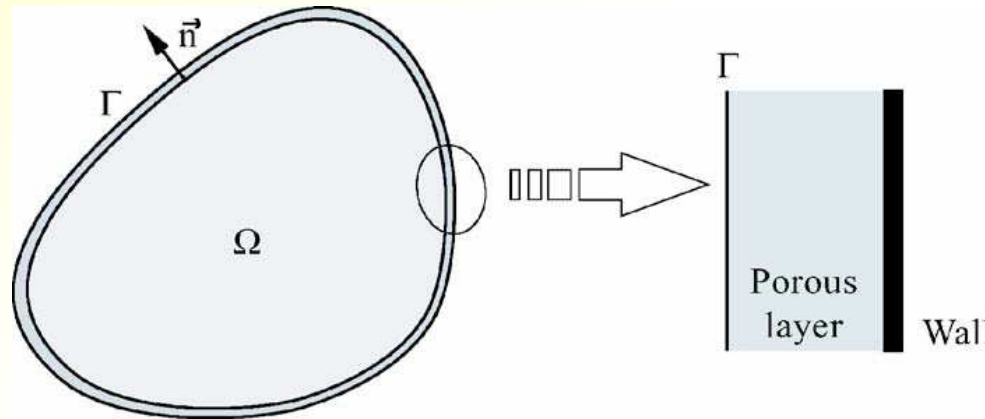
2.2. Flat piston

2.3. Circular plate

3. Application to multilayer

Conclusion

Problem to be solved



Introduction

Part I

Part II

Part III

Conclusion

$$\text{in } \Omega, \quad \nabla^2 p + k^2 p = 0$$

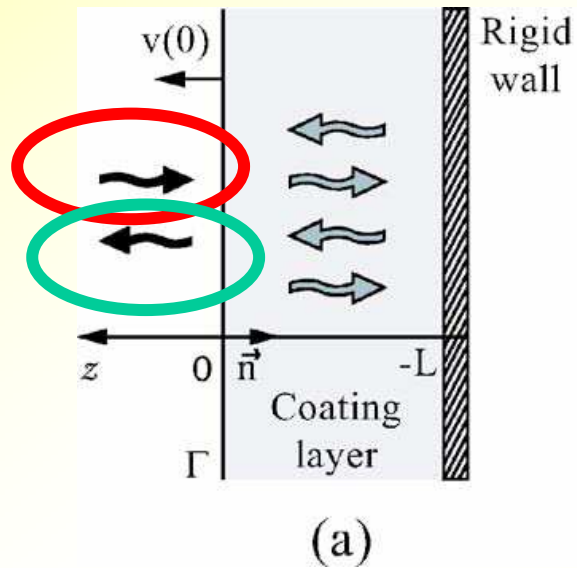
$$\text{at } \Gamma, \quad A p + B \frac{\partial p}{\partial n} = h(r)$$

Source at boundary

Dirichlet
(imposed pressure)

Neumann
(imposed velocity)

Problem divided into 2 cases :



- a) Acoustic excitation : amplitude of reflected wave ?

Introduction

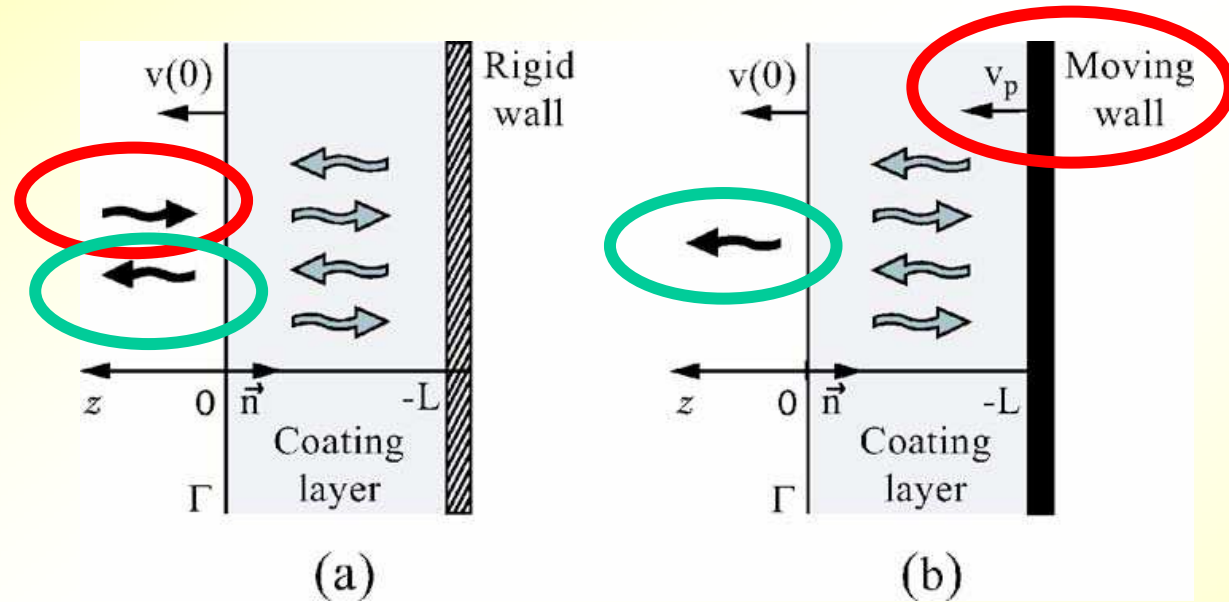
Part I

Part II

Part III

Conclusion

Problem divided into 2 cases :



- a) Acoustic excitation : amplitude of reflected wave ?
- b) Vibratory excitation : amplitude of transmitted wave ?

Introduction

Part I

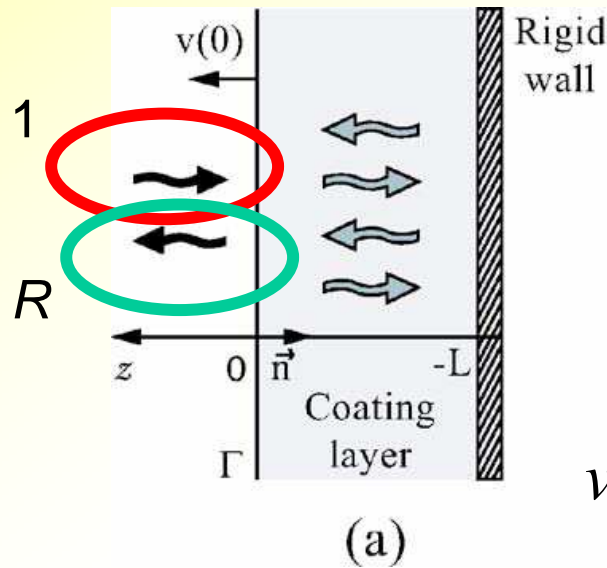
Part II

Part III

Conclusion

- Introduction
- Part I**
- Part II
- Part III
- Conclusion

Surface Impedance Z_S



Porous layer characterized by

$$Z_S = \frac{p(0)}{v(0)}$$

$$p(z) = 1 e^{jkz} + R e^{-jkz}$$

$$v(z) = \frac{1}{Z_0} (e^{jkz} - R e^{-jkz})$$

The reflected wave is function of surface impedance Z_S :

$$R = \frac{Z_S - Z_0}{Z_S + Z_0} \quad \text{with} \quad Z_0 = \rho_0 c_0$$



Z_S is measured in a Kundt tube

- Introduction
- Part I**
- Part II
- Part III
- Conclusion

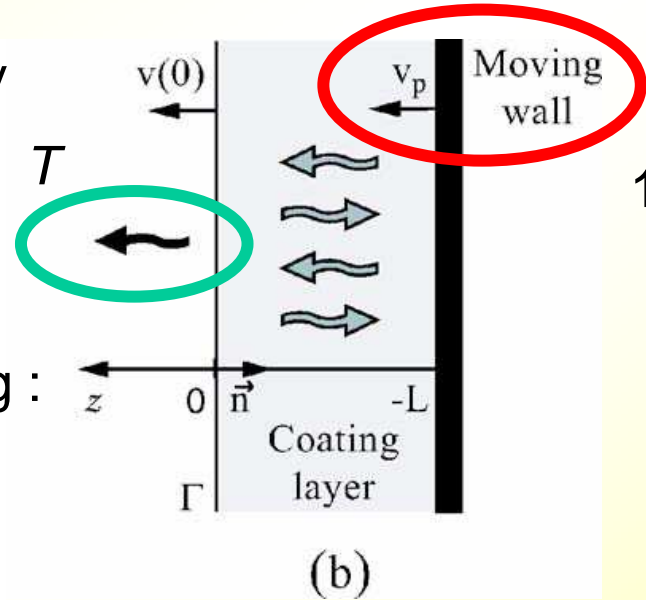
Transfert impedance Z_T

Porous layer characterized by

$$Z_T = \frac{p(0)}{v_p - v(0)}$$

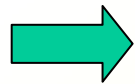
Non-porous massless coating :

$$Z_T = \frac{\text{Bulk modulus}}{j\omega \text{ Thickness}}$$



The transmitted wave is function of transfert impedance Z_T :

$$p(z) = T e^{-jkz} \quad \text{gives} \quad T = \frac{Z_T}{Z_T + Z_0} Z_0 v_p$$



Can Z_T be measured in a Kundt tube ?

Is Z_T equivalent to Z_s ?

A simple coating : Spring

Introduction

Part I

Part II

Part III

Conclusion

Static law:

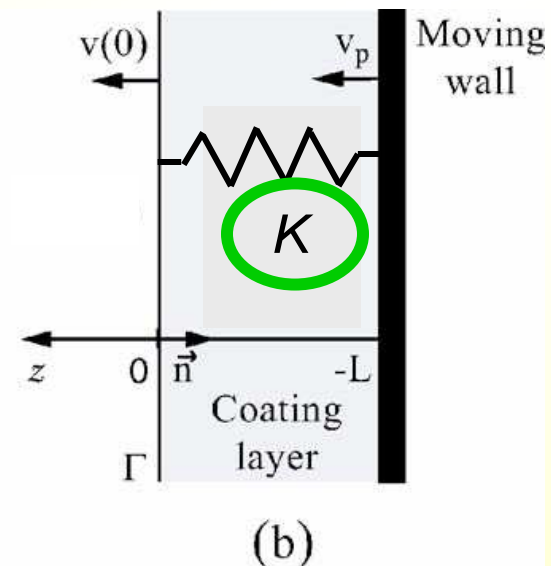
$$p(0) - \frac{K}{j\omega} (v - v_p) = 0$$

Case a: $v_p = 0$

$$Z_S = \frac{p(0)}{v} = K / j\omega$$

Case b: $Z_T = \frac{p(0)}{v - v_p} = K / j\omega$

→ $Z_T = Z_S = \frac{K}{j\omega} = \frac{\text{Bulk modulus}}{j\omega L}$



Elastic layer

Adding a layer defined by its mass

Introduction

Part I

Part II

Part III

Conclusion

Dynamic law:

$$j\omega M_s v = p(0) - \frac{K}{j\omega} (v - v_p)$$

Case a: $v_p = 0$

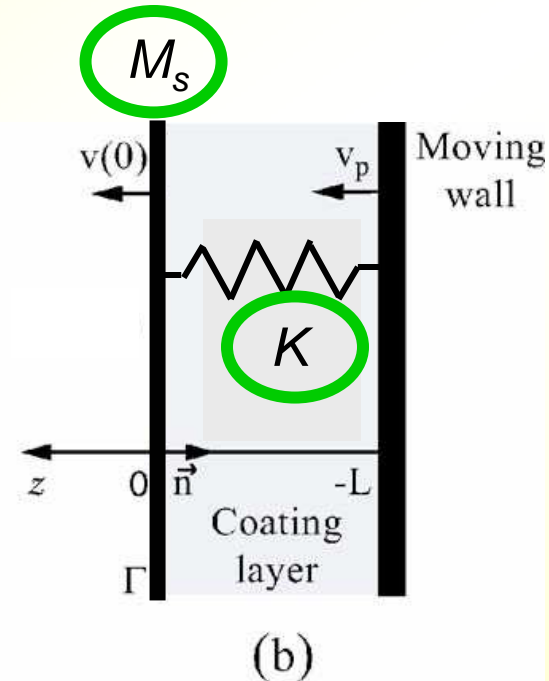
$$Z_S = \frac{p(0)}{v} = j\omega M_s + K / j\omega$$

Case b: $\frac{p(0)}{v} = Z_0 \Rightarrow Z_T = \frac{K}{j\omega Z_0 - j\omega M_s} Z_0$

➔ $Z_T \neq Z_S$ excepted when $\omega \rightarrow 0$:

$$Z_T = Z_S = \frac{K}{j\omega} = \frac{\text{Bulk modulus}}{j\omega L}$$

Elastic layer at low frequency



Monophasic continuous layer

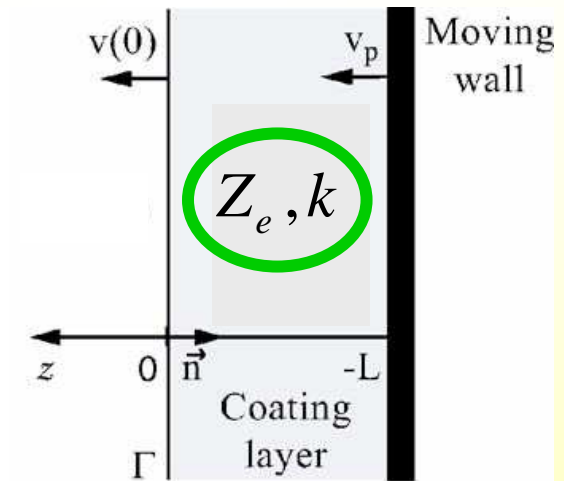
Transfert matrix:

$$\begin{pmatrix} p(0) \\ v(0) \end{pmatrix} = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \begin{pmatrix} p(-L) \\ v(-L) \end{pmatrix}$$

$$\begin{bmatrix} a & b \\ c & d \end{bmatrix} = \begin{bmatrix} \cos kd & jZ_e \sin kd \\ j/Z_e \sin kd & \cos kd \end{bmatrix}$$

Case a: $v_p = 0 \Rightarrow Z_S = \frac{a}{c}$

Case b: $\frac{p(0)}{v} = Z_0 \Rightarrow Z_T = \frac{Z_0}{a - 1 - cZ_0}$



(b)

➔ $Z_T \neq Z_S$ excepted when $\omega \rightarrow 0$: $a \rightarrow 1$ and $c \rightarrow j \frac{kd}{Z_e}$

$$Z_T = Z_S = j \frac{Z_e}{kd} = \frac{\text{Bulk modulus}}{j\omega L}$$

Monophasic layer at low frequency

Introduction

Part I

Part II

Part III

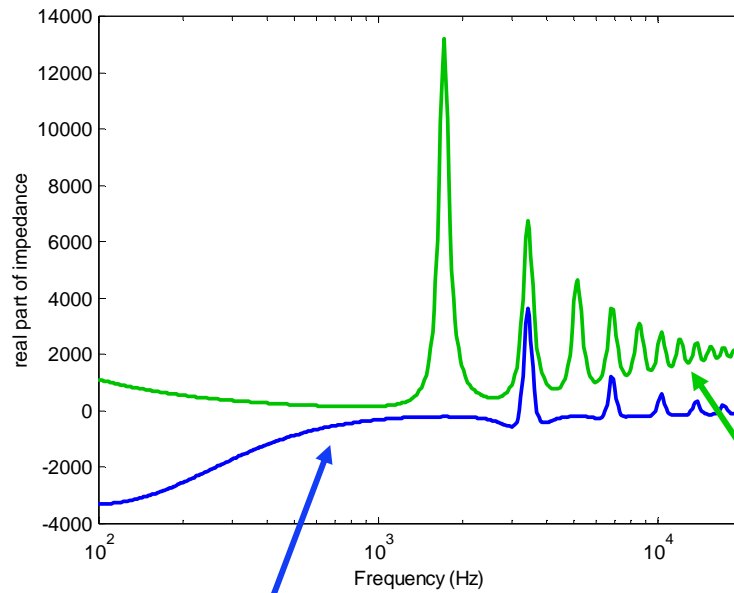
Conclusion

- Introduction
- Part I**
- Part II
- Part III
- Conclusion

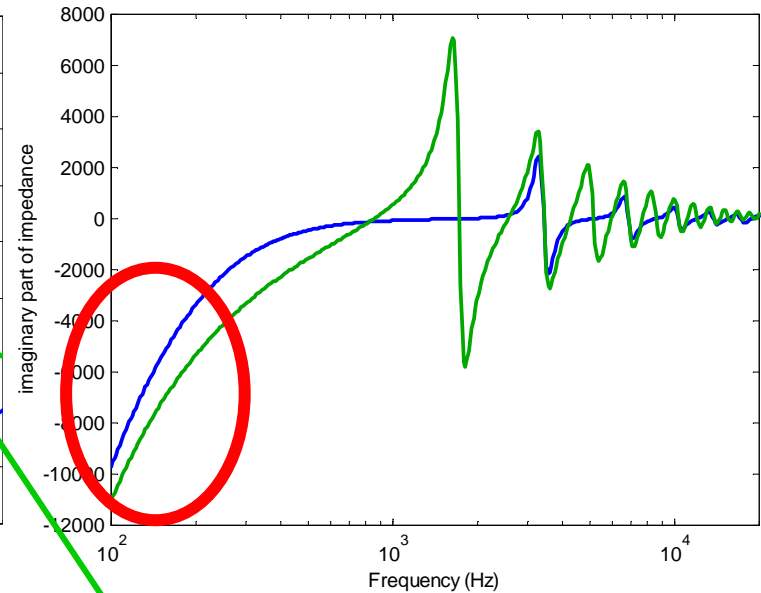
Impedance curves

Thickness	20 mm
Bulk modulus	140(1+j0.1) kPa
Density	30 kg.m ⁻³

Real part



Imaginary part



Z_T transfert impedance **≠** **Z_S surface impedance**

Poroelastic layer: Surface impedance Z_s

Biot-Allard theory in 1D in the porous layer :
2 waves travelling in each direction

Two waves in fluid medium

Boundary conditions : Continuity of Stress and displacement

at $x=0$ $v^s(0) = v^f(0) = 0$

at $x=d$ $v(d) = (1 - \phi)v^s(0) + \phi v^f(0)$

$\sigma^f(d) = -\phi p(d)$ and $\sigma^s(d) = -(1 - \phi)p(d)$

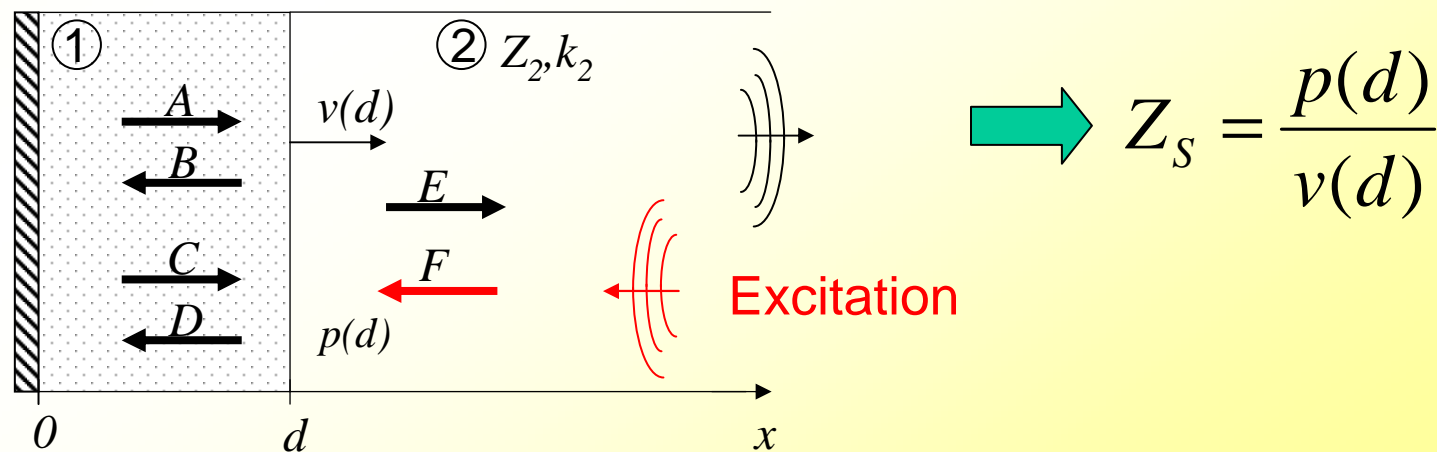
Introduction

Part I

Part II

Part III

Conclusion



Poroelastic layer: Transfert impedance Z_T

Biot-Allard theory in 1D in the porous layer :
2 waves travelling in each direction

One waves in fluid medium $\Rightarrow \frac{p(d)}{v(d)} = Z_0$

Boundary conditions : Continuity of Stress and displacement

at $x=0$ $v^s(0) = v^f(0) = v_p$

at $x=d$ $v(d) = (1 - \phi)v^s(0) + \phi v^f(0)$

$\sigma^f(d) = -\phi p(d)$ and $\sigma^s(d) = -(1 - \phi)p(d)$

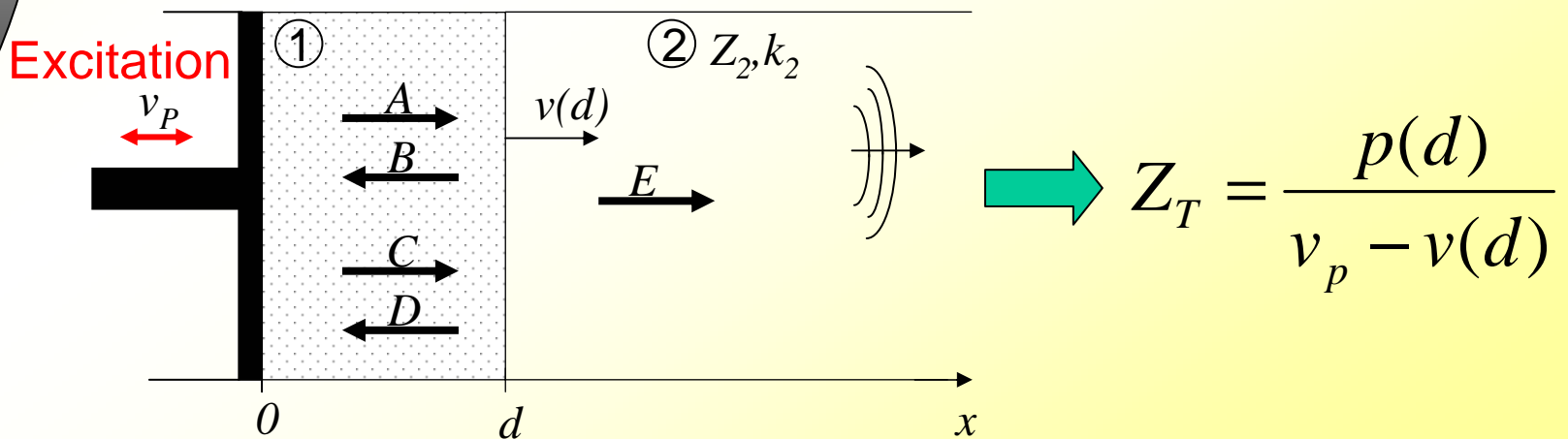
Introduction

Part I

Part II

Part III

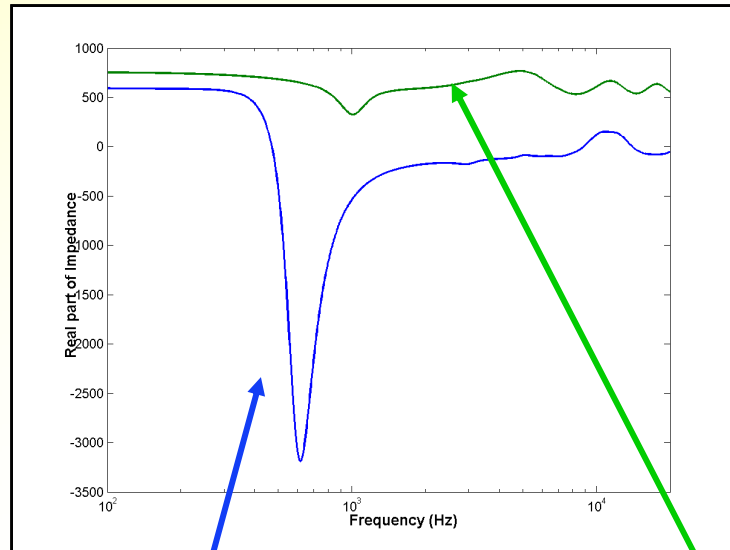
Conclusion



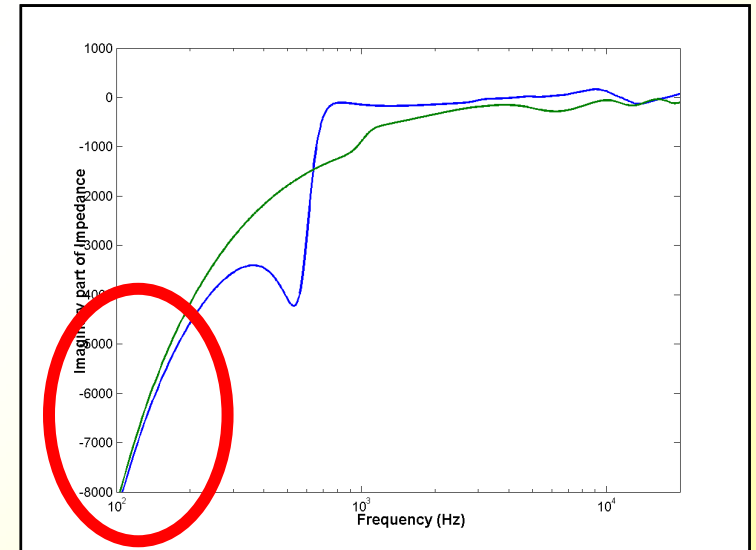
Thickness (<i>mm</i>)	20.17
Air flow resistivity (<i>N sm⁻⁴</i>)	75000
Skeleton Young's Modulus at 5Hz (<i>Pa</i>)	285000
Skeleton density (<i>kgm⁻³</i>)	59

Impedance curves

Real part



Imaginary part



Z_T transfert impedance

\neq

Z_S surface impedance

Introduction

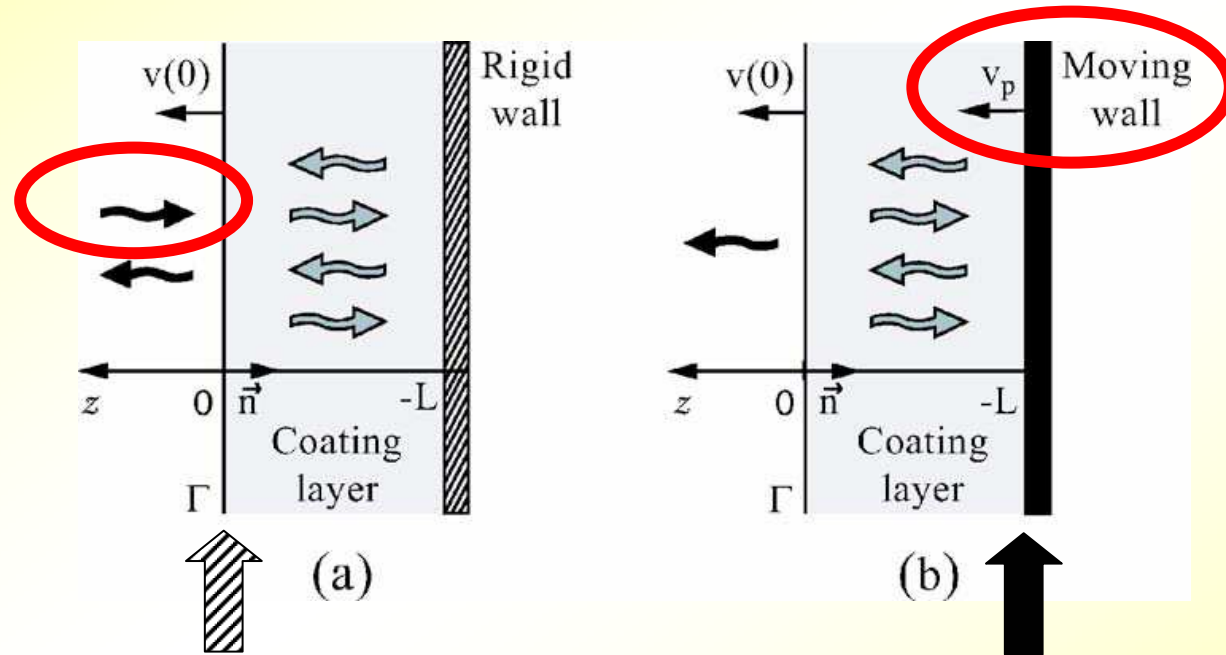
Part I

Part II

Part III

Conclusion

Boundary conditions at excitation interface



Fluid-porous interface :

$$v(0) / j\omega = (1 - \phi)u^s + \phi u^f$$

Wall-porous interface :

$$v_p / j\omega = u^s = u^f$$

\Rightarrow The skeleton is much more excited in case b)

Introduction

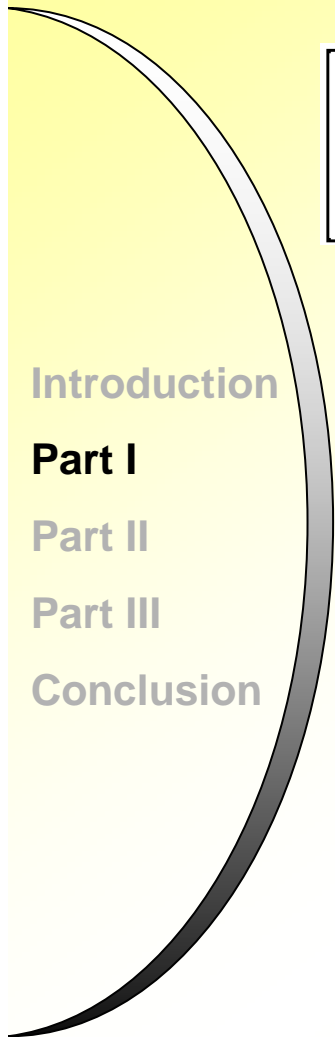
Part I

Part II

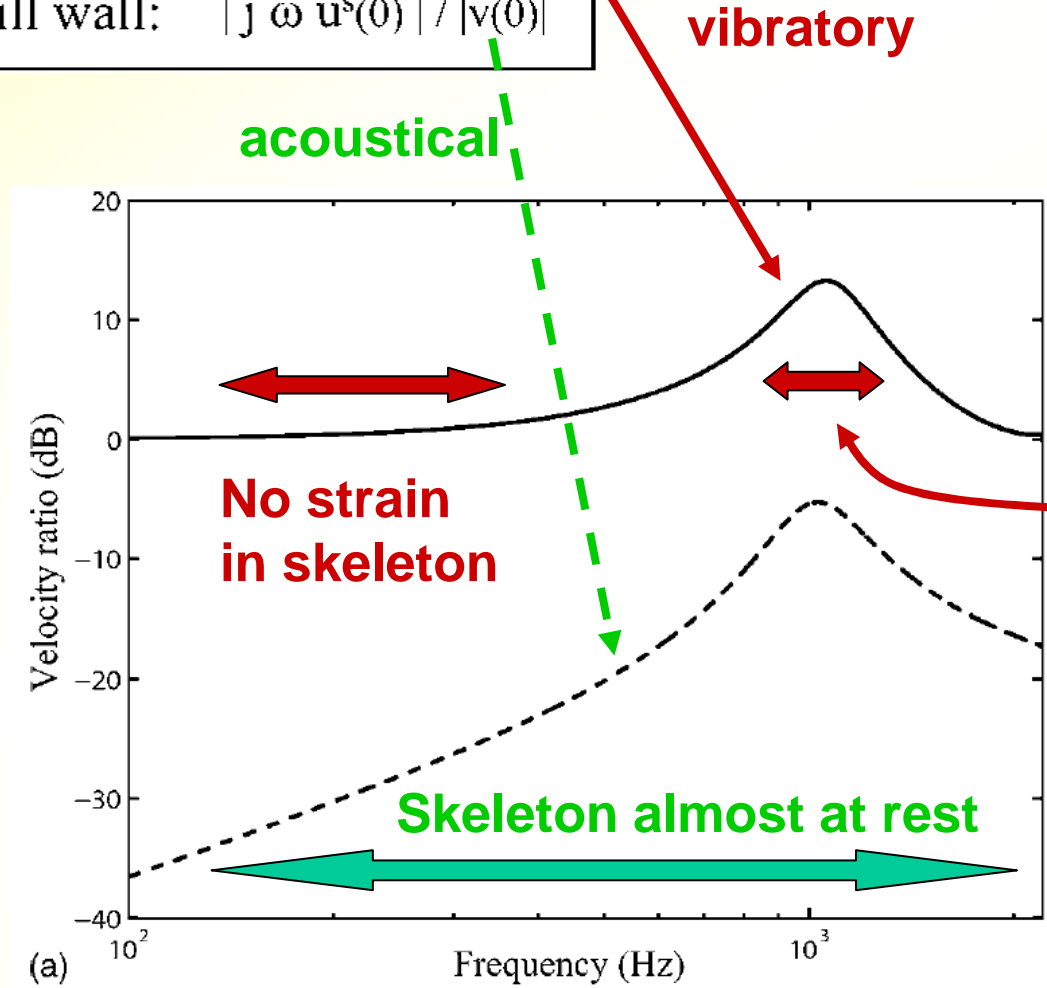
Part III

Conclusion

Relative skeleton velocity for both excitation at fluid-porous interface



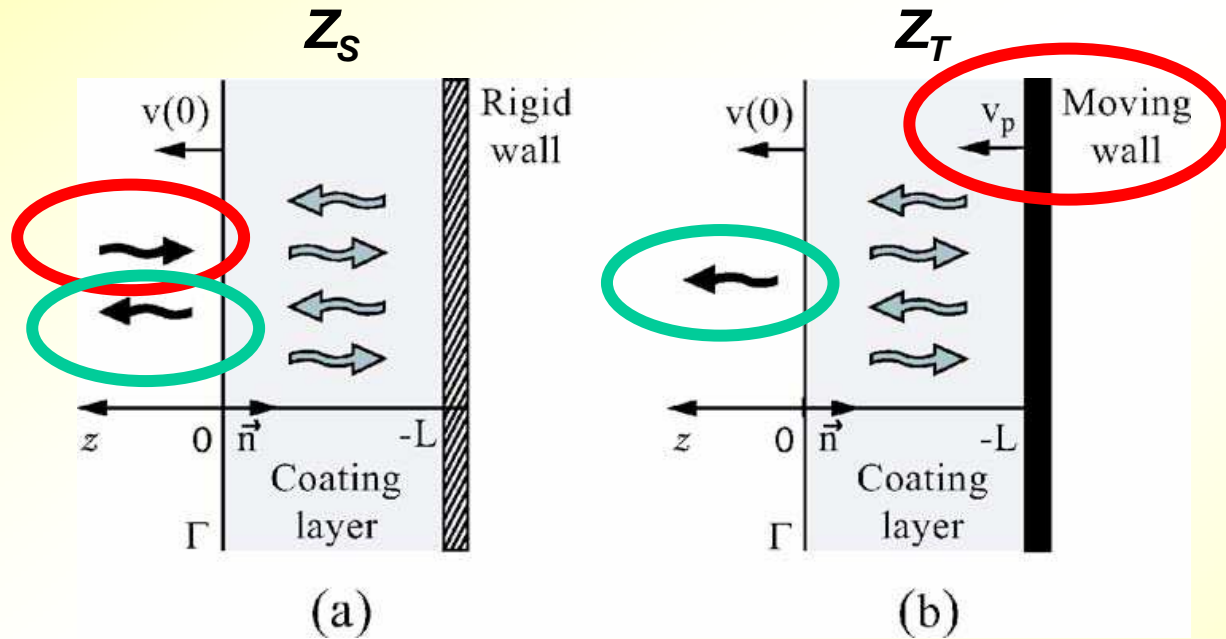
— Moving wall: $|j \omega u^s(0)| / |v_p|$
 - - Still wall: $|j \omega u^s(0)| / |v(0)|$



Strong influence of frame borne wave

(a)

Conclusion of Part I



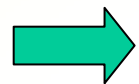
Introduction

Part I

Part II

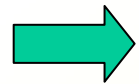
Part III

Conclusion



Excepted at low frequency,

$$Z_T \neq Z_S$$



Z_T can not be measured in a Kundt tube

1 impedance for each problem

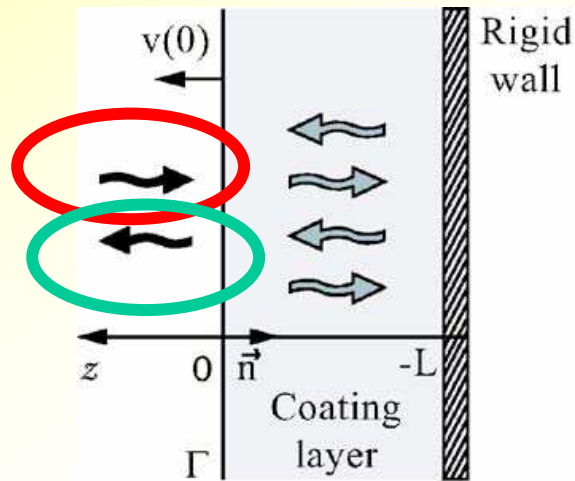
Introduction

Part I

Part II

Part III

Conclusion



(a)

$$(a) \quad jk \frac{Z_0}{Z_S} p + \frac{\partial p}{\partial z} = 0$$

$$(b) \quad jk \frac{Z_0}{Z_T} p - \frac{\partial p}{\partial z} = j\omega\rho_0 v_p$$



Introduction

Part I

Part II

Part III

Conclusion

Introduction

1. Transfert impedance concept

2. Acoustic radiation efficiency

2.1. Infinite plate

2.2. Flat piston

2.3. Circular plate

3. Application to multilayer

Conclusion

2. Acoustic radiation efficiency

Introduction

Part I

Part II

Part III

Conclusion

$$\sigma_R = \frac{\Pi_a}{\Pi_v}$$

Radiated acoustic
power

Injected vibratory
power

$$\Pi_a = \iint_S \langle I \rangle r^2 \sin \vartheta d\vartheta d\varphi$$

with:

$$\langle I \rangle = \frac{|p|^2}{2\rho_0 c_0}$$

$$\Pi_v = \rho_0 c_0 \iint_S \frac{|w(\vec{r})|^2}{2} dS$$

2.1. Acoustic radiation efficiency

Infinite plate at normal incidence (1D) :
$$\sigma_R = \frac{|T|^2}{Z_0 v_p^2} = \left| \frac{Z_T}{Z_T + Z_0} \right|^2$$

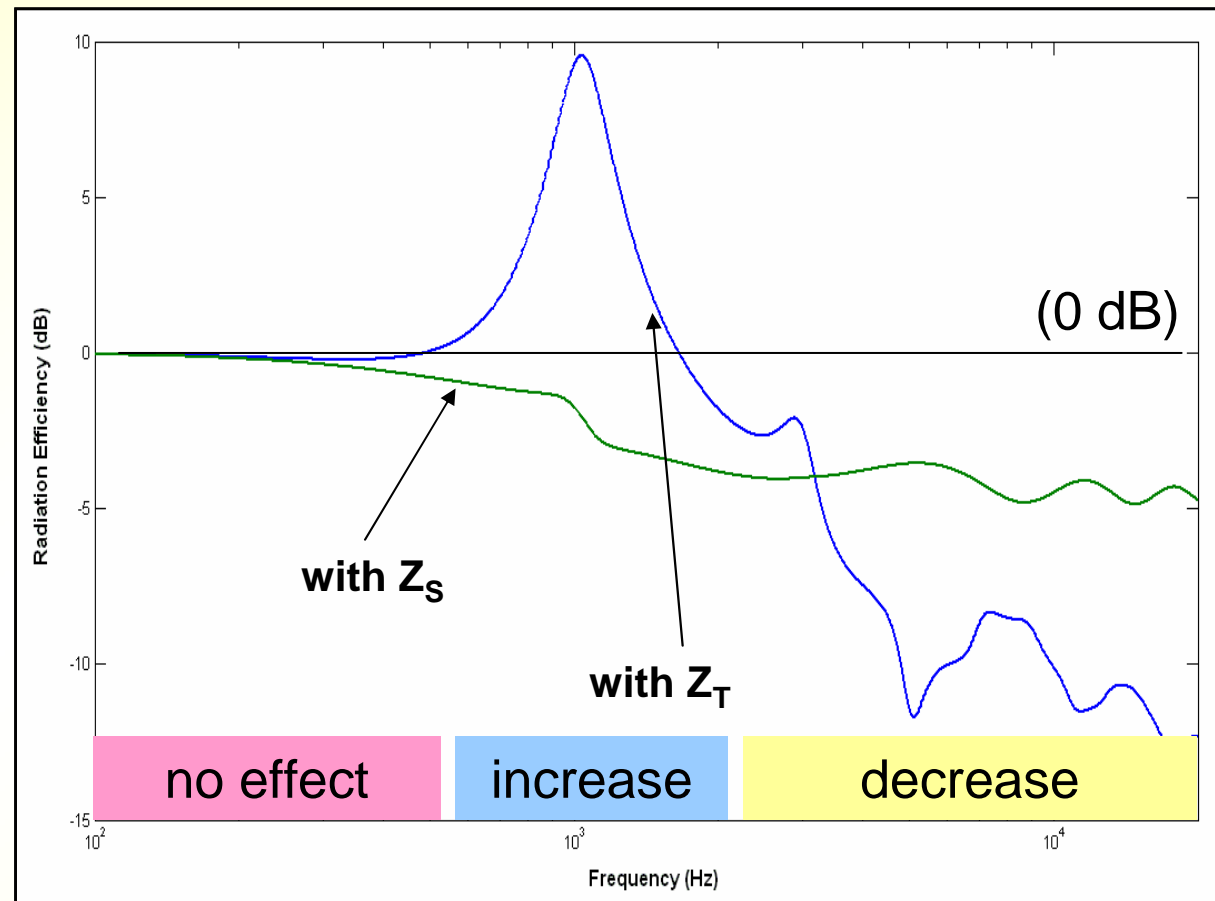
Introduction

Part I

Part II

Part III

Conclusion



2.2. Acoustic radiation of a flat piston in semi-infinite field

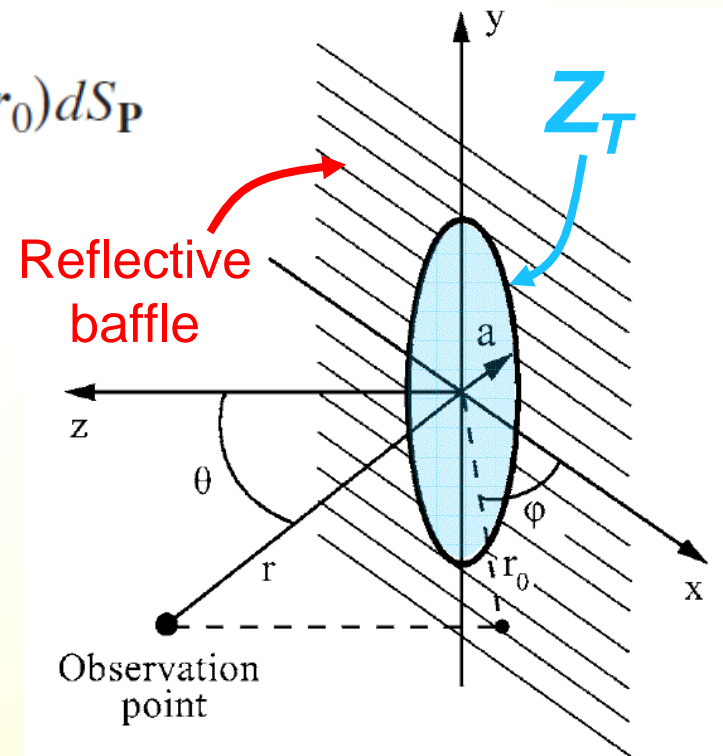
Calculation of the far field pressure: Rayleigh Integral

$$p(\mathbf{r}) = j\omega\rho_0 \frac{Z_t}{Z_t + Z_0} \int_{S_P} v_p G(\mathbf{r}, \mathbf{r}_0) dS_P$$

with $G(\mathbf{r}, \mathbf{r}_0) = \frac{e^{-jk|\mathbf{r}-\mathbf{r}_0|}}{2\pi|\mathbf{r}-\mathbf{r}_0|}$

For a flat piston of radius a
in far field :

$$p(\mathbf{r}) = \frac{Z_t}{Z_t + Z_0} jk\rho_0 c_0 v_p \frac{e^{-jkr}}{2\pi r} \pi a^2 \frac{2J_1(ka \sin \theta)}{ka \sin \theta}$$



Introduction

Part I

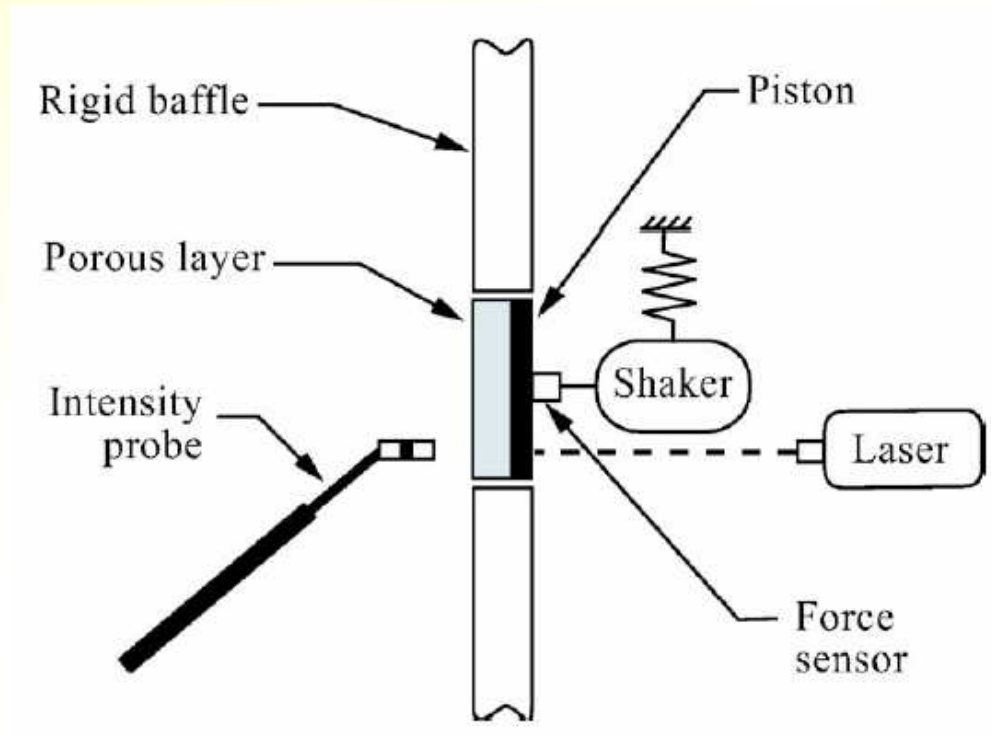
Part II

Part III

Conclusion

2.2. Acoustic radiation of a flat piston in semi-infinite field

Measurement :



2 materials :

A : polymer foam

B : soft fibrous

Introduction

Part I

Part II

Part III

Conclusion

2.2. Acoustic radiation of a flat piston in semi-infinite field

A : polymer foam

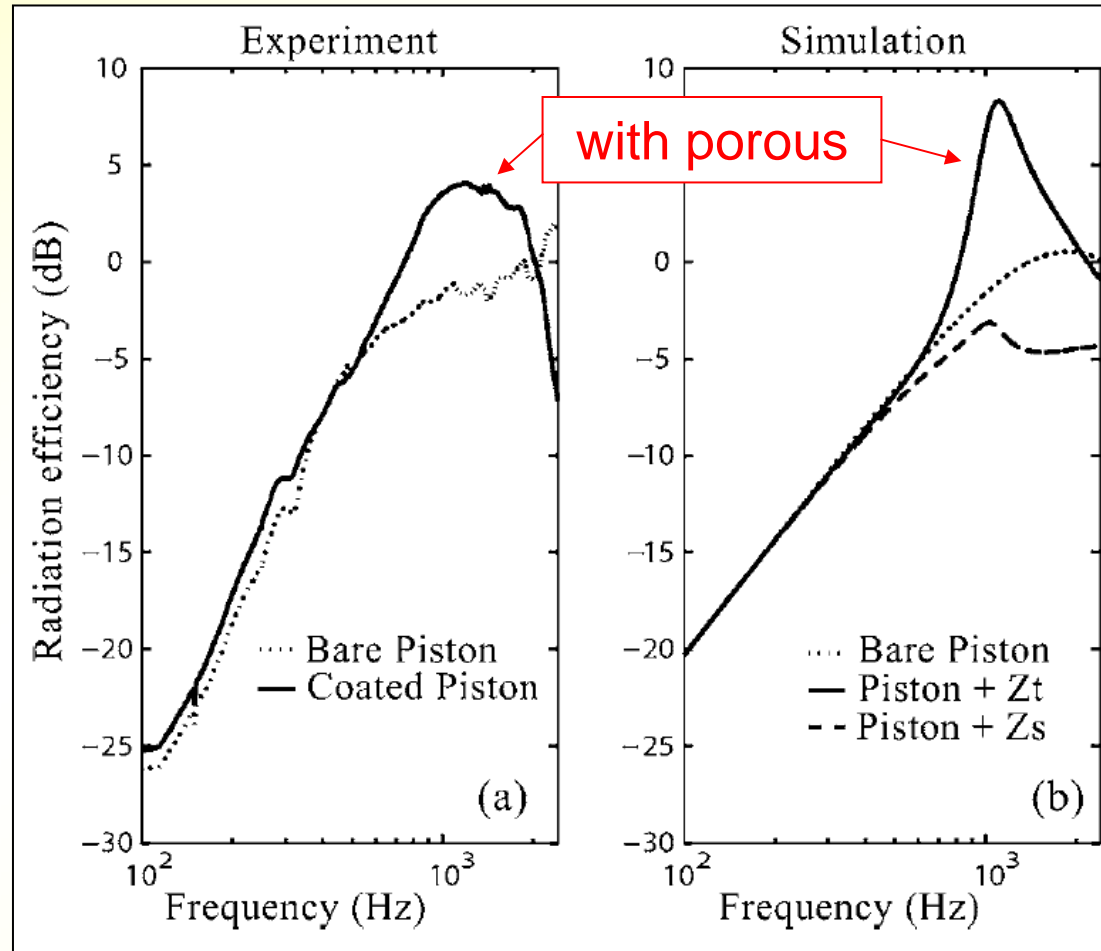
Introduction

Part I

Part II

Part III

Conclusion



2.2. Acoustic radiation of a flat piston in semi-infinite field

B : soft fibrous

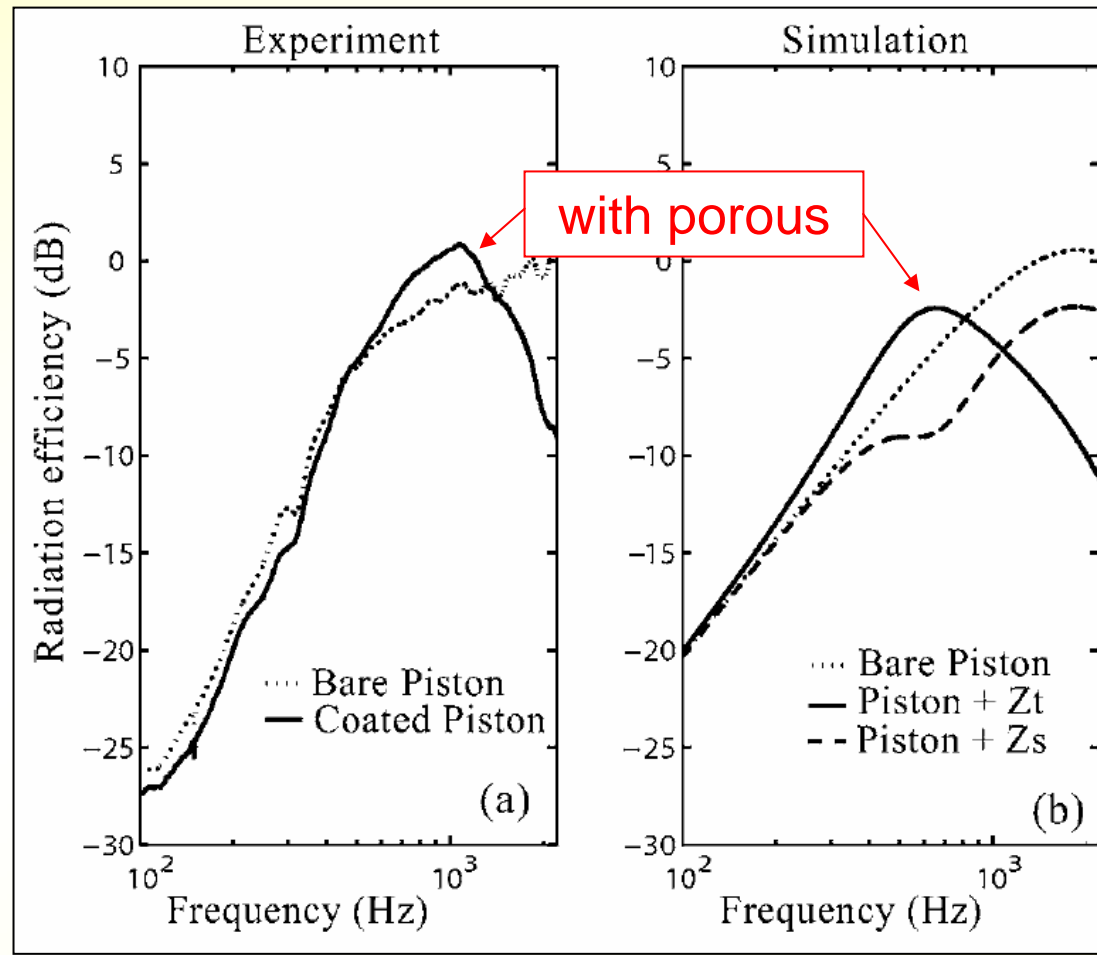
Introduction

Part I

Part II

Part III

Conclusion





Introduction

Part I

Part II

Part III

Conclusion

Introduction

1. Transfert impedance concept

2. Acoustic radiation efficiency

2.1. Infinite plate

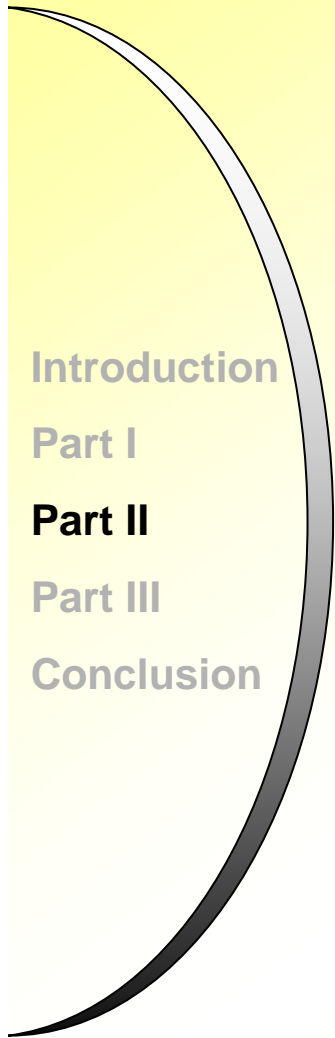
2.2. Flat piston

2.3. Circular plate

3. Application to multilayer

Conclusion

2.3. Radiation efficiency of a circular plate



Introduction

Part I

Part II

Part III

Conclusion

Plate equation



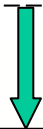
Axisymmetrical modes



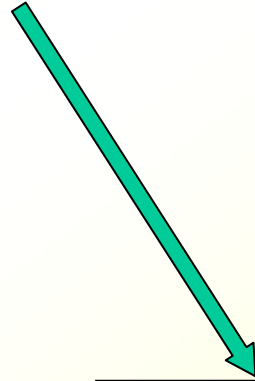
Modal synthesis



Rayleigh integration



Acoustic power



Vibratory power

$$\nabla^4 w - \frac{\rho h \omega^2}{D} w = 0$$

$$w_n(r) = J_0(\beta_{0n} r) - \frac{J_0(\beta_{0n} a)}{I_0(\beta_{0n} a)} I_0(\beta_{0n} r)$$

$$w(r) = \frac{F}{D} \sum_n \frac{w_n(r_s) w_n(r)}{(\beta_n^4 - \beta^4) \pi a^2 \Lambda_n}$$

with: $\beta_{0n}^4 = \frac{\rho h}{D} \omega_n^2$

Modal contribution to the acoustic radiation

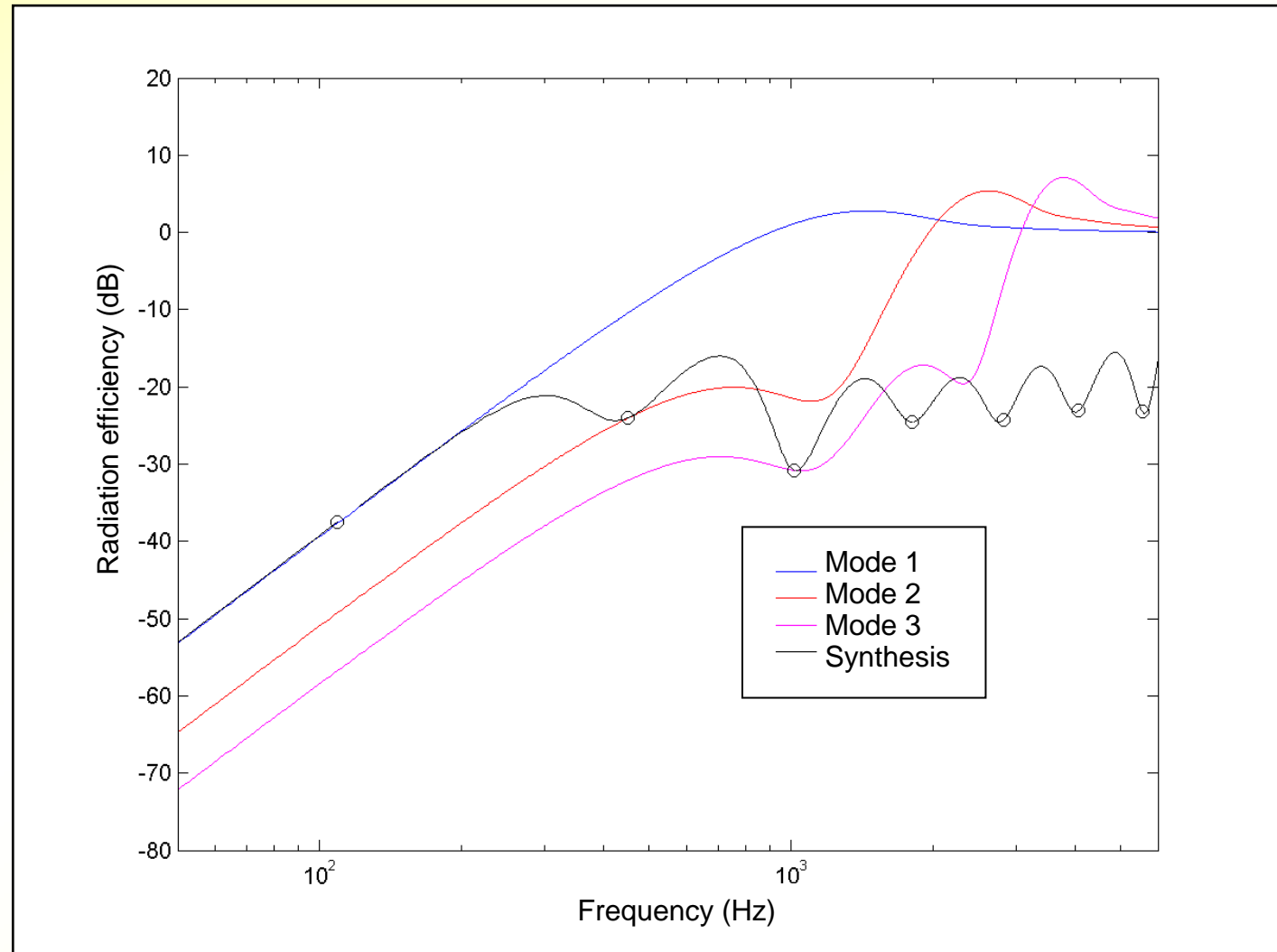
Introduction

Part I

Part II

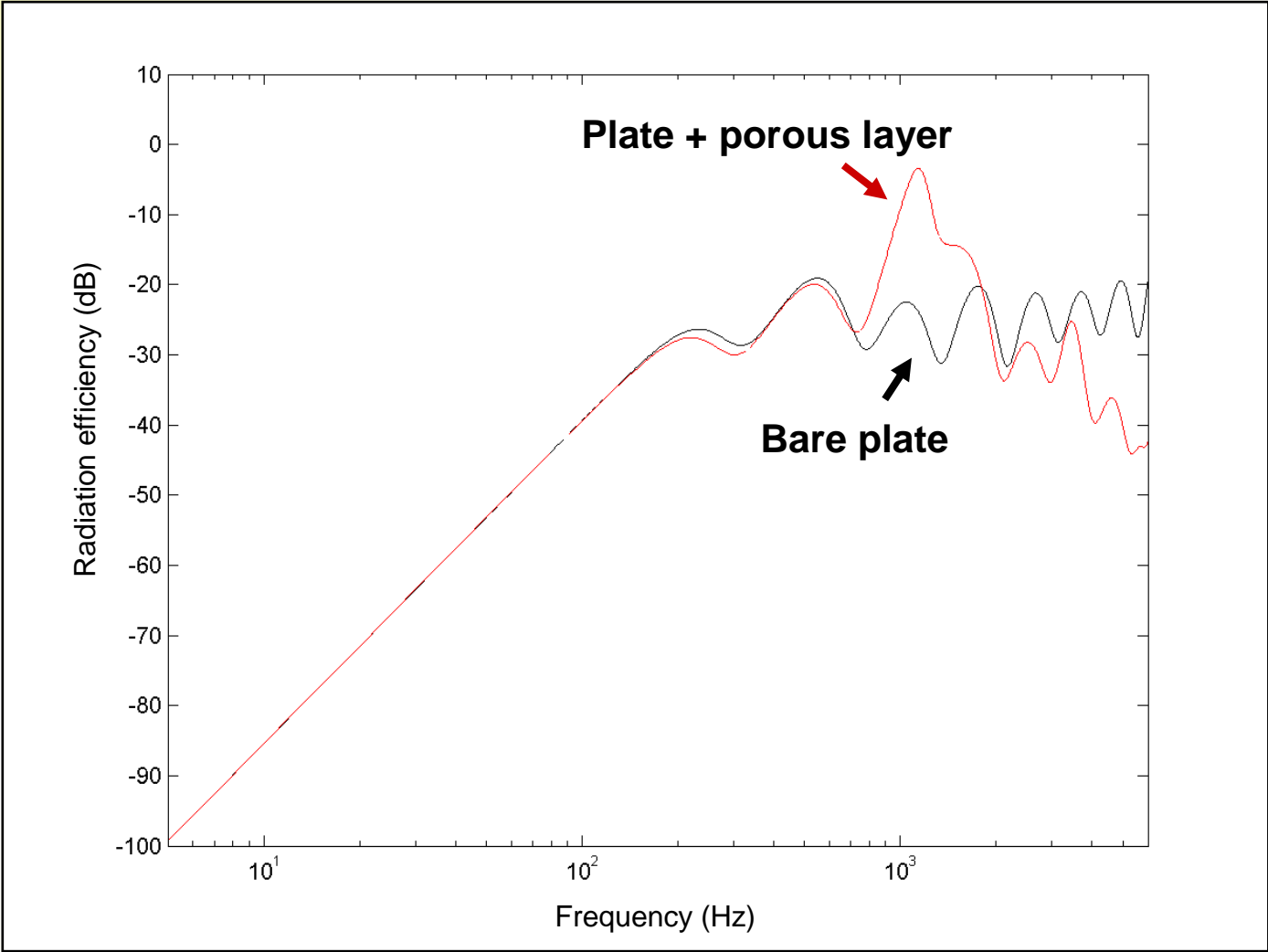
Part III

Conclusion



Influence of the porous material on the radiation of the plate

- Introduction
- Part I
- Part II**
- Part III
- Conclusion



Relative radiation efficiency

$$\sigma_{Rn} = \frac{\sigma_{R \text{ plate+porous}}}{\sigma_{R \text{ plate}}}$$

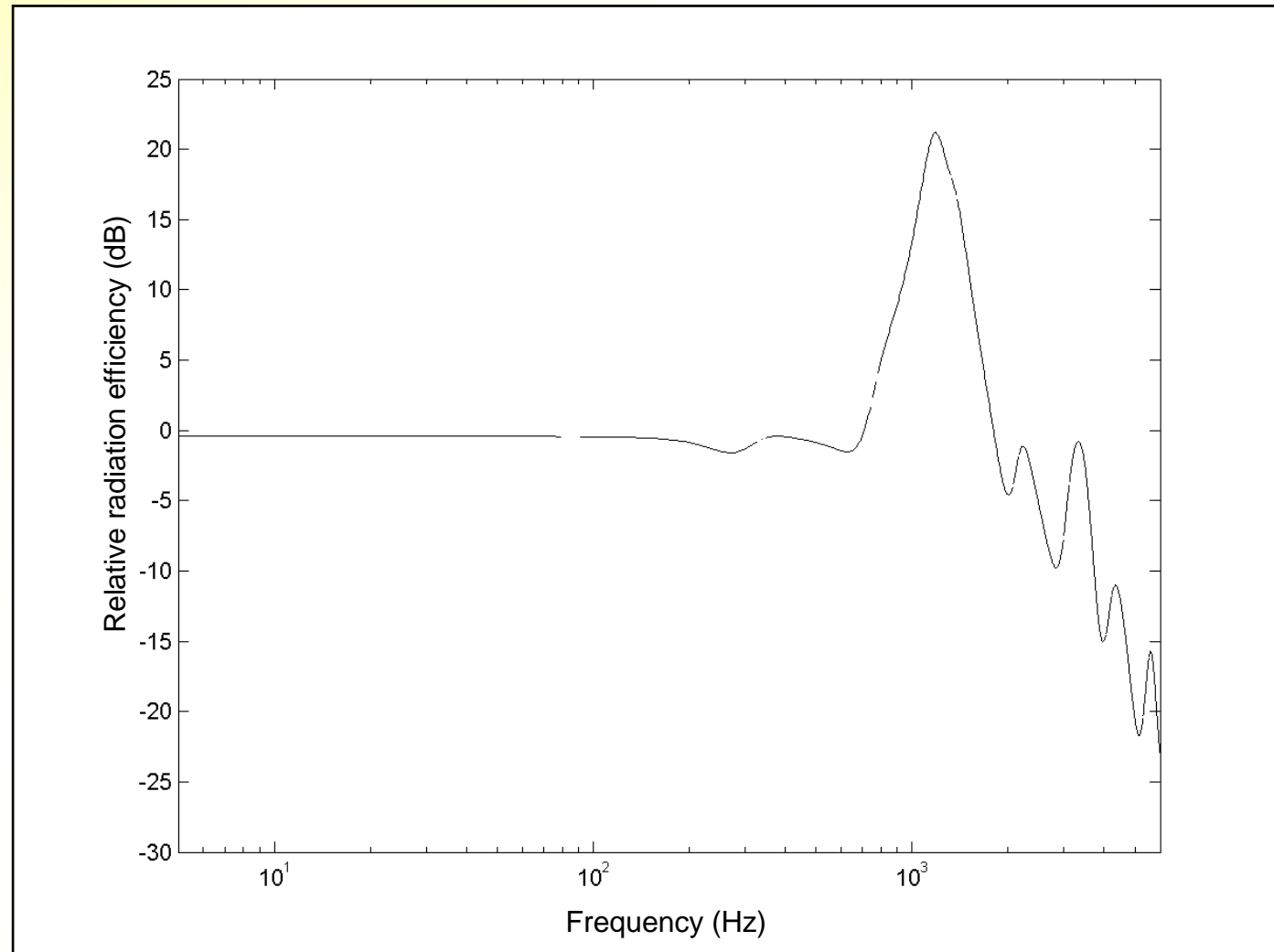
Introduction

Part I

Part II

Part III

Conclusion



Experimental validation

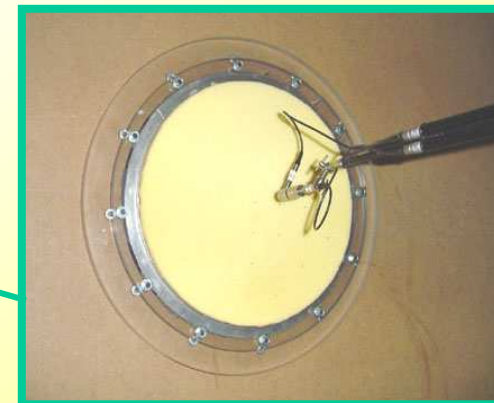
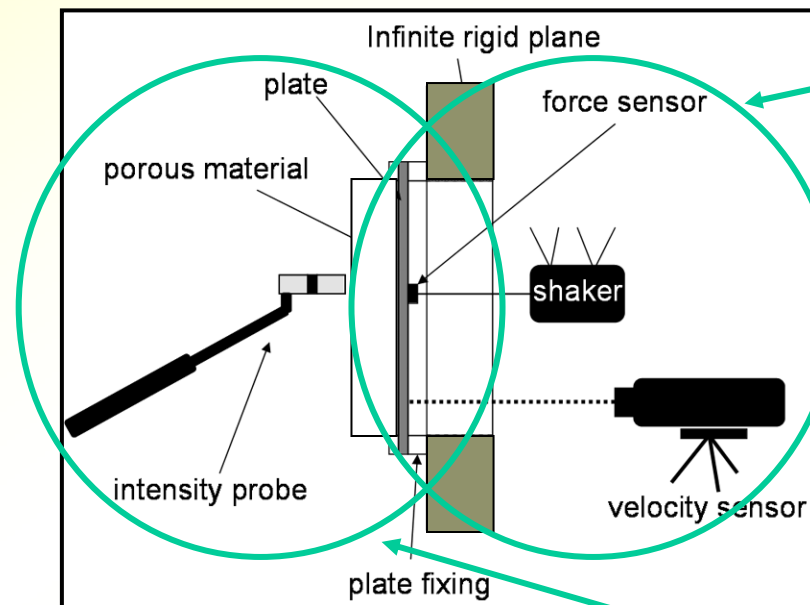
Introduction

Part I

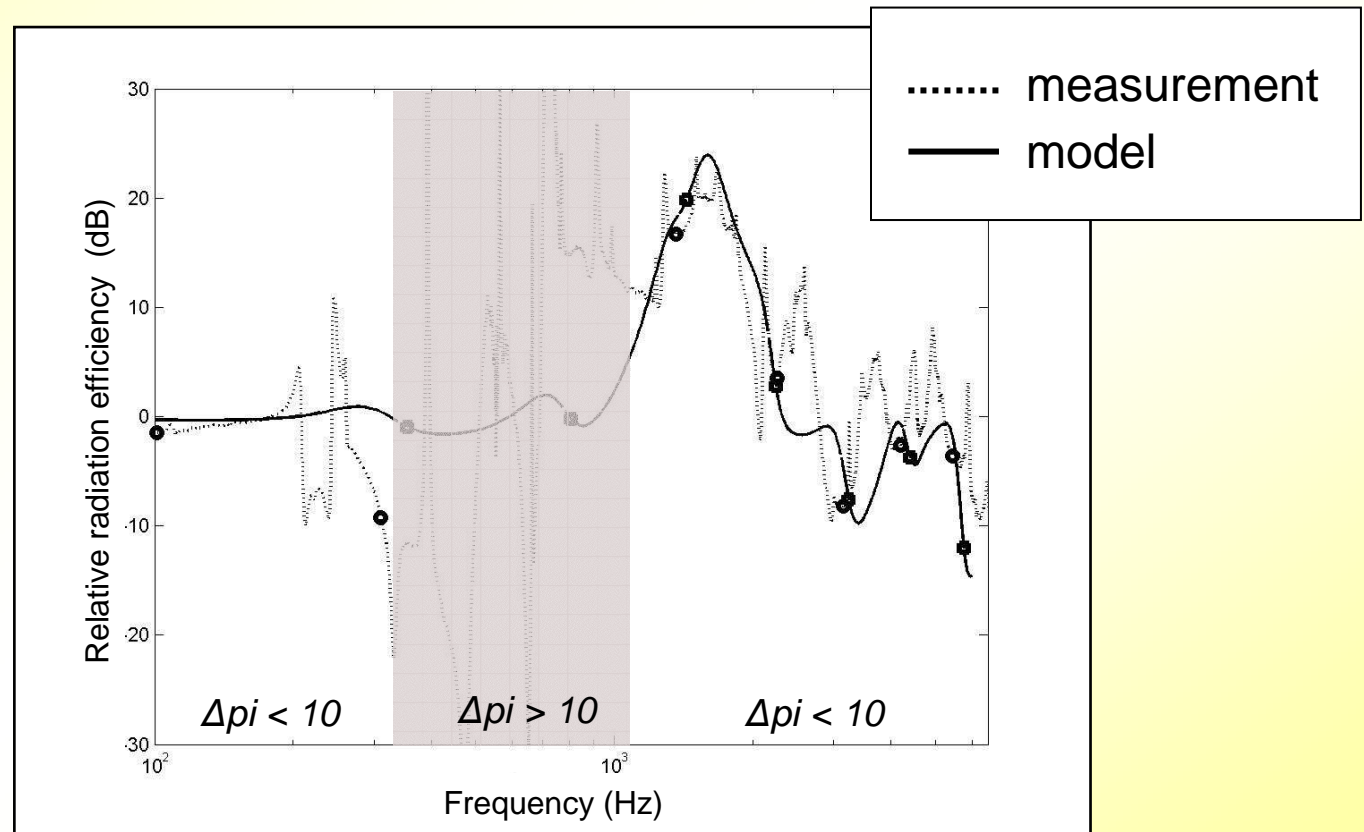
Part II

Part III

Conclusion



Results: Radiation efficiency of the covered plate



➔ **Good prediction of the acoustic radiation**

Introduction

Part I

Part II

Part III

Conclusion



Introduction

Part I

Part II

Part III

Conclusion

Introduction

1. Transfert impedance concept

2. Acoustic radiation efficiency

2.1. Infinite plate

2.2. Flat piston

2.3. Circular plate

3. Application to multilayer

Conclusion

Calculation of Z_T for a multilayer

- Use of transfert matrix method
- Maine3A software

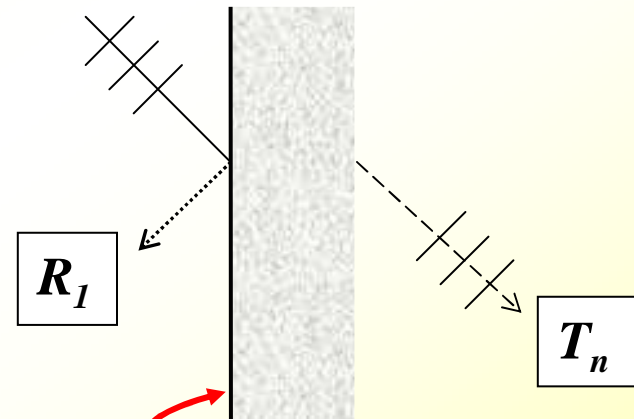
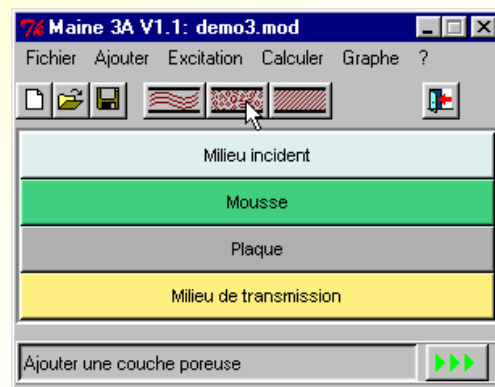
Introduction

Part I

Part II

Part III

Conclusion



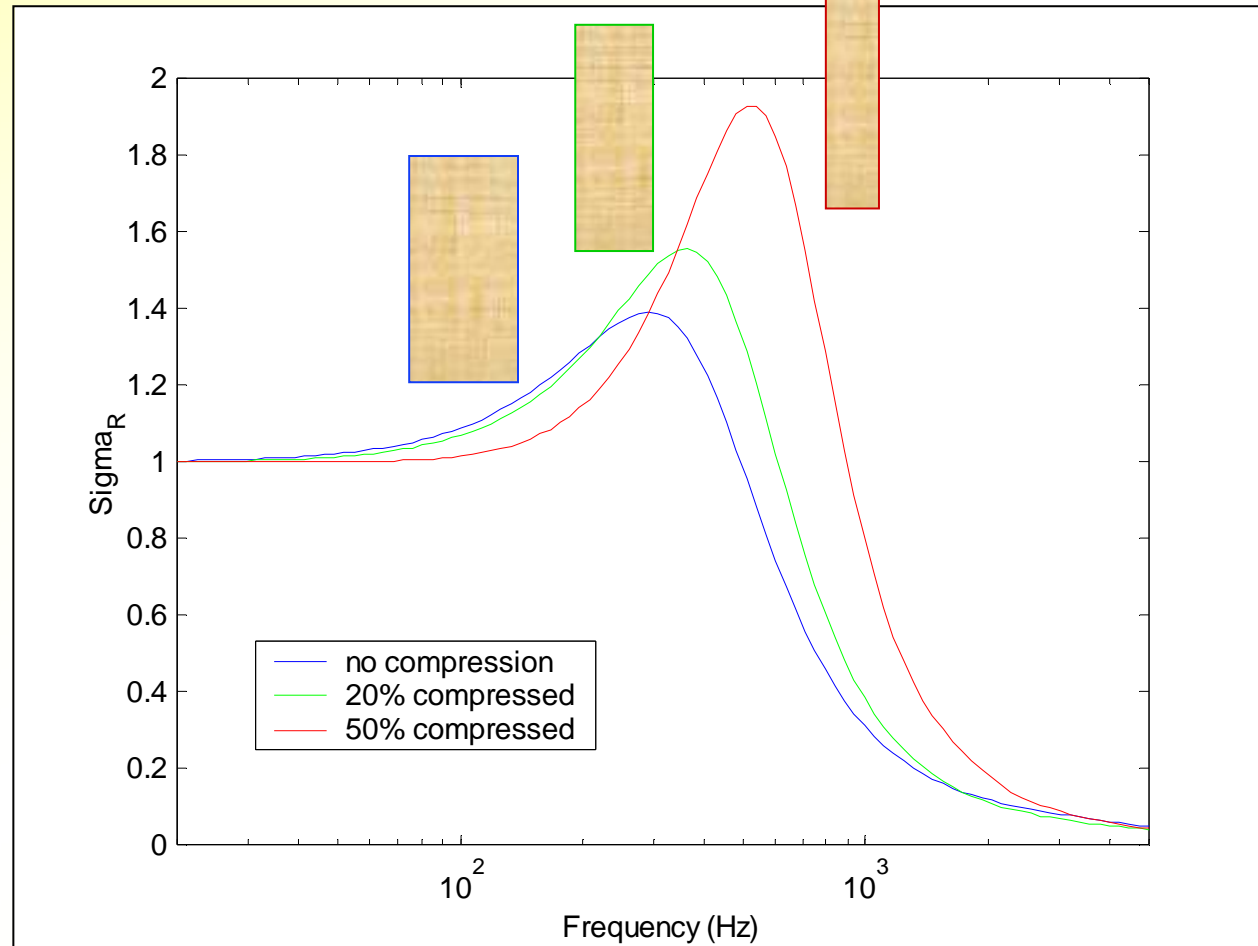
Impervious film

$$\Rightarrow u^s = u^f$$

$$\rightarrow Z_T = \frac{Z_0 T_n}{1 - R_1 - T_n}$$

Effect of fibrous material compression on the radiation

- Fibrous material compressed to **20%** and **50%**



Introduction

Part I

Part II

Part III

Conclusion

- Increase of the radiation with compression

Effect of a light film on the radiation

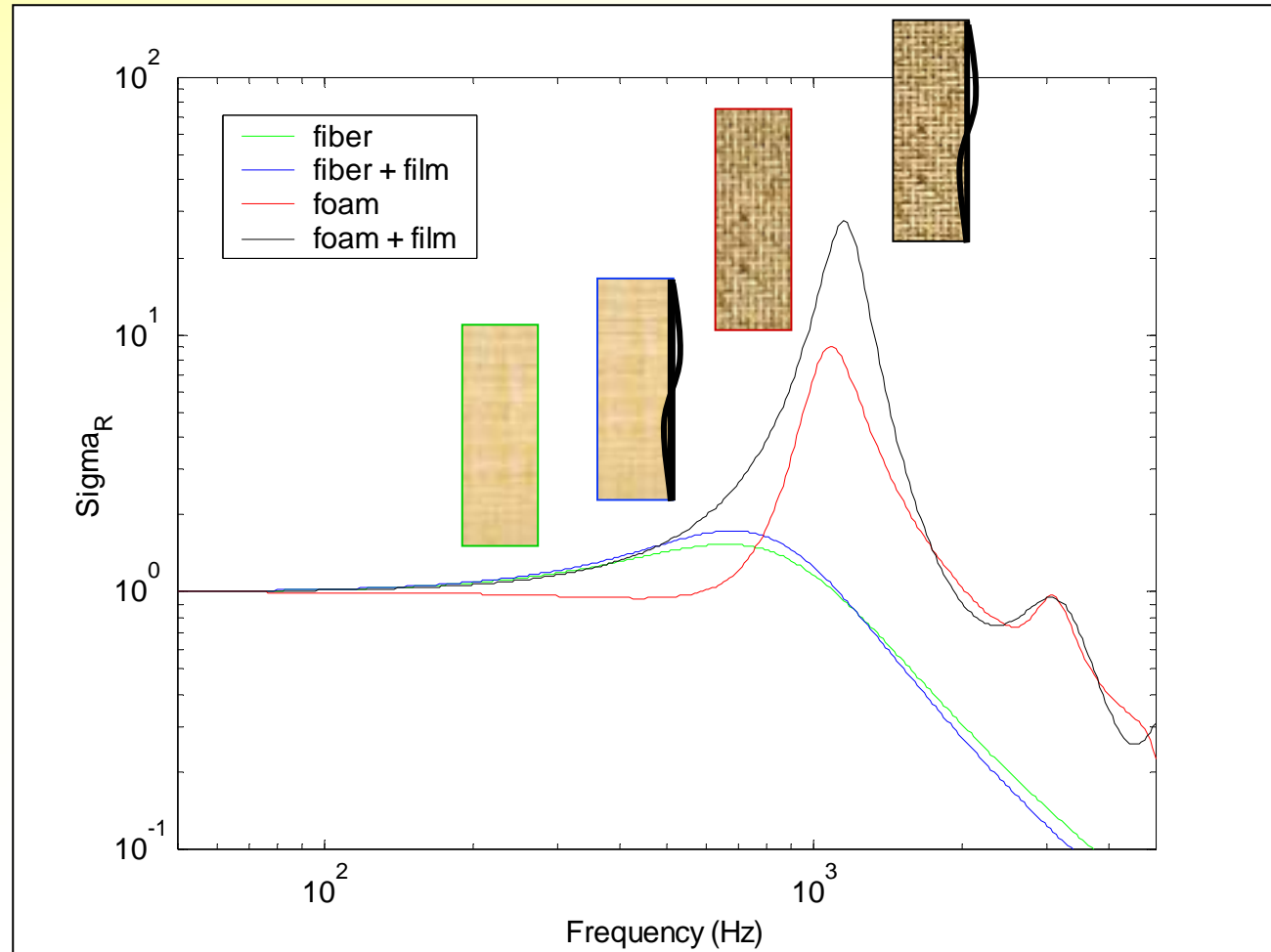
Introduction

Part I

Part II

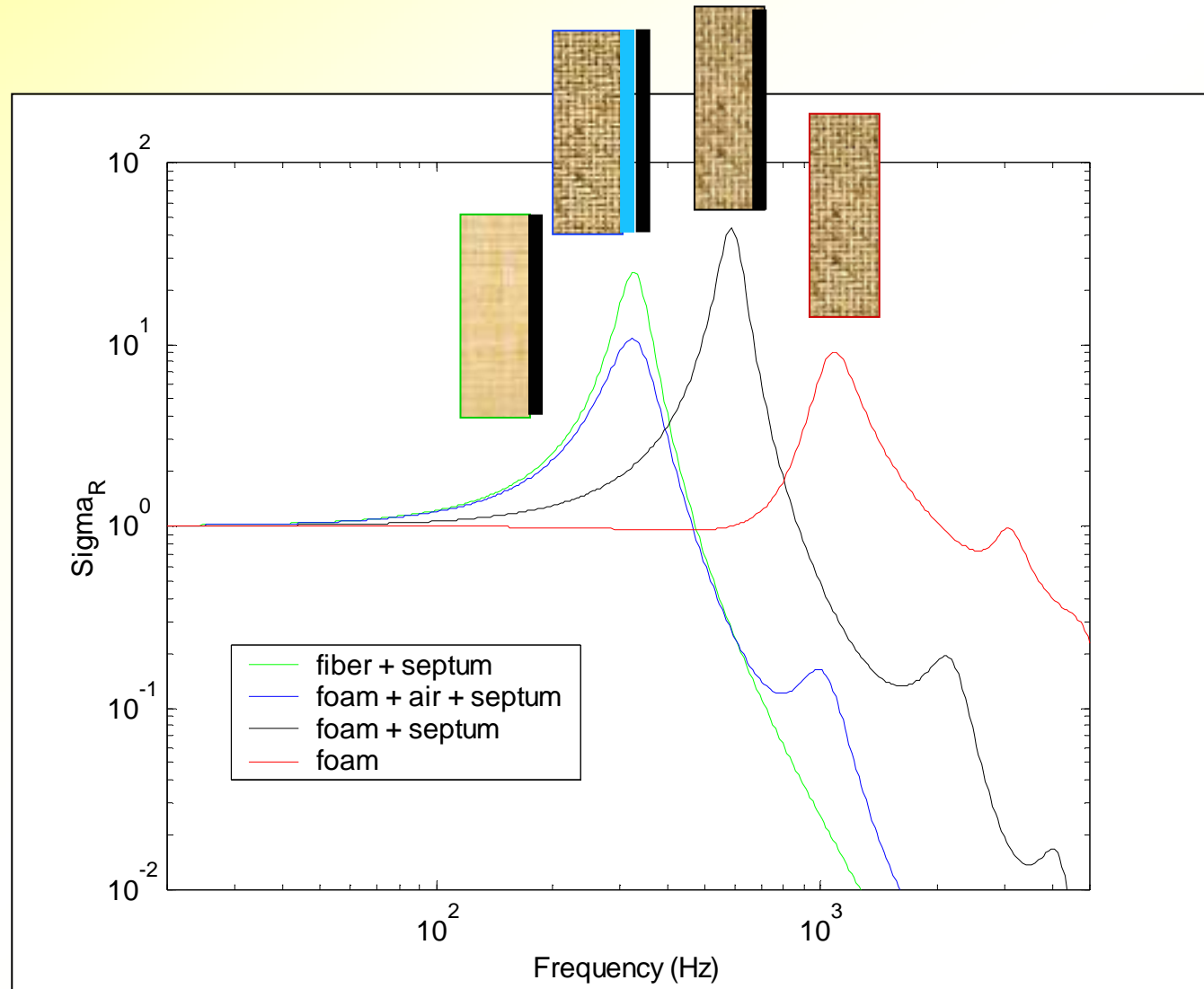
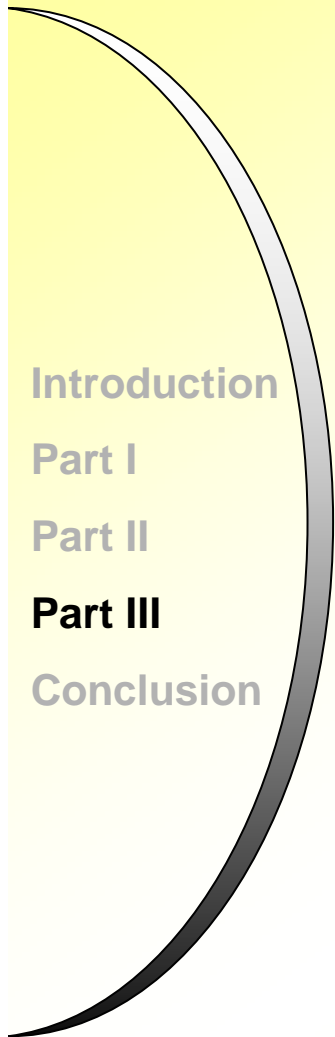
Part III

Conclusion



- **Increase of the radiation for the foam**
- **Almost no effect for fibrous**

Effect of a septum on the radiation



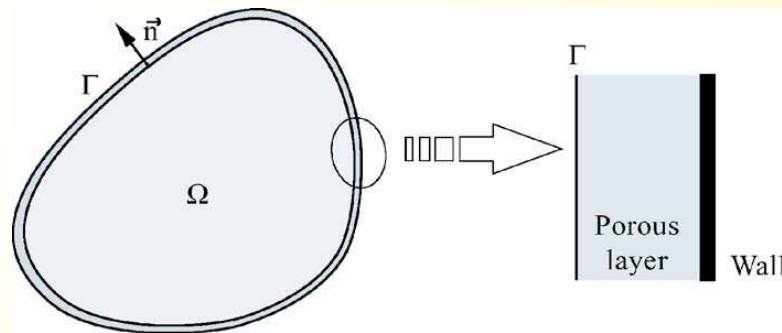
- **Decrease of the frequency**
- **Effect of air layer for the foam** → **skeleton bypass**

Conclusion

- **Acoustic radiation of a covered piston and plate**
➔ good agreements with measurements
- using **transfert impedance concept**
for porous materials **(NOT SURFACE IMPEDANCE !)**

Prospects

- **Effect of mounting conditions ?**
- **Absorption versus transmission coefficient ?**



Introduction

Part I

Part II

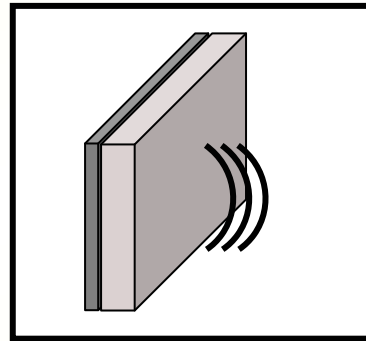
Part III

Conclusion

158th Meeting of the Acoustical Society of America
San Antonio, Texas

Acoustic radiation of a vibrating wall covered by a porous layer

Transfer impedance concept and effect of compression



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